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Additional Information

An experimental study of the effects of fuel properties on reactive spray evolution using Primary Reference Fuels

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Abstract

An experimental study on the ignition and combustion of Diesel-type sprays using n-heptane, iso-octane and four intermediate blends is presented. The choice of components was done in order to represent the transition from conventional diesel fuel (n-heptane) to a gasoline-like one (iso-octane) in terms of ignition behaviour. The experiments have been carried out in a high pressure high temperature vessel using specifications from the Engine Combustion Network (ECN). Parametric variations of oxygen concentration and air temperature have been performed for each fuel. In order to investigate the spray development, schlieren imaging for the quantification of spray penetration and ignition delay, OH* chemiluminescence imaging for the lift-off length, and broadband radiation imaging for the soot intensity and flame length have been applied. The results show the large effect of mixture reactivity on the ignition times and lift-off length values. Regarding the effect

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of the octane number of the blends on the ignition delay times, a linear effect has been found in the lower half of the blend range, while an exponential trend is evident in the top one. On the other hand, a scaling law for the stabilized flame length based upon momentum-controlled assumptions has shown that results are comparable to those obtained in the literature. Finally, the applicability of the results obtained on the performance and efficiency in real engines is discussed.

Keywords: Primary Reference Fuel, Diesel Spray, Ignition, Combustion, Lift-off Length

1 1. Introduction

The role of fuel properties on diesel engine performance and emission has been an active field of research for a long time. In recent years, more detailed studies in combustion vessels have helped understanding the fundamental processes dealing with fuel effects on spray development and combustion. Different fuel types are within reseach focus, such as bio-diesel, synthetic Fischer-Tropsch fuels, oxygenated fuels, diesel-gasoline blends and surrogate fuels, to mention a few. The end goal has been the matching of different combustion strategies and fuel types to improve the engine efficiency and reduce the emission of pollutants. For example, a study with direct diesel injection and port gasoline injection showed that the higher amount of gasoline retarded the combustion phasing. This in turn lowered the heat transferred and allowed for the lowest fuel consumption, which also resulted in low NO_X and PM emissions [1]. Another study replaced entirely the diesel fuel with gasoline and also obtained less specific fuel consumption together with lower

smoke levels [2]. Blends of gasoline and ethanol have also been used in diesel engines in order to study the effects on combustion and reduce the amount of diesel injected [3]. Results have shown reduction of soot levels, but at the cost of higher NO_X and compromised combustion stability under some conditions. Furthermore, a similar study with blends and changes in the contour conditions also presented improvements in the combustion efficiency [4].

Among the most used fuels used for the evaluation of diesel-gasoline 22 blends are those known as Primary Reference Fuels (PRFs), which represent both ends of the octane rating scale, namely n-heptane (ON = 0) and iso-octane (ON = 100). Such a difference in octane (and thus cetane) numbers makes them good representatives of diesel and gasoline, respectively, in terms of ignition and combustion. A very recent study [5] presents the ignition mechanism of n-heptane and iso-octane under various conditions, showing the expected increase of autoignition delay times when moving from n-heptane to iso-octane. Such results are also verified by experimental investigations [6], as the ignition delay of four different PRF blends was shown to increase as the composition of *iso*-octane augmented in the blends. Furthermore, engine tests also report an increase in the ignition delay times as the octane number of the fuel is larger [7, 8, 9]. Another work done with PRFs, diesel and gasoline, investigated the combustion recession after the end of injection [10]. The results showed that the fuels with highest cetane numbers presented the least lifted flames and the shortest combustion recession times. Analogously, the lift-off length and ignition delay time relationship was also studied under a wider range of conditions using conventional diesel and a PRF blend of similar properties [11]. Although the trend was not very

clear, there was a interaction implying that the condition with longest lift-off length also presented the longest ignition delay times.

The present work reports a fundamental investigation on the ignition 43 and combustion behaviour of fuel sprays under Spray A conditions [12], with a detailed focus on fuel properties. For that purpose, binary blends of nheptane and iso-octane, which are Primary Reference Fuels (PRFs), have been tested. In terms of ignition behaviour, such blends should be representative of a detailed transition from a conventional diesel fuel (n-heptane) to a gasoline-like one (iso-octane). The main objective of this work is the experimental characterization of spray mixing, ignition and sooting processes. For each fuel blend, the test plan includes parametric variations of ambient temperature and oxygen concentration, which have the largest effect on spray ignition behaviour. Experiments have been conducted in a constant pressure vessel, where high ambient pressure and temperature conditions typical of Diesel engines can be reproduced. For each condition, high speed schlieren and broadband luminosity imaging have been used to compare ignition, combustion and sooting behaviour of all fuel blends; OH* radical imaging has also been performed to measure lift-off length. The present contribution is structured as follows. After this introduction, both the experimental setup and the optical techniques employed will be presented. Test conditions will be summarized, and results of fuel effects will be analysed. Finally, the main conclusions from this study are drawn.

2. Experimental setup

64 2.1. High Temperature and High Pressure vessel

Tests have been performed in a high temperature and high pressure test chamber where the thermodynamic conditions obtained in a Diesel engine at the time of injection can be obtained with a maximum ambient temperature of 1000K and a maximum pressure of 15MPa. Compared to similar facilities [13], it is possible to obtain nearly quiescent and steady thermodynamic conditions in the test chamber. More details can be found in [14].

71 2.2. Optical setup

Different optical techniques have been employed in these experiments.

Schlieren imaging has been used to measure spray penetration and autoignition delay, broadband luminosity imaging has been used to assess sooting intensity, and OH* chemiluminescence imaging has been used to measure flame lift-off length. Figure 1 presents a schematic of the optical arrangement.

78 2.2.1. Schlieren imaging

The spray evolution inside the combustion chamber has been recorded by schlieren imaging [15]. This technique is sensitive to the first spatial derivative of density within the combustion chamber, which makes it useful to detect spray boundaries and thus evaluate macroscopic spray scales, whether vaporizing or non-vaporizing, inert or reactive. This technique shows the boundary between vaporized liquid and background gas because of the refractive index differences that exist between them, additionally, density gradients are also created in the chamber as the vaporized liquid cools the

ambient gas [16, 17]; such refractive index gradients are also present during combustion, as the high temperate creates low density regions. Therefore, this method is valid for inert and reactive conditions. For this technique, the spray has to be illuminated from one side by a collimated beam. The shadow produced by the spray is then gathered with a lens and at its focal length a diaphragm is positioned to produce the schlieren effect by eliminating the diverted light beams. The recorded image is then captured by a high speed camera in order to obtain a time resolved evolution of the spray.

An example of an schlieren image can be seen in Figure 2. The routine used for the processing was developed by Sandia National Laboratories as part of the ECN group and is available on-line [18]. The code is based on the successive calculation of two standard deviation images to remove the schlieren effect of the hot ambient gases and to detect the spray boundary. For every instant (I_t) , the processing routine subtracts the two preceding images $(I_{t-1} \& I_{t-2})$, then a two-step derivative process highlights the zones 101 where the pixels have changed due to either spray or background gas move-102 ment, then the image is segmented [19, 20, 21]. Once the boundaries have 103 been defined, the spray penetration is calculated by the procedure shown in [17]. Additionally, the ignition delay based upon schlieren images has also been calculated using the method described in [12, 16, 22], which relies 106 on the change in refractive index of the mixture and the shortening of the 107 penetration. 108

A schematic of the optical arrangement is shown in Figure 1. The spray has been illuminated from the left window and the light has been captured by a CMOS camera from the opposite side. The light from a Xenon lamp passed

through a 1mm diameter pinhole that simulates a single-point light source. The light was reflected on a 150mm parabolic mirror to obtain collimated light. After passing through the spray, the light has been collected by a biconvex lens and at its focal point a 4mm diaphragm has been positioned to produce the schlieren effect. A BG39 bandpass filter (360-580nm) was used to minimize soot radiation effects. A 10-bit Photron SA-5 CMOS high-speed camera equipped with a Nikon 50mm f=1:8 lens were used, image acquisition frequency was 42000FPS with exposure time of $4.18\mu s$ and a pixel/mm ratio of 5.26 leading to a field of view of 97mm.

2.2.2. Broadband radiation imaging

The soot distribution and intensity inside the combustion chamber have been recorded by direct imaging. This technique records the flame broadband radiation, which corresponds to the soot thermal radiation during the diffusion combustion phase.

Figure 3 shows a sample image of broadband radiation. An average back-126 ground image is calculated for each repetition and based on that image and a 127 two constant thresholds, a mask is generated to define the region of interest. 128 It is important to clarify that given the strong difference between the intensity levels in the soot onset area (in the vicinity of the lift-off) and further downstream, two threshold values (a lower one for the onset and a higher one 131 for the tip) were selected in order to accurately detect the contour. Then, the area of the flame and cumulative digital intensity are computed, and based 133 on those two values the specific intensity is obtained. Also, and in a similar way to the spray penetration, the soot onset length and soot penetration (or flame length) are obtained as the distance from the nozzle to the points

closest and furthest away from the injector, respectively. After all cases have been processed, the normalized intensity and specific intensity are calculated 138 taking into account the camera settings of exposure time and aperture. It is worth noting that the soot onset length is not a replacement of the lift-off length calculated by means of OH* chemiluminescence, because significant 141 amounts of soot do not form until downstream from the lift-off location [23]. 142 A schematic of the optical arrangement is shown in Figure 1 as well. A 143 50/50 beam splitter has been positioned in between the parabolic mirror and the left window with a 45° angle so the camera axis could be parallel to the spray axis but the reflected image perpendicular to it. A 10-bit Phantom 146 V12 CMOS high-speed camera equipped with a Zeiss 100mm f=2 lens were used, image acquisition frequency was 42000FPS with exposure times ranging from 0.4 to $23.29\mu s$ and a pixel/mm ratio of 7.23 leading of a field of view of 105mm.

$_{51}$ 2.2.3. OH^* Chemiluminescence imaging

The lift-off length (LoL) was recorded by OH* chemiluminescence imaging. This technique records radiation at 310nm, which is controlled by the OH* radical, and thus of the lift-off length for diesel sprays [23]. The OH* radical is a marker of the high temperature combustion [24], for example the one occurring at the lift-off length. Additionally, the band at 310nm is the strongest one, and therefore the best for determining the lift-off length [23, 24]. Although soot radiation could also be contributing to the recorded signal in the images, previous studies in the literature [23] also indicate that there should exist a spatial separation in both contributions. Therefore, while the downstream radiation will most probably be dominated by soot, the most

upstream one is assumed to be essentially related to OH* chemiluminiscence.

Figure 4 shows an image of OH* radiation. These images are similar in 163 appearance to the broadband ones, but with different spectral information. Additionally, these images have been taken with a long exposure, from the start of combustion until the end of injection ($\approx 3000 \mu s$); as opposed to the 166 broadband radiation where the exposure times were much shorter ($\approx 20\mu s$). 167 The field of view in this case is located closer to the injector, in order to be 168 able to measure the lift-off length with an adequate resolution. The processing algorithm for OH* images is similar, but lift-off length is calculated as the closest zone to the nozzle where flame can be found [23]. The algorithm 171 divides the image in two parts (upper and lower in Figure 4), based on the 172 injector axis, because the time averaged images show two bright bands of 173 light corresponding to the OH* chemiluminescence from the diffusion flame. Nevertheless, these bands do not imply that the combustion extends further on the edges of the spray, but instead, they correspond to line-of-sight averaging of the light from the flame [25, 26]. Carrying on with the processing, the values per row and side are added up so two curves Digital Level vs Axial Distance can be obtained. Finally, the lift-off is calculated as the average level between valley and first peak of each curve, as this gives the mean location of the turbulently fluctuating lift-off length [23]. 181

A schematic of the optical arrangement is also shown in Figure 1. Due to the lack of available optical accesses, the camera has been placed next to the parabolic mirror and at a small angle in order to image the spray properly.

A 16-bit Andor iStar ICCD intensified camera equipped with a 100mm focal length f=2 UV objective (by Bernhard Halle Nachfolger GmbH) and a 310nm

interference filter (FWHM = 10nm) was used to eliminate any additional radiation such as the soot luminosity. Because of the quasi-steady nature of the diffusion flame only one image was recorded per injection event, with an exposure time of 3ms starting 2ms after SoI and pixel/mm ratio of 5.85 leading to a field of view of 78mm.

92 3. Test conditions

Keeping in mind the objective of the present work, six different blends of n-heptane and iso-octane in increments of 20% have been used, including the pure components. The nomenclature of each blend is PRF#, where the # stands for the percentage in volume of iso-octane in the blend, being the rest n-heptane.

The so-called Spray A injector has been used; this is a standard within 198 the Engine Combustion Network (ECN) group, which is an international collaboration among different research laboratories in the world [12]. The 200 injector has a single-hole nozzle with a nominal diameter of $90\mu m$ (Serial 201 No. 210675). Reference injection pressure, chamber density, temperature 202 and oxygen mole fraction for Spray A are 150MPa, 22.8kg/m³, 900K and 15%, respectively. The only difference of the present study compared to the standard Spray A is the use of n-heptane/iso-octane blends, instead of the 205 standard n-dodecane. As previously described, the reactive spray has been 206 characterized in terms of spray tip penetration, lift-off length and soot radi-207 ation. Parametric variations of temperature and oxygen concentration have been performed as shown in Table 1. Additionally, for some experimental conditions (Table 2) the study was also carried out in an inert atmosphere,

i.e. with no oxygen in the ambient with the main purpose of having a reference for the comparison of spray tip evolution.

Figure 5 shows a sample of the scattering among repetitions based on the inert spray penetration, for each operation condition 15 repetitions were performed. It shows that all the values are within two standard deviations (blue lines), which should enclose approximately 95% of the values according to the two-sigma (2σ) rule of the Normal distribution. Since the tests have proven to be repetitive, the results shown throughout the rest of the section will only present the average values of all the repetitions (red line). For timeaveraged values such as lift-off lengths and ignition delay times, the error bars presented will also correspond to two standard deviations (95%).

22 4. Results

4.1. Analysis of the baseline case

Figure 6 shows a sequence of half-images that describes the evolution of the vapour penetration throughout the whole injection event for the inert and reactive sprays. For this example Spray A nominal conditions and PRF0 fuel have been selected.

From the start of injection (SoI), the spray penetrates as a clearly defined shadow under inert and reactive conditions. At around 500μ s after SoI, the tip of the reactive spray starts becoming transparent, this is due to the low temperature pre-reactions (also known as "cool flames") that precede the start of high temperature heat release. Such reactions result in the spray having similar refraction index as the surrounding air, and therefore becoming transparent [16, 27]. At 750μ s after SoI, the spray tip of the reactive

spray becomes visible again, even though some parts are still transparent; also a widening of the spray tip is evident, this indicates that the ignition has begun as the premixed combustion of the air-fuel mixture has caused a sudden increase in the spray volume due to the lower density of the products from the high temperature reactions taking place. After 1000μ s, the whole reactive spray shadow is visible again.

The reactive spray is narrow in the non-reacting region, closer to the noz-241 zle and up to the lift-off length; after this point its width increases abruptly due to combustion, which expands the original flow as explained in the previous paragraph. This difference can clearly be seen by comparing the shapes of 244 the inert and reactive sprays after the ignition has taken place and the spray is fully visible again $(1500\mu s)$. The inert spray maintains a constant angle 246 and a steadily increasing radius all the way to the transient tip, while the reactive one is very similar to the inert up to the lift-off length (which is 23mm for the present operating conditions), then the radius increases suddenly and the heat release region extends up to the spray tip. The combustion also causes the tip of the reactive spray to accelerate and penetrate at a faster 251 rate than its inert counterpart, which can be observed on the images from $1500\mu s$ onwards. This means that the subsequent spray evolution is also governed by the combustion process. The reactive spray structure remains 254 similar until the end of injection, with the spray tip penetrating steadily with time. 256

Figure 7 shows the broadband radiation images corresponding to the same spray evolution sequence, broadband luminosity is an indicator of the presence of soot in a reactive spray. At around 1000μ s after SoI combustion

has already started, as explained in the previous paragraph. However, soot is not visible yet up to this time, either due to the low flame temperature or because not enough soot is formed yet. After 1500μ s from the SoI, soot starts to become visible at the tip of the spray; and as the spray keeps on penetrating further, the sooting region also grows axially. Nevertheless, the soot region widths are narrower than those of the spray due to the fact that schlieren marks the whole spray region, while soot radiation mainly occurs within the stoichiometric diffusion flame front [25].

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Figure 8 (top) shows the average spray tip penetration for inert and reactive conditions and the flame length and soot onset length measured from image sequences such as the ones shown in Figures 2 and 3. From Figure 8 (top), it can be seen that the measured inert and reactive penetrations are close to one another until autoignition. Looking back at the $500\mu s$ image in Figure 6, corresponding to the cool flame period, one can already discern some differences, which become evident after the high-temperature start of combustion at $666\mu s$. After this point, the reactive spray enters the acceleration phase and its penetration separates from the the inert case. This can be appreciated in Figure 8 (bottom), where the ratio between the reacting and inert penetration gas been plotted. This variable has a peak at the SoC, it decreases slightly and then starts increasing steadily until it reaches a maximum value and stabilizes under the given conditions. Although for this nominal case the stabilized maximum is found close to the end of the observation window, the parametric variations of ambient temperature and oxygen in the following sections will show a similar behaviour, with a first well defined acceleration followed by a second more stable period. This means that after the acceleration period of the reactive spray has ended, the ratio between reactive/inert velocities is approximately constant and no longer increasing [28].

As explained earlier, the sooting region of the spray tends to reach a steady value after a certain time. This can be seen in Figure 8 (top), where the soot onset length (SOL) reaches a constant value fairly quickly while the flame length (FL) keeps on penetrating along with the vapour until it starts deviating to come to a steady length which under this nominal condition occurs close to 4000μ s. Although this is close to the end of the observation window, the subequent study of this FL under oxygen concentration variation (Figure 16) will confirm that also for the present nominal condition stabilization occurs.

The OH* lift-off length (LoL) is also shown in Figure 8 (top), but unlike
the other parameters shown, this value is not time resolved as explained in
Section 2. A remark worth to be made is the difference between the LoL and
SOL, while both parameters measure the position of the flame, their spectral
information is different. The broadband luminosity comes from the radiation
of the soot, as opposed to the chemiluminescence of the OH* radical. This
latter value is the best indicative that combustion is taking place within a
spray and is always present. The soot, on the other hand, depends strongly
on the conditions and fuel used, and may not be present even if the spray is
burning.

Summarizing, the reactive spray evolution can be broken down into four well defined stages. First, there is no combustion and its behaviour is like an inert spray (0-250 μ s). Second, pre-reactions start taking place, the spray

becomes transparent due to the changes in density, broadband radiation is not visible due to low temperatures or not enough soot formed $(250-500\mu s)$. 311 Third, $(500-2500\mu s)$ high temperature ignition happens. In a first moment, 312 the spray expands radially but not axially, but this is followed by the acceleration of the spray tip velocity compared to the inert case. Soot becomes 314 visible, if conditions are adequate, SOL starts rising and so does the flame 315 front. Fourth and last, the acceleration phase of the reacting case compared 316 to the inert one ends, and the spray reaches a quasi-steady period where, its 317 tip speed is similar to its inert counterpart (2500-4000 μ s). This four stages are consistent with observations of the authors in previous studies [28]. 319

20 4.2. Spray tip penetration

Air temperature has no effect at all on the non-reacting part of the spray 321 under reactive conditions, the density has kept at a constant value by adjusting the ambient pressure. This can be seen in Figure 9 (top) where all 323 penetration curves overlap the inert one up to $600\mu s$ approximately. At this point all reactive curves separate from the inert reference due to the prereactions occurring at the spray front, this is shown more clearly in Figure 9 (bottom) where the ratio between reactive and inert penetrations is depicted. 327 Once the ignition has taken place, the reactive spray will continue deviating 328 from the inert case at a higher rate during the acceleration phase; but when it reaches the quasi-steady phase, the acceleration will cease and the ratio between Reactive/Inert penetration will remain constant. It can be seen that 331 the higher the temperature of the air, the sooner the acceleration period will 332 begin and, therefore, the stabilization period will be reached earlier as the combustion is faster. As mentioned earlier, the temperature does not have

an effect on the non-reacting region, but will affect the ignition delay time and hence the point at which the acceleration of the tip commences. Concurrently, Figure 9 (top) shows that the higher air temperature causes the lift-off length to become shorter for the same injection pressure.

The oxygen concentration has an effect on the start of combustion of 339 a spray and it doesn't affect the mixing or evaporation of the non-reactive 340 phase. Figure 10 shows the comparison of the penetrations under the para-341 metric variation of oxygen. It can be seen that not only the reactivity of the air causes the curves to depart at different positions, but the more pronounced slopes as the oxygen fraction increases is an indicator of a combustion hap-344 pening more aggressively. Estimations of adiabatic flame temperature for a stoichiometric mixture, initially at the adiabatic mixing conditions, indicate that it changes from around 2337 K for the lowest oxygen case to 2524 K and 2681 K for the medium and highest cases. This approximate 350 K temperature variation with oxygen is much stronger than that observed when 340 changing the ambient air temperature, 2256, 2337 and 2405 K for the 800, 350 900 and 1000 K cases shown in Figure 9. Therefore it can be hypothesized 351 that the corresponding acceleration of the reacting vs. the inert case is dependent on the combustion-induced temperature variation, which will create a corresponding density drop. Finally, Figure 10 (top), also shows the lift-off 354 length to be shorter as the oxygen increases, this proves that the increased reactivity of the mixture is shifting the reactive region of the spray closer to the nozzle.

The fuel effect on the reactive penetration is very similar to that of the oxygen, because it will modify the reactivity of the mixture. The fuel has

very little effect on the inert spray [29, 30]; on the other hand, the reactive penetration shows a strong dependency on the fuel due to different ignition 361 delay times. Figure 11 (top) shows the penetration of the baseline case for all fuels and the inert penetration of PRF0 as a reference point. It can be seen that as the percentage of iso-octane increases (higher PRF), the 364 penetration starts deviating form the inert case at a later time and with a 365 less pronounced slope, and for some cases it doesn't seem to separate at all. 366 This last phenomenon could be due to two factors, either the ignition takes place much later than the time window studied, or because it doesn't occur at all. This also explains why the stabilized penetrations for all fuels do 369 not converge on the same values. Figure 11 (bottom) shows the penetration 370 ratios for all PRFs. In a similar trend to that of the oxygen, being that a 371 more reactive mixture will cause the curve to deviate earlier and faster, not all curves start deviating, meaning that the ignition has not been achieved yet. Moreover, the blends that burn do not seem to reach the same steady 374 value. Finally, the effect of the fuel on the lift-off length can be appreciated in Figure 11 (top), where the lower the percentage of n-heptane (lower PRF) the longer the lift-off length is. This is analogue with increasing the air temperature or the oxygen concentration, which results in a more reactive mixture and hence a shorter lift-off length. 379

The previous results also help at the understanding of lift-off stabilization mechanism. Although originally thought as based upon a flame front velocity balance, extensive evidences starting from [31] indicate that lift-off stabilizes based upon the autoignition of fuel elements being injected. In this particular case, the fuel composition has been swept from *n*-heptane to *iso*- octane, both of which have very similar laminar flame speeds [32] but highly
different autoignition properties. From a point of view of flame stabilization
by velocity equilibrium, both fuels should produce the same lift-off length.
However, experimental results presented here confirm that both lift-off and
autoignition delay change with operating conditions in a highly similar fashion, which indicates the leading role of autoignition on lift-off stabilization.

91 4.3. Ignition delay

During the study of the penetration it was observed that some parameters
had a stronger effect on the ignition delay. This was seen where the Reactive/Inert penetration ratio started deviating from unity. In this subsection,
the analysis of the ignition delay time under the parametric variations (Table
1) is presented.

The oxygen concentration appears to have an effect on the start of com-397 bustion. This is because while the same amount of air is being entrained, more oxygen is available, leading to a more reactive air-fuel mixture, so the ignition can happen sooner [33]. Figure 12 shows the effect of varying the 400 oxygen concentration for each blend tested. Altering the fuel reactivity and 401 the oxygen concentration at the same time has a combined effect, where the 402 higher PRF (more iso-octane) causes the differences in ignition delay to be 403 greater among oxygen concentrations. As mentioned previously, this combined effect is related to the reactivity of the air-fuel mixture and it seems to have a non-linear trend. 406

Air temperature has the strongest effect on the ignition delay than any of the other parameters. Figure 13 shows that the difference among 950K and 900K is smaller (around 200 μ s for PRF20) compared to the effect of

reducing the temperature from 900K to 800K (approx. 1000 μ s for PRF0).

The explanation for this non-linear effect is that at low temperatures the rate
of chemical reaction is slow enough to ignore the time required to form the
ignitable mixture, while at high temperatures the ignition delay chemistry
progresses very fast, and delay is mainly dominated by physical processes
[34]. Missing points for the low temperature cases for PRF60 to PRF100,
as well as the intermediate temperature for PRF100 are due to combustion
not occurring within the investigated conditions. Additionally, it can be seen
that the difference in ignition delay between 15% and 18% of oxygen is similar
to changing the temperature from 900K to 950K for PRF20.

20 4.4. Flame lift-off length

As was also seen during the analysis of the penetration, the boundary conditions not only affected the time at which the reactive curve started deviating from the inert one (ignition delay time), but also on the flame lift-off length, which will be analyzed in the present section.

The lift-off length seems to be dramatically affected by changes in the air temperature. This is very clear to see in Figure 14, where the lower the temperature the longer the lift-off length is going to be [23, 33, 35, 36, 31], but more important is that the effect is non-linear. A statistical regression fitting of the results shown in Figure 14 indicates that the temperature dependence of the lift-off length occurs with an exponent of -4.27, -4.57, -4.84 and -5.28 for PRF0 to PRF60, which indicates that the higher the fuel reactivity, the more important the temperature dependence is. The fitted lines are also presented in Figure 14 as dashed lines along with the expressions. Compared to a reference work in [31], with a temperature dependence of -3.74, the

temperature sensitivity is higher for all cases in the present study. However, the present results are coherent with previous results obtained in the same experimental facility (-5.0 to -5.4 as reported in [21]).

The oxygen concentration causes a similar effect to that of the air temperature. Figure 15 shows that the higher the oxygen concentration, the shorter the lift-off length is due to the higher reactivity of the mixture causing the spray to burn earlier (shorter ignition delay) and since the velocity of the spray is not altered (same injection pressure), this happens closer to the injector [23, 35].

Figures 14 and 15 also shows the effect of the fuel composition, and it can
be seen that it has a strong effect on the lift-off length. Like it was mentioned
earlier, the percentage of *iso*-octane on the blend reduces the reactivity of the
mixture tremendously, especially when the percentage is over 50. Moreover,
the results shown only present values up to PRF80, due to the fact that
the camera setup recorded the OH* chemiluminescence in a fixed temporal
window, and greater stabilization times of the lift-off length were not able
to be registered. Additionally, some conditions were not favourable for the
ignition of PRF100 such as 800K and 900K at 15% oxygen. In a similar
way as the effect on the ignition delay, the effects between 15% and 18% of
oxygen are similar to changing the temperature from 900K to 950K.

5 4.5. Stabilized flame length

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The length of the flame was measured by means of broadband radiation, but stabilized values are not achieved for all operating conditions due to either spatial or temporal limitations in the acquisition setup.

Radiation onset occurs after ignition delay. If resulting radiation is strong

enough to be detected, measured flame penetration is similar to that of the reacting spray tip penetration. Figure 16 shows the evolution of the flame penetration under the oxygen concentration variation. For all three cases it can be seen that the flame penetrates faster than the corresponding spray tip of the non-reacting references. Also, and according to what was seen in the study of the spray penetration, the deviation of the flame penetration occurs earlier and more aggressively as the oxygen concentration increases. Finally, the higher concentration of oxygen also causes the flame to stabilize closer to the nozzle. The reason for this behaviour is that the more reactive the air-fuel mixture the closer the stoichiometric surface will be to the injector, resulting in a shorter stabilized flame length.

Stabilized flame length measurements under Diesel engine conditions are rarely available in the literature, mainly due to the length, which is usually larger than most optical accesses, and the time needed for stabilization, which is quite long in terms of engine scales. Comparable results to the ones shown here are those reported in [37], which use long-exposure OH* chemilumines-cence images to quantify the stabilized FL for different nozzle diameters. To compare the present results with those measurements, the obtained FL values have been normalized for the different experimental conditions according to the following scaling law based on a momentum-controlled diffusion flame description:

$$FL_{norm} = FL \cdot \frac{z_{st}}{d_0 \cdot \sqrt{\frac{\rho_f}{\rho_a}} \cdot \sqrt{\frac{\rho_a}{\rho_{st,r}}}}$$
(1)

where $d_{eq} = d_0 \cdot \sqrt{\frac{\rho_f}{\rho_a}}$ is the equivalent diameter, which was already used

for the normalization under inert conditions, z_{st} is the stoichiometric mixture fraction, and $\sqrt{\frac{\rho_a}{\rho_{st,r}}}$ is a term that accounts for entrainment reduction due to the combustion-induced density decrease, where $\rho_{st,r}$ is the stoichiometric mixture density under reacting conditions. This term has been demonstrated to justify entrainment decrease under atmospheric pressure gas flames [38]. Figure 17 presents the results of the normalized flame length (FL_{norm}) , which shows the values to be in the range of 4-6 for most operating conditions, in agreement with those from [37]. Therefore, it confirms that the stabilization of the diffusion flame is controlled by mixing arguments, and not by fuel reactivity.

92 5. Conclusions

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The conclusions drawn from this investigation are as follows:

• The ambient temperature and oxygen concentration have very similar effects. They have no apparent effect on spray penetration during the inert spray evolution, but they do have an effect on the reacting spray evolution. As temperature rises or oxygen concentration increases, the reactivity of the mixture increases dramatically, so the fuel will start to burn earlier and closer to the nozzle, resulting in shorter ignition delay times and lift-off lengths. Where they have different effects is on the stabilized flame lengths, as a higher oxygen concentration will results in shorter lengths while the temperature does not have an effect on the stabilized flame. Finally, even though they have no effect on the inert vapour penetration, the shorter ignition delay times will cause the

start of the acceleration phase to start sooner and its slope to be more pronounced, therefore resulting in longer penetrations.

• Increasing the amount of *iso*-octane in the mixture has a completely opposite effect on the combustion behaviour in comparison to increasing oxygen content or temperature in the air. As *iso*-octane is less reactive than n-heptane, increasing its concentration will only cause the mixture to start to burn later and further away from the injector tip. Even though *iso*-octane has a slightly higher sooting tendency than n-heptane, the increase in lift-off causes much leaner mixtures, resulting in lower soot formation and a concurrent decrease in broadband radiation intensity. Two trends are defined regarding the octane number of the blend. On one hand, for the lower half of the octane scale (PRF0-20-40) the reactivity of the mixture (lift-off lengths and ignition delay times) decreases linearly with the fuel proportion, while on the the upper half (PRF60-80-100) a non linear trend can be seen.

A scaling law for the stabilized flame length based on momentum-controlled considerations has been validated for those conditions where such stabilization occurs. This scaling law agrees with similar results in the literature.

Finally, a small projection on the applicability of the present results for engine performance will be made. For that purpose, ignition delay and lift-off results (Figures 12 and 15) for PRF0 with 15% of oxygen and PRF40 with 21% of oxygen will be revisited. They could be taken as representative of a Diesel engine running with pure Diesel and an EGR level of 30% for a high load case, or with a Diesel-gasoline blend and no EGR. One can see that both cases have an ignition delay time of 600μ s and approximately

the same lift-off length, 23mm and 22mm. This would mean a similar combustion timing for both fuels, which is always interesting in terms of engine calibration. Furthermore, due to the higher oxygen amount in ambient air, the equivalence ratio at the lift-off for the gasoline (PRF40) would be much leaner, which should result in lower soot formation within the flame, and therefore lower exhaust particulates. The aforementioned comparisons agree with general trends presented in recent studies using primary reference fuels and similar strategies to reduce emissions while maintaining engine performance [2, 7, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48]. Nevertheless, further research on this topic is needed, to confirm the previous conclusion.

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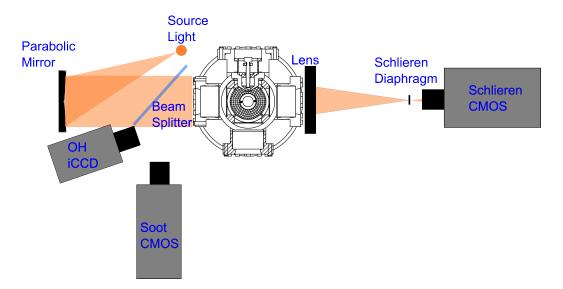


Figure 1: Schematic of the experimental setup for the study of the reactive spray

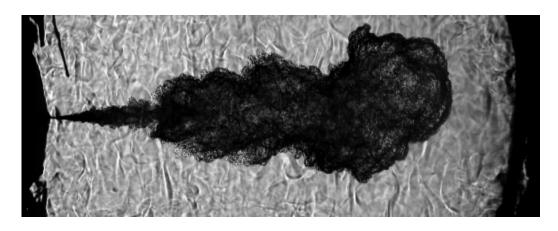


Figure 2: Sample image of the schlieren technique at $3000 \mu s$ after SoI



Figure 3: Sample of the broadband radiation imaging at $3000\mu s$ after SoI

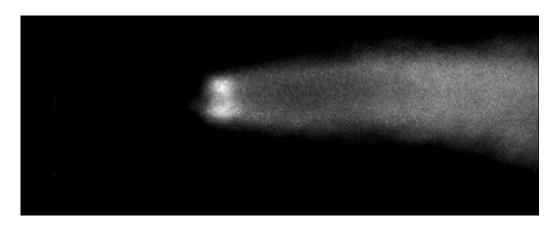


Figure 4: Sample image of the OH* chemiluminescence technique from $2000\mu s$ to $5000\mu s$ after SoI

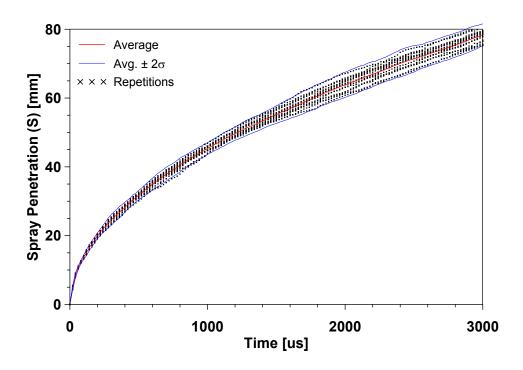


Figure 5: Spray penetration scattering of all 15 repetitions and average value, for PRF0 at 900K, $15\%O_2$ and 150MPa, time after SoI.

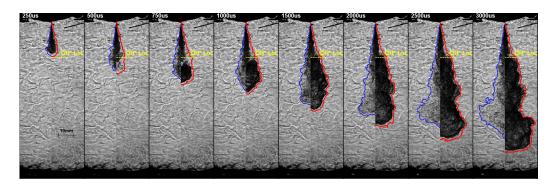


Figure 6: Evolution of inert (blue) and reactive (red) spray penetration for PRF0 at 900K, $15\%O_2$ and $150\mathrm{MPa},$ time after SoI.

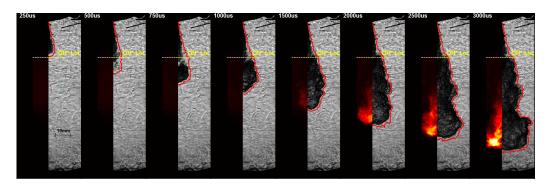


Figure 7: Evolution of soot luminosity and spray penetration simultaneous images for PRF0 at 900K, $15\%O_2$ and 150MPa, time after SoI.

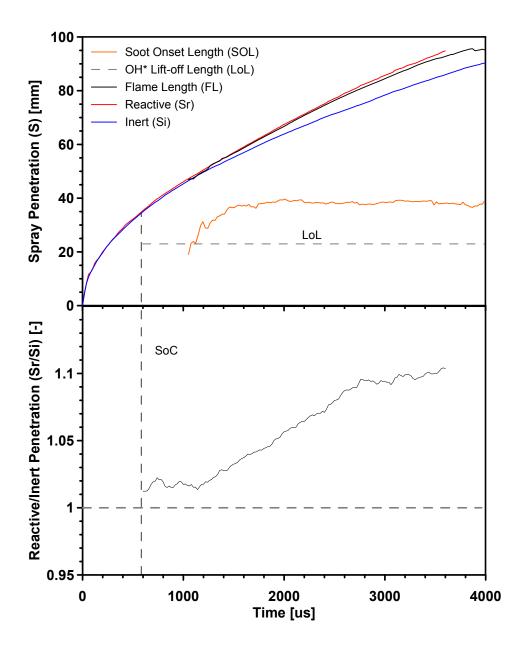


Figure 8: Time resolved evolution of the inert and reactive penetrations, along with flame length, soot onset length and lift-off length (top) and Reactive-Inert spray penetration ratio (bottom) for PRF0 at 900K, $15\%O_2$ and 150MPa, time after SoI.

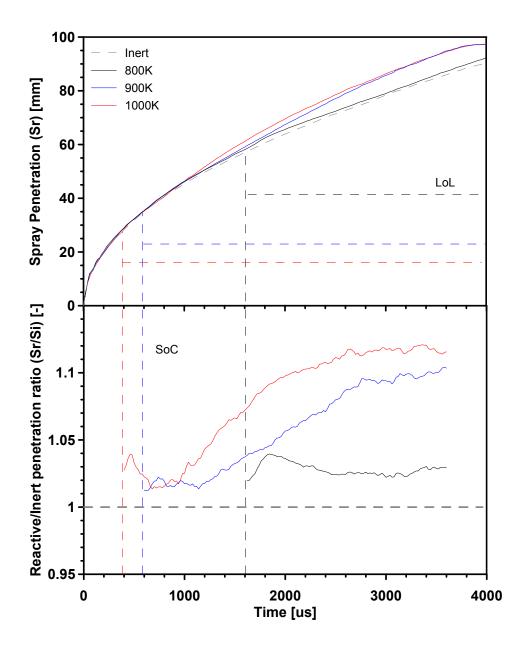


Figure 9: Reactive spray penetration (top) and Reactive/Inert spray penetration ratio (bottom) comparison for PRF0 under air temperature variation at $15\%O_2$, time after SoI.

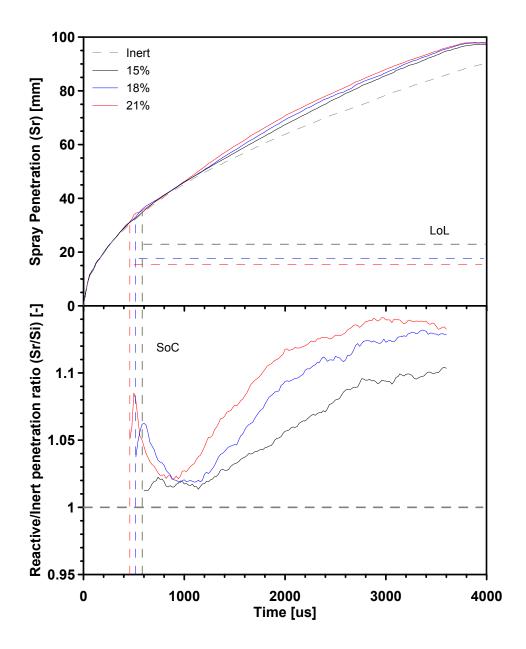


Figure 10: Reactive spray penetration (top) and Reactive/Inert spray penetration ratio (bottom) comparison for PRF0 under oxygen concentration variation at 900K, time after SoI.

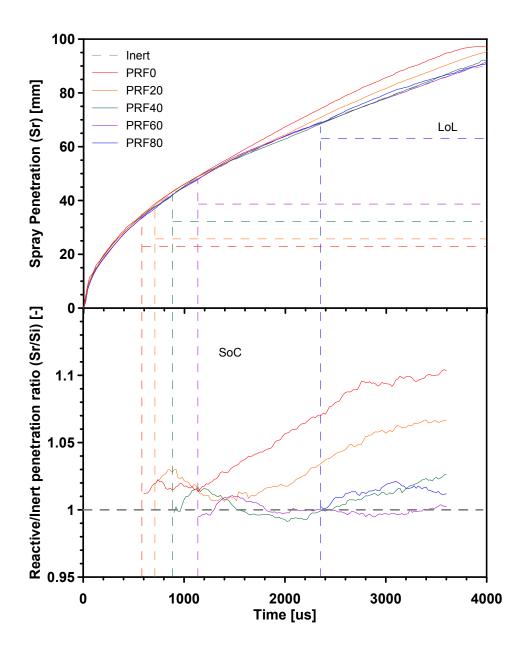


Figure 11: Reactive spray penetration (top) and Reactive/Inert spray penetration ratio (bottom) comparison for PRF0 to PRF100, at 900K and 15% O_2 , time after SoI.

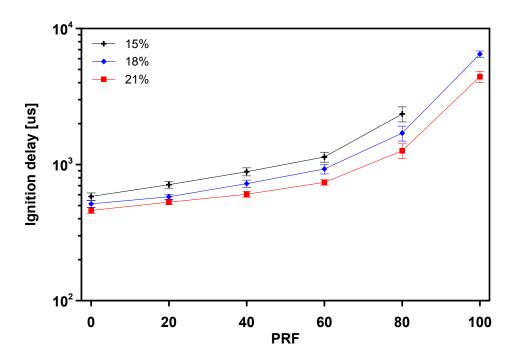


Figure 12: Ignition delay comparison under oxygen concentration variation, time after SoI

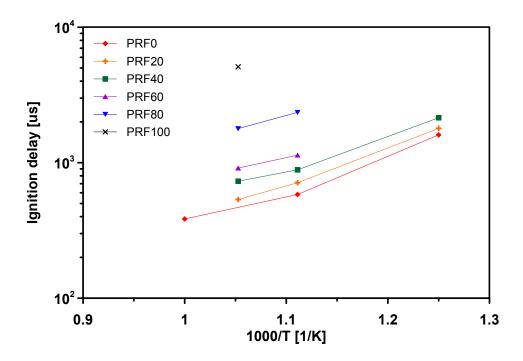


Figure 13: Ignition delay comparison under air temperature variation, time after SoI. Oxygen concentration 15%

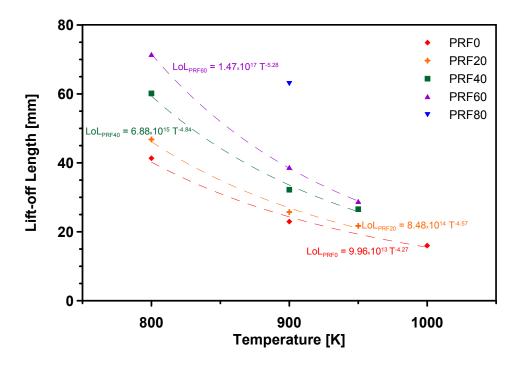


Figure 14: Lift-off length comparison under air temperature variation.

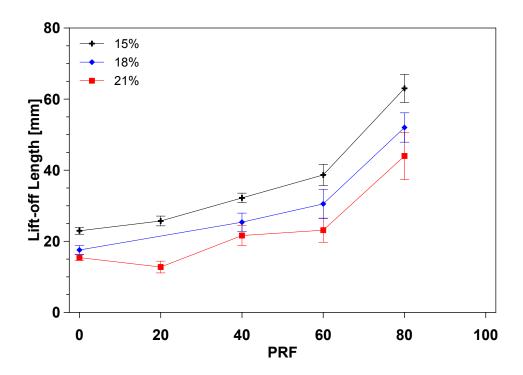


Figure 15: Lift-off length comparison under oxygen concentration variation.

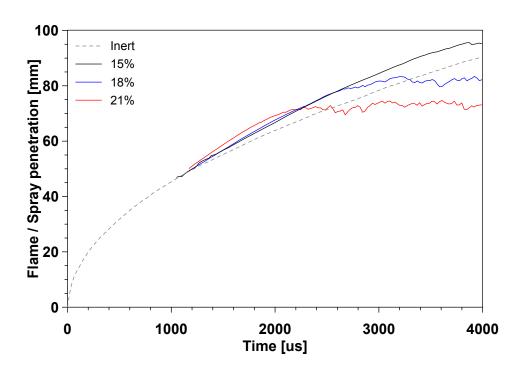


Figure 16: Flame length comparison for PRF0 under oxygen concentration variation at $900 \, \mathrm{K}$ with inert spray penetration, time after SoI.

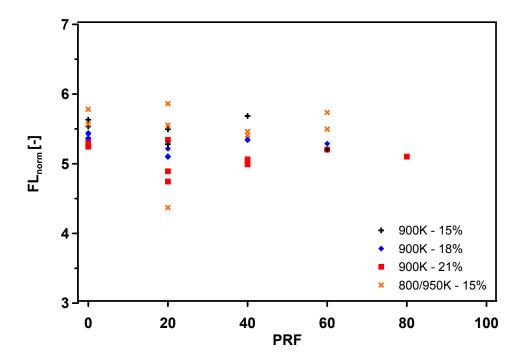


Figure 17: Normalized flame length (FL_{norm}) versus fuel composition at ambient temperature and oxygen concentration variations.

Table 1: Test matrix for parametric variation under reactive atmosphere. "A" indicates Spray-A standard conditions. All points performed at $150 \mathrm{MPa}$ of injection pressure and $22.8 \mathrm{kg/m^3}$ of ambient density. *For PRF0, high temperature tests were performed at $1000 \mathrm{K}$.

$O_2 \backslash T$	800K	900K	950K*
15%	X	A	X
18%		X	
21%		X	

Table 2: Test matrix for inert conditions. "A" indicates Spray-A standard conditions. All points performed at $150 \mathrm{MPa}$ of injection pressure and 0% of oxygen concentration.

$\rho \backslash T$	700K	800K	900K
$15.2kg/m^3$		X	
$22.8kg/m^3$	X	X	A
$45.6kg/m^{3}$		X	