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Additional Information

Individual wave overtopping volumes on mound breakwaters in breaking wave conditions and gentle sea bottoms

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Abstract:

Mound breakwaters are usually designed to limit the mean wave overtopping rate (q) or the maximum individual wave overtopping volume (V_{max}). However, rarely do studies focus on wave overtopping volumes on breakwaters in depth-limited breaking wave conditions. This study analyzes 2D physical tests on mound breakwaters with relevant overtopping rates ($0.33 \leq R_c/H_{m0} \leq 2.83$) and three armor layers (Cubipod®-1L, rock-2L and cube-2L) in depth-limited breaking wave conditions ($0.20 \leq H_{m0}/h_s \leq 0.90$) and with two bottom slopes ($m = 1/25$ and $m = 1/50$). The 2-parameter Weibull distribution was used to estimate $V_{max}^* = V_{max}/(g H_{m0} T_{01}^2)$ with coefficient of determination $R^2 = 83.3\%$. In this study, the bottom slope ($m = 1/50$ and $m = 1/25$) did not significantly influence V_{max} or the number of overtopping events, N_{ow} . During the design phase of a mound breakwater, q is required to use the methods given in the literature to estimate V_{max} . Thus, q must be estimated for design purposes when measured q is not available. In this study, CLASH Neural Network (CLASH NN) was used to estimate q with $R^2 = 63.6\%$. If the 2-parameter Weibull distribution proposed in this study is used to estimate V_{max} with q estimated using CLASH NN, the prediction error of V_{max}^* is $R^2 = 61.7\%$. With the method presented in this study, the ratio between estimated and measured V_{max}^* falls within the range 1/2 to 2 (90% error band) when q is estimated with CLASH NN. The new estimators derived in this study provide good predictions of N_{ow} and V_{max} with a method simpler than those in the literature on overtopped mound breakwaters in depth-limited breaking wave conditions on gentle sea bottoms ($1/50 \leq m \leq 1/25$).

Keywords: mound breakwater; wave overtopping; individual wave overtopping volumes; depth-limited breaking wave conditions; bottom slope; proportion of overtopping events

1. Introduction

Crest elevation is a key parameter when designing mound breakwaters due to its direct effect on construction costs as well as visual and environmental impact. Climate change effects (e.g., sea level rise) and increasing social concern about the visual impact of coastal structures are leading to reductions in crest freeboards and increases in the overtopping hazard. In this situation, coastal structure designs with reduced crest freeboards and relevant overtopping discharges become significant. In addition, most mound breakwaters are built in the surf zone, where they are attacked by waves breaking on the sea bottom.

Tolerable mean overtopping discharges, q ($\text{m}^3/\text{s}/\text{m}$), are commonly considered to design crest elevation of coastal structures. However, the mean individual overtopping volume (\bar{V}) may be much lower than the maximum individual wave overtopping volume, V_{max} (m^3/m). For this reason, Franco et al. (1994) suggested that overtopping hazard should be directly related with individual wave overtopping events, rather than the mean overtopping rate.

Several prediction methods exist to estimate q (e.g. Molines and Medina, 2015a and EurOtop, 2018), the number of overtopping events ($N_{ow} = N_w P_{ow}$) and V_{max} (see Molines et al., 2019) on mound

46 breakwaters in non-breaking conditions. Victor et al. (2012) conducted 2D physical tests on smooth
 47 impermeable structures under depth-limited breaking wave conditions with horizontal bottom slope
 48 and concluded that wave breaking had a significant impact on N_{ow} and V_{max} . Gallach (2018) carried
 49 out 2D physical tests on smooth impermeable steep sloped structures in depth-limited breaking wave
 50 conditions and two bottom slopes ($m = 0$ and $m = 1/100$). However, Gallach (2018) did not find a
 51 significant effect of the breaking waves conditions on V_{max} . Nørgaard et al. (2014) performed 2D
 52 physical tests on rubble mound breakwaters ($\cot\alpha = 1.5$, where $\cot \alpha$ is the armor slope) in depth-
 53 limited breaking wave conditions with horizontal bottom slope. They evaluated the performance of
 54 the existing formulations, valid for non-breaking wave conditions, for observations in breaking wave
 55 conditions and concluded that existing formulas were underpredicting N_{ow} and overpredicting V_{max} .
 56 Therefore, the depth-limited breaking wave conditions of the incoming waves may be a significant
 57 factor to consider.

58 The bottom slope highly affects the type of wave breaking at the toe of the structure. Herrera et
 59 al. (2017) pointed out that bottom slope directly influences mound breakwater design; the optimum
 60 point where wave characteristics are estimated needs to be determined when in depth-limited
 61 breaking wave conditions. Mares-Nasarre et al. (2020) found a significant effect of bottom slope on
 62 the overtopping layer thickness (h_c) and the overtopping flow velocity (u_c). Mares-Nasarre et al. (2020)
 63 also determined that the optimum point to estimate wave characteristics to calculate h_c and u_c was
 64 located at a distance of 3 times the water depth from the toe of the structure. However, the influence
 65 of the bottom slope on V_{max} has not yet been analyzed.

66 This research focuses on the distribution of individual wave overtopping volumes $F(V)$ in depth-
 67 limited breaking wave conditions for mound breakwaters and the influence of bottom slope on V_{max} .
 68 This paper is organized as follows. In section 2, the literature on individual wave overtopping
 69 volumes is examined. Neither the optimum point to estimate wave characteristics nor the effect of
 70 bottom slope on N_{ow} and V_{max} was assessed by the studies in the literature. Section 3 describes the
 71 experimental setup with two bottom slopes ($m = 1/25$ and $m = 1/50$) and the experimental data
 72 analysis. 2D small-scale tests on mound breakwaters in depth-limited breaking wave conditions (0.20
 73 $\leq H_{m0}/h_s \leq 0.90$) and three armor layers (Cubipod®-1L, rock-2L and cube-2L) were conducted. Section
 74 4 assesses existing estimators for N_{ow} and V_{max} . None of the existing estimators for mound breakwaters
 75 satisfactory describes N_{ow} for very low q . In section 5, the optimum point to estimate wave
 76 characteristics when calculating N_{ow} is determined and a new N_{ow} estimator is developed. In section
 77 6, the 2-parameter Weibull distribution is considered to fit $F(V)$; the influence of bottom slope on the
 78 two-parameter Weibull distribution is also investigated. The quadratic utility function proposed by
 79 Molines et al. (2019) is used in this study to take into account the higher relevance of the largest
 80 individual wave overtopping volumes for practical applications. In section 7, the performance of the
 81 new N_{ow} and V_{max} estimators is validated using q estimators given in the literature. Finally, in section
 82 8, conclusions are drawn.

83 2. Literature review

84 2.1. Individual wave overtopping volumes

85 Van der Meer and Janssen (1994) and Franco et al. (1994) first introduced the Weibull
 86 distribution to describe individual wave overtopping volumes for dikes, and vertical and composite
 87 breakwaters, respectively. Later, the 2-parameter Weibull distribution was proposed by different
 88 authors (e.g., Besley, 1999 or Victor et al., 2012) to analyze individual wave overtopping volumes in
 89 a variety of coastal structures. The 2-parameter Weibull distribution is given by

$$F(V) = F(x \leq V) = 1 - \exp\left[-\left(\frac{V}{a}\right)^b\right] \quad (1)$$

90 where $F(x \leq V)$ is the non-exceedance probability of the individual wave overtopping volume per
 91 wave, x is the individual wave overtopping volume, a is the dimensional scale factor and b is the
 92 shape factor. Eq. (1) can also be found as:

$$F(V) = 1 - \exp\left[-\left(\frac{V/\bar{V}}{A}\right)^b\right] \quad (2)$$

93 where $A = a/\bar{V}$ is the scale factor and \bar{V} is the measured mean individual wave overtopping volume.
 94 If all the measured data were used for the analysis and they followed a perfect Weibull
 95 distribution, the mean individual wave overtopping volume, \bar{V} , would be equal to the mean value
 96 of the Weibull distribution, μ ($\mu = \bar{V}$). Under the previous hypothesis, a relationship between A and
 97 b exists and is described by

$$A = \frac{a}{\bar{V}} = \frac{1}{\Gamma\left(1 + \frac{1}{b}\right)} \quad (3)$$

98 where Γ is the gamma function, given by $\Gamma(z) = \int_0^{\infty} t^{z-1} e^{-t} dt$. Van der Meer and Janssen (1994) and
 99 Franco et al. (1994) recommended a value of $b = 0.75$ for dikes, and vertical and composite
 100 breakwaters, respectively, which corresponds to $A = 0.84$ according to Eq. (3).

101 Besley et al. (1999) studied individual wave overtopping volumes for sloped structures, vertical
 102 walls and composite breakwaters. These authors also referred to the results reported by Franco et al.
 103 (1996), who highlighted the influence of wave steepness on shape factor b for vertical walls. Franco
 104 et al. (1996) also noticed that the shape factor b was around 0.1 higher for sloping structures than for
 105 vertical walls. Regarding sloped structures, Besley et al. (1999) recommended values for the shape
 106 factor b as a function of the offshore wave steepness, $s_{op} = 2\pi H_{s0}/(gT_{p0}^2)$, where H_{s0} is the significant
 107 offshore wave height and T_{p0} is the deep water peak period. These authors suggested $b = 0.76$ for $s_{op} =$
 108 0.02 and $b = 0.92$ for $s_{op} = 0.04$.

109 Bruce et al. (2009) carried out 2D physical tests on mound breakwaters with horizontal bottoms,
 110 $0.8 \leq R_c/H_{m0} \leq 1.3$ and $0.33 \leq H_{m0}/h_s \leq 0.40$, where R_c is the crest freeboard, $H_{m0} = 4(m_0)^{1/2}$ is the significant
 111 wave height, and h_s is the water depth at the toe of the structure. These authors tested a wide variety
 112 of armor units, both double- and single-layer armors, and analyzed the individual wave overtopping
 113 volumes higher than \bar{V} . Bruce et al. (2009) suggested a shape factor $b = 0.74$ and concluded that no
 114 significant differences could be observed between the different armor units.

115 Victor et al. (2012) investigated individual wave overtopping volumes on smooth impermeable
 116 steep slopes ($0.36 \leq \cot \alpha \leq 2.75$) with horizontal bottoms and $0.11 \leq R_c/H_{m0} \leq 1.69$. The authors observed
 117 that the wave heights during the tests with large H_{m0} did not fit a Rayleigh distribution ($0.04 \leq H_{m0}/h_s$
 118 ≤ 0.37), but a Composite Weibull distribution, and they concluded that deviations were caused by
 119 depth-induced breaking of the largest waves. Victor et al. (2012) also observed that this wave
 120 breaking process limited the value of the maximum individual wave overtopping volumes and
 121 decreased the shape factor b . Moreover, these authors investigated the effect of the relative crest
 122 freeboard, R_c/H_{m0} , slope angle, α , and s_{op} . They concluded that the effect of s_{op} ($0.012 \leq s_{op} \leq 0.041$) was
 123 negligible and proposed Eq. (4) to estimate the shape factor b considering the individual wave
 124 overtopping volumes higher than \bar{V} .

$$b = \exp\left(-2.0 \frac{R_c}{H_{m0}}\right) + (0.56 + 0.15 \cot \alpha) \quad (4)$$

125 Zanutigh et al. (2013) analyzed the shape factor b for rough and smooth low-crested structures
 126 ($0 \leq R_c/H_{m0} \leq 2$), using the individual wave overtopping volumes higher than \bar{V} . These authors
 127 reported higher scatter in the data for rubble mound breakwaters than in the data for smooth slopes.
 128 Zanutigh et al. (2013) also pointed out that, even if formulas considering the dimensionless crest
 129 freeboard (R_c/H_{m0}) gave good results for smooth structures, they were not adequate for rubble mound
 130 breakwaters. The shape factor b for rubble mound breakwaters was found to be related to a
 131 dimensionless mean wave overtopping discharge, $Q^{**} = q/(g H_{m0} T_{m-1,0})$ (where $T_{m-1,0} = \frac{m-1}{m_0}$ and m_k is
 132 the k -th spectral moment, $m_k = \int_0^{\infty} S(f) f^k df$, and $S(f)$ is the wave spectrum) similar to $Q^{**} = q/(g H_{m0}$
 133 $T_{01})$ (where $T_{01} = \frac{m_0}{m_1}$) proposed by Besley (1999), and Eq. (5) was derived. EurOtop (2018) also
 134 recommends Eq. (5) for estimating the shape factor b for armored rubble slopes and mounds.

$$b = 0.85 + 1500 Q^{**1.3} \quad (5)$$

135 Nørgaard et al. (2014) conducted 2D physical tests on rock-armored mound breakwaters with
 136 crown wall both in non-breaking and breaking wave conditions ($0.18 \leq H_{m0}/h_s \leq 0.50$) with horizontal
 137 bottoms and $0.9 \leq R_c/H_{m0} \leq 2.0$. These authors assessed the existing formulas in the literature for the
 138 shape factor b in non-breaking wave conditions and concluded that they were overpredicting the
 139 largest overtopping wave volumes in depth-limited breaking wave conditions. Nørgaard et al. (2014)
 140 proposed Eq. (6) based on 30% of the highest individual wave overtopping volumes.

$$b = \begin{cases} 0.75 & \text{for } \frac{H_{m0}}{H_{1/10}} \leq 0.848 \text{ or } \frac{H_{m0}}{h_s} \leq 0.2 \\ -6.1 + 8.08 \frac{H_{m0}}{H_{1/10}} & \text{for } \frac{H_{m0}}{H_{1/10}} > 0.848 \text{ and } \frac{H_{m0}}{h_s} > 0.2 \end{cases} \quad (6)$$

141 where $H_{1/10}$ is the average of 10% of the highest waves in the test run.

142 Gallach (2018) carried out thousands of 2D physical tests using bottom slopes $m = 0$ and $m =$
 143 $1/100$ for steep slopes and vertical structures in a wide range of crest freeboards ($0.0 \leq R_c/H_{m0} \leq 3.25$).
 144 The author investigated the effect of depth-limited breaking wave conditions ($0.03 \leq H_{m0}/h_s \leq 0.50$) on
 145 the shape factor b and found it negligible, contrary to results published by Victor et al. (2012) and
 146 Nørgaard et al. (2014). Gallach (2018) also noticed that the shape factor b was not affected by the
 147 roughness of the structure and proposed a new formula to estimate b as function of R_c/H_{m0} and the
 148 structure slope, using the largest 10% individual wave overtopping volumes. Regarding the scale
 149 factor A , Gallach (2018) found that the best fit values were significantly different than those given by
 150 Eq. (3).

151 Molines et al. (2019) analyzed the 2D physical tests conducted by Smolka et al. (2009) on
 152 conventional mound breakwaters ($1.25 \leq R_c/H_{m0} \leq 4.78$) with crown wall in non-breaking conditions
 153 ($0.10 \leq H_{m0}/h_s \leq 0.32$) and reported the inconsistencies in the selection criteria of the number of
 154 overtopping events used to fit the scale and shape factors identified by Pan et al. (2016). Molines et
 155 al. (2019) compared the fitting of A and b of the 2-parameter Weibull distribution using 10%, 30%
 156 50%, and 100% (with quadratic utility function) of the highest individual wave overtopping volumes.
 157 Utility functions are used to consider the relative relevance of the observed data; using the whole
 158 dataset with a quadratic utility function, all the observations are used but special attention is paid to
 159 the highest volumes. The relationship between A and b was not given any more by Eq. (3). Note that
 160 small overtopping events significantly affect \bar{V} and N_{ov} ; the estimations of A based on Eq. (3) are
 161 sensitive to small overtopping events which are not significant for practical applications. The shape
 162 factor, b , is given as function of the dimensionless mean wave overtopping discharge, $Q^* = q/(g H_{m0}$
 163 $T_{01})$, whereas the scale factor A depends on the shape factor b , as shown in Eqs. (7) and (8),
 164 respectively, when applying the quadratic utility function to all observed individual wave
 165 overtopping volumes.

$$b = 0.63 + 1.25 \exp(-3.0 \cdot 10^5 Q^*) \quad (7)$$

$$A = 1.4 - 0.4 \frac{1}{b} \quad (8)$$

166 Additionally, Molines et al. (2019) proposed a 2-parameter Exponential distribution to describe
 167 individual wave overtopping volumes, given by

$$F(V) = 1 - \exp\left[-\left(\frac{V/\bar{V} - C}{D}\right)\right] \quad (9)$$

168 where

$$D = 2.6 - 2.6 \exp(-3.0 \cdot 10^5 Q^*) \quad (10)$$

$$C = 1.2 - D - 0.2 D^2 \quad (11)$$

169 2.2. Number of overtopping events

170 In order to assign an exceedance probability to every individual wave overtopping volume,
 171 Makkonen (2006) recommended the Weibull plotting position formula, given by

$$F(V) = 1 - \frac{i}{N_{ow} + 1} \quad (12)$$

172 where $F(V)$ is the exceedance probability of the individual wave overtopping volume per wave, i is
 173 the rank of the individual volume, sorted in descending order ($i = 1$ corresponds to V_{max}) and N_{ow} is
 174 the number of overtopping events.

175 Lykke-Andersen et al. (2009) applied Eq. (12) to rewrite the Weibull distribution function as:

$$V_i = A\bar{V} \left[-\ln\left(\frac{i}{N_{ow} + 1}\right) \right]^{1/b} = A\bar{V} [\ln(N_{ow} + 1) - \ln(i)]^{1/b} \quad \text{with } i = 1 \text{ to } N_{ow} \quad (13)$$

176

177 By setting $i = 1$ in Eq. (13), V_{max} can be calculated as

$$V_{max} = A\bar{V} [\ln(N_{ow} + 1)]^{1/b} \quad (14)$$

178 Besley (1999), EurOtop (2007) and EurOtop (2018) proposed Eq. (15), which uses N_{ow} instead of
 179 $N_{ow} + 1$.

$$V_{max} = A\bar{V} [\ln(N_{ow})]^{1/b} \quad (15)$$

180 Lykke-Andersen et al. (2009) warned that Eq. (15) would predict an inconsistent $V_{max} = 0$ for N_{ow}
 181 $= 1$. To estimate V_{max} , not only N_{ow} has to be estimated, using either Eq. (14) or Eq. (15), but also the
 182 mean individual wave overtopping volume ($\bar{V} = q T_{01} N_w / N_{ow}$, where q is the mean overtopping
 183 discharge). Thus, q has to be estimated in order to calculate V_{max} . To this end, Besley (1999) proposed
 184 Eqs. (16) and (17) for simple slopes, and complex slope structures with return walls or berms,
 185 respectively.

$$P_{ow} = \frac{N_{ow}}{N_w} = \exp \left[-K_1 \left(\frac{R_c}{T_{01} \sqrt{g H_{m0}} \gamma_f} \frac{1}{\gamma_f} \right)^2 \right] \quad (16)$$

186

$$\begin{cases} P_{ow} = 55.4 Q^{*0.634} & \text{for } 0 < Q^* < 8 \cdot 10^{-4} \\ P_{ow} = 2.5 Q^{*0.199} & \text{for } 8 \cdot 10^{-4} < Q^* < 10^{-2} \\ P_{ow} = 1 & \text{for } Q^* > 10^{-2} \end{cases} \quad (17)$$

187 where P_{ow} is the proportion of overtopping waves, N_{ow} is the number of overtopping events and N_w
 188 is the number of incoming waves, γ_f is the roughness factor, H_{m0} is the spectral significant wave
 189 height, $Q^* = q/(g H_{m0} T_{01})$ and q are the dimensionless and dimensional mean overtopping discharges,
 190 respectively. Besley (1999) recommended $K_1 = 37.8$ for structure slope $\cot \alpha = 2$ and $K_1 = 63.8$ for $\cot \alpha$
 191 $= 1$. Besley (1999) proposed Eq. (18) to estimate q .

$$\frac{q}{g T_{01} H_{m0}} = K_2 \exp \left[-K_3 \frac{R_c}{T_{01} \sqrt{g H_{m0}} \gamma_f} \frac{1}{\gamma_f} \right] \quad (18)$$

192 where K_2 and K_3 are experimental coefficients function of α . For $\cot \alpha = 1.5$, $K_2 = 8.84 \cdot 10^{-5}$ and $K_3 = 19.9$.

193 Nørgaard et al. (2014) proposed a variation of Eq. (19) to extend the application of this equation
 194 to depth-limited breaking wave conditions for $0.006 \leq P_{ow} \leq 0.120$ and $7.3 \cdot 10^{-7} \leq Q^* \leq 6.2 \cdot 10^{-5}$, given by

$$N_{ow} = C_1 \text{ Eq. (17)}$$

$$C_1 = \begin{cases} 1 & \text{for } H_{m0}/H_{1/10} \leq 0.848 \text{ or } H_{m0}/h_s \leq 0.2 \\ -6.65 + 9.02 \frac{H_{m0}}{H_{1/10}} & \text{for } H_{m0}/H_{1/10} > 0.848 \text{ and } H_{m0}/h_s > 0.2 \end{cases} \quad (19)$$

195 Nørgaard et al. (2014) recommended using CLASH Neural Network (CLASH NN), described
 196 by Van Gent et al. (2007), for q estimation.

197 EurOtop (2018) recommended Eq. (20) for mound breakwaters with permeable crest berms.

$$P_{ow} = \exp \left[- \left(\sqrt{-\ln 0.02} \frac{R_c}{Ru_{2\%}} \right)^2 \right] \quad (20)$$

198 where $Ru_{2\%}$ is the wave run-up height exceeded by 2% of the incoming waves, calculated as

$$\frac{Ru_{2\%}}{H_{m0}} = 1.65 \gamma_f \gamma_\beta \gamma_b \xi_{-1,0} \quad (21a)$$

199 with a maximum value of

$$\frac{Ru_{2\%}}{H_{m0}} = \min \left(1.00 \gamma_{f \text{surging}} \gamma_\beta \left[4.00 - \frac{1.50}{\sqrt{\xi_{-1,0}}} \right], 2.0 \right) \quad (21b)$$

200 where $\gamma_{f \text{surging}} = \gamma_f + (I_{r-1,0} - 1.8)(1 - \gamma_f)/8.2$; γ_β , the oblique wave attack factor and $\xi_{-1,0} =$
 201 $\tan \alpha / \sqrt{2\pi H_{m0} / (g T_{m-1,0}^2)}$.

202 EurOtop (2018) suggested Eq. (22) to estimate q .

$$\frac{q}{\sqrt{g H_{m0}^3}} = 0.09 \exp \left(-1.5 \frac{R_c}{H_{m0} \gamma_f \gamma_\beta} \right) \quad (22)$$

203

204 Molines et al. (2019) recently proposed to estimate the proportion of overtopping waves, P_{ow} ,
 205 valid for $0.001 \leq P_{ow} \leq 0.20$ and $7.0 \cdot 10^{-8} \leq Q^* \leq 6.4 \cdot 10^{-5}$, using

$$P_{ow} = 480 Q^{*0.8} \quad (23)$$

206 Similar to Nørgaard et al. (2014), Molines et al. (2019) recommended using CLASH NN for
 207 estimating q . Table 1 summarizes the experimental ranges of the methods found in the literature.

Author	Structure	Crown wall	R_c/H_{m0} [-]	H_{m0}/h_s [-]	m [-]	$\cot \alpha$ [-]
Bruce et al. (2009)	Mound breakwaters	Yes	0.80 – 1.03	0.33 – 0.40	0	2
Victor et al. (2012)	Smooth impermeable steep slopes	No	0.11 – 1.69	0.04 – 0.37	0	0.36 – 2.75
Zanuttigh et al. (2013)	Smooth slopes and rubble mound breakwaters	-	0 – 2	-	-	2 – 4
Nørgaard et al. (2014)	Rubble mound breakwaters	Yes	0.9 – 2	0.18 – 0.50	0	1.5
Gallach (2018)	Steep slopes and vertical structures	No	0 – 3.25	0.03 – 0.50	0, 1/100	0 – 0.27, 1.5 – 2.75
Molines et al. (2019)	Mound breakwaters	Yes	1.25 – 4.78	0.10 – 0.32	0	1.5

208 Table 1. Summary of the experimental ranges of the methods to estimate V_{max} in the literature.

209 Table 2 presents a summary of the methods in literature to estimate V_{max} on mound breakwaters or
 210 permeable slopes.

Author	q [m ³ /s/m]	P_{ow} [-]	b [-]	A [-]	V_{max} [l/m]
--------	---------------------------	--------------	---------	---------	-----------------

Besley (1999)	$g T_{01} H_{m0} K_2 \exp \left[-K_3 \frac{R_c}{T_{01} \sqrt{g H_{m0}}} \right]$	<p>Simple slopes:</p> $\exp \left[-K_1 \left(\frac{R_c}{T_{01} \sqrt{g H_{m0}}} \frac{1}{\gamma_f} \right)^2 \right]$ <p>Complex slopes:</p> $\begin{cases} 55.4 Q^{*0.634} & 0 < Q^* < 8 \cdot 10^{-4} \\ 2.5 Q^{*0.199} & 8 \cdot 10^{-4} < Q^* < 10^{-2} \\ 1 & Q^* > 10^{-2} \end{cases}$	0.76 for $s_{op} = 0.02$ 0.92 for $s_{op} = 0.04$	$\frac{1}{\Gamma \left(1 + \frac{1}{b} \right)}$	$A \bar{V} [\ln(N_{ow})]^{1/b}$
EurOtop (2018)	$0.09 \exp \left(-1.5 \frac{R_c}{H_{m0} \gamma_f \gamma_\beta} \right) \sqrt{g H_{m0}^3}$	$\exp \left[- \left(\sqrt{-\ln 0.02} \frac{R_c}{R u_{2\%}} \right)^2 \right]$	$0.85 + 1500 Q^{*1.3}$	$\frac{1}{\Gamma \left(1 + \frac{1}{b} \right)}$	$A \bar{V} [\ln(N_{ow})]^{1/b}$
Nørgaard et al. (2014)	CLASH NN	$C_1 \times (\text{Besley, 1999; compl. sl.})$ Where: $C_1 = -6.65 + 9.02 \frac{H_{m0}}{H_{1/10}}^{*2}$	$-6.1 + 8.08 \frac{H_{m0}^{*3}}{H_{1/10}}$	$\frac{1}{\Gamma \left(1 + \frac{1}{b} \right)}$	$A \bar{V} [\ln(N_{ow} + 1)]^{1/b}$
Molines et al. (2019)	CLASH NN	$480 Q^{*0.8}$	$0.63 + 1.25 \exp(-3 \cdot 10^5 Q^*)$	$1.4 - 0.4 \frac{1}{b}$	$A \bar{V} [\ln(N_{ow} + 1)]^{1/b}$

Notes:

*1 K_2 and K_3 are empirical coefficients function of α . Here, for $\cot(\alpha)=1.5$, $K_2 = 8.84 \cdot 10^{-5}$ and $K_3 = 19.9$

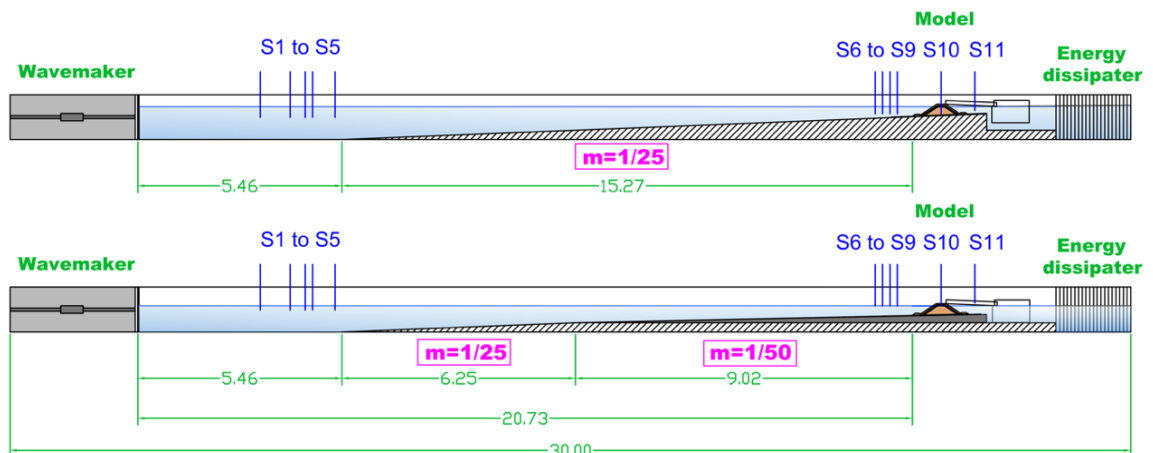
*2 for $\frac{H_{m0}}{H_1} > 0.848$ and $\frac{H_{m0}}{h_s} > 0.2$; otherwise $C_1=1$.

*3 for $\frac{H_{m0}}{H_1} > 0.848$ and $\frac{H_{m0}}{h_s} > 0.2$; otherwise $b=0.75$.

211 Table 2. Summary of the methods in the literature to estimate V_{max} for mound breakwaters or slopes
212 structures.

213 3.1. Experimental setup

214 2D physical tests were conducted in the wave flume (30.0 m × 1.2 m × 1.2 m) of the Laboratory
215 of Ports and Coasts at the *Universitat Politècnica de València* (LPC-UPV), with two bottom slope
216 configurations and a piston-type wave maker. The first configuration involved a continuous ramp of
217 4% slope ($m = 1/25$) all along the flume. The second configuration was formed by two ramps: a 6.3 m-
218 long $m = 1/25$ bottom slope, and a 9.0 m-long $m = 1/50$ bottom slope. Figure 1 shows the longitudinal
219 cross-sections of the LPC-UPV wave flume for both configurations as well as the locations of the wave
220 gauges.



221
222 Figure 1. Longitudinal cross-section of the LPC-UPV wave flume (dimensions in meters).
223

224 In order to measure the water surface elevation, 11 capacitive wave gauges were placed along
225 the flume. Wave gauges S1 to S5 were located in the wave generation zone, and these were used to

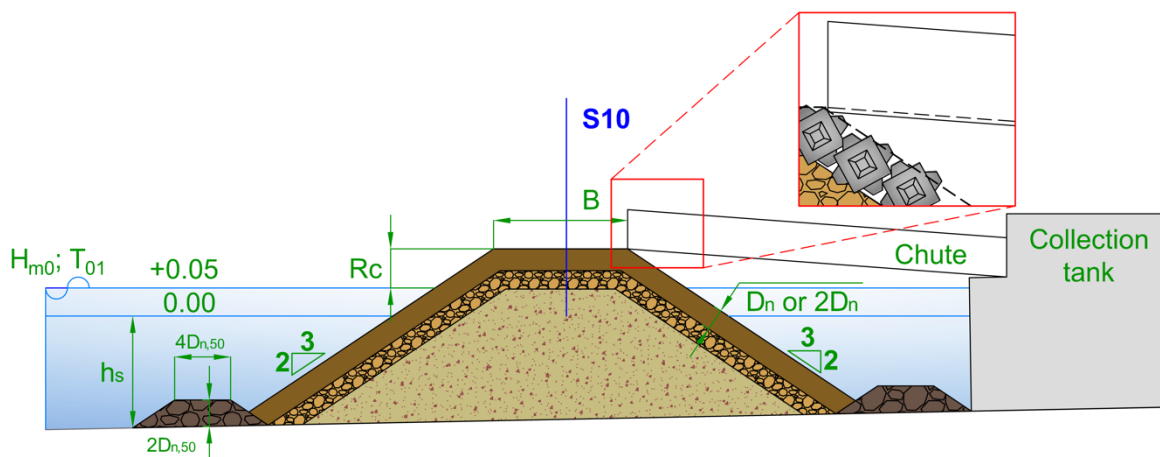
226 separate incident and reflected waves in the wave generation zone using the LASA-V method (see
 227 Figueres and Medina, 2004). Wave gauges S6 to S9 were placed close to the model, where depth-
 228 limited wave breaking occurs and existing methods to separate incident and reflected waves are not
 229 reliable. The distances from the structure toe to S6, S7, S8 and S9 were modified with the water depth
 230 at the toe of the structure, h_s . S6, S7, S8 and S9 were located at distances $5h_s$, $4h_s$, $3h_s$ and $2h_s$ from the
 231 toe of the breakwater, respectively, following the recommendations given by Herrera and Medina
 232 (2015). Wave gauge S10 was placed in the middle of the structure crest, and S11 was located behind
 233 the model.

234 Random wave runs of 1,000 waves were generated following a JONSWAP spectrum ($\gamma=3.3$). The
 235 AWACS wave absorption system was activated to avoid multireflections. Neither low-frequency
 236 oscillations nor piling-up (wave gauge S11) were significant during the tests. Piling-up consists of an
 237 increase of the water depth around the model caused by the accumulation of water when high
 238 overtopping rates occur. The LPC-UPV wave flume allows the water to be recirculated through a
 239 double floor of 25 cm to prevent it.

240 Test series were associated to the water depth at the toe of the model (h_s). For each h_s , T_p and H_{m0}
 241 $= 4(m_0)^{0.5}$ at the wave generation zone were calculated so as to keep the wave steepness approximately
 242 constant throughout each test series ($s_{op} = 0.018$ and 0.049). For each s_{op} , H_{m0} at the wave generation
 243 zone ($H_{m0,g}$) was increased in steps of 1 cm from no damage to failure of the armor layer or wave
 244 breaking at the wave generation zone. The water depths at the toe of the structure were $h_s = 20$ and
 245 25 cm for the tests carried out with a bottom slope $m = 1/25$. The water depths were $h_s = 20$ and 25 cm
 246 for the single-layer Cubipod® and double-layer rock armored models with a bottom slope $m = 1/50$.
 247 The water depths were $h_s = 25$ and 30 cm for the double-layer cube armored model with a bottom
 248 slope $m = 1/50$.

249 Due to the importance of crest freeboard on overtopping, two corrections were made: (1) the
 250 natural evaporation and facility leakages during the tests and (2) the extracted accumulated
 251 overtopping volumes during the working day (overtopping volumes in the collection tank were
 252 pumped out of the flume). These lead to a small increase in the crest freeboard along time of the order
 253 of 10 mm for a long working day.

254 The tested breakwater model corresponds to a mound breakwater with $\cot\alpha = 1.5$ slope and toe
 255 berms (see Figure 2). Three armor layers were tested: a single-layer Cubipod® armor, a double-layer
 256 rock armor and a double-layer randomly-placed cube armor. The nominal diameters or equivalent
 257 cube sizes were: $D_n = 3.79$ cm (Cubipod®-1L), $D_n = 3.18$ cm (rock-2L) and $D_n = 3.97$ cm (cube-2L).
 258 Tests with the bottom slope $m = 1/50$ were conducted with a medium size rock toe berm ($D_{n,50} = 2.6$
 259 cm), while tests with the bottom slope $m = 1/25$ were carried out with a larger rock toe berm ($D_{n,50} =$
 260 3.9 cm) to guarantee the toe berm stability. The range of the variables in the test is shown in Table 3.
 261 Note that wave conditions in the model zone are estimated using the SwanOne model (see Verhagen
 262 et al., 2008), as explained in section 3.2.



263
 264 Figure 2. Cross-section of the breakwater model tested in LPC-UPV wave flume (dimensions in
 265 meters).

m	Armor	#tests	B [m]	h_s [m]	R_c [m]	$H_{m0,g}$ [m]	H_{m0} [m]	T_{01} [s]
1/50	CC-1L	47	0.24	0.20	0.12	0.06 – 0.21	0.06 – 0.15	0.83 – 1.97
				0.25	0.07	0.06 – 0.21	0.06 – 0.16	0.87 – 2.02
	CB-2L	45	0.27	0.25	0.11	0.06 – 0.19	0.05 – 0.16	0.86 – 2.03
				0.30	0.06	0.06 – 0.20	0.06 – 0.18	0.83 – 1.88
	CE-2L	13	0.26	0.20	0.15	0.06 – 0.13	0.06 – 0.12	0.81 – 1.66
0.25				0.10	0.06 – 0.13	0.06 – 0.13	0.82 – 1.71	
1/25	CC-1L	46	0.24	0.20	0.12	0.07 – 0.21	0.06 – 0.15	0.84 – 1.65
				0.25	0.07	0.07 – 0.21	0.06 – 0.18	0.82 – 2.11
	CB-2L	47	0.27	0.20	0.11	0.06 – 0.20	0.05 – 0.16	0.87 – 1.69
				0.25	0.06	0.06 – 0.21	0.06 – 0.17	0.87 – 2.12
	CE-2L	21	0.26	0.20	0.15	0.06 – 0.16	0.05 – 0.14	0.84 – 1.86
0.25				0.10	0.06 – 0.13	0.05 – 0.14	0.80 – 1.88	

268 Table 3. Dimensions and wave conditions at the toe of the structure in 2D physical tests at the
269 LPC-UPV wave flume: CC-1L, CB-2L and CE-2L correspond to Cubipod® - 1L, cube-2L and rock-2L
270 armors.

271 Overtopping discharges were measured using a weighing system placed in a collection tank
272 behind the model during each test. Overtopping was collected using a chute in the rear side line of
273 the crest. The inner border of the base of the chute was aligned with the armor layer to prevent too
274 much wave overtopping losses. Individual wave overtopping volumes were identified following the
275 method developed by Molines et al. (2019), based on a continuous record of accumulated overtopping
276 volume. This method uses the derivative of the overtopping record to identify the overtopping
277 volumes. Figure 3 shows a photo of the experimental set up with the Cubipod®-1L armored
278 breakwater model.



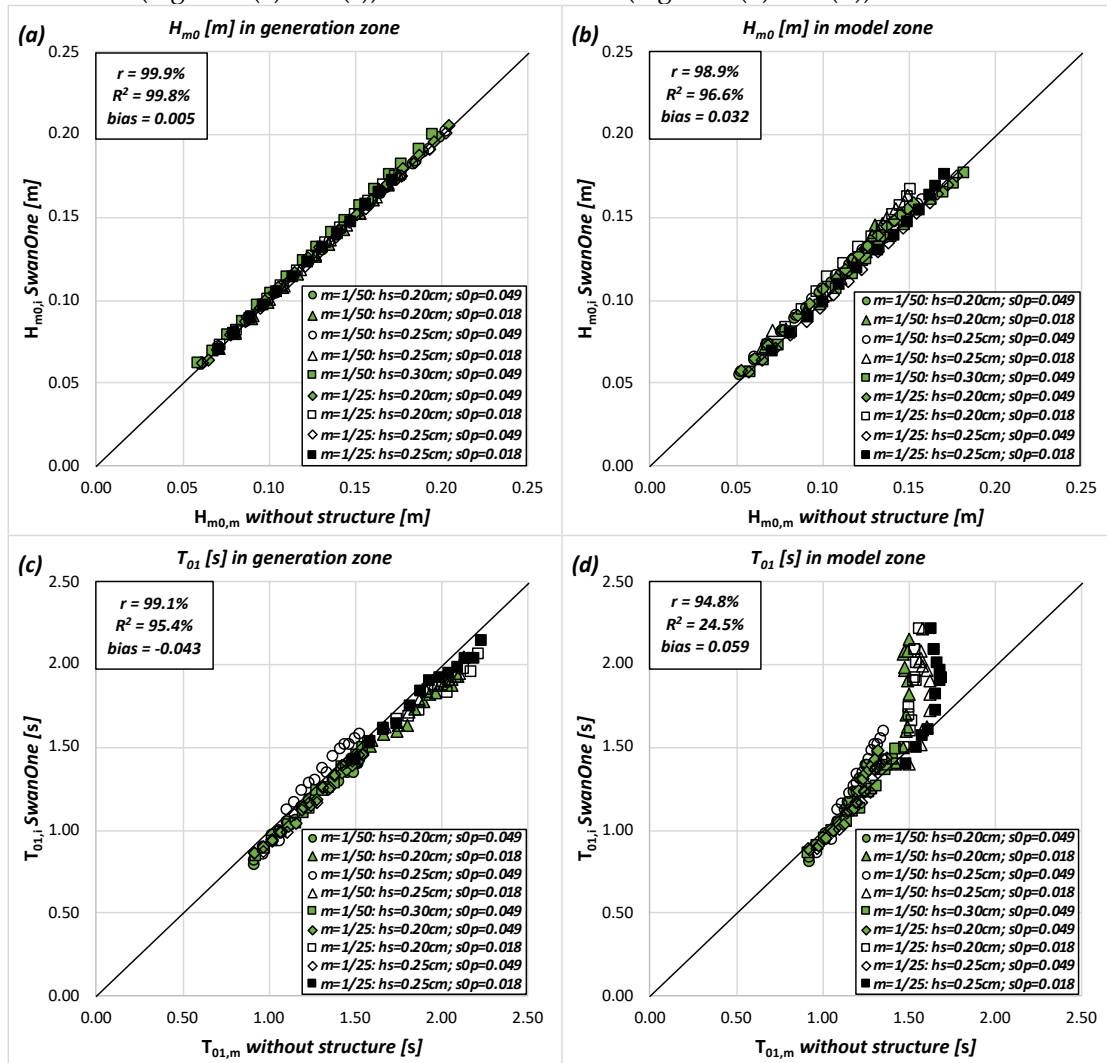
279 Figure 3. Experimental set up with the Cubipod®-1L armored breakwater model.
280

281 3.2. Wave analysis

282 Incident and reflected waves were separated in the wave generation zone using wave gauges S1
283 to S5 applying the LASA-V method (see Figueres and Medina, 2004). Although the LASA-V method
284 is applicable to nonstationary and nonlinear irregular waves, it is not valid for breaking waves.

285 In order to estimate the incident wave conditions in the model zone, where wave breaking takes
286 place, SwanOne software was used. This model assumes a Composite Weibull distribution to
287 describe the wave height distribution in shallow foreshores, as suggested by Battjes and Groenendijk

288 (2000). Following the methodology proposed by Herrera and Medina (2015), the incident wave height
 289 in the depth-induced wave breaking zone was estimated with the SwanOne model using the incident
 290 waves at the wave generation zone. SwanOne model fits a JONSWAP spectrum ($\gamma=3.3$) based on the
 291 given incident wave conditions in the wave generation zone and propagates such fitted wave
 292 conditions along a given bathymetry. Herrera and Medina (2015) validated this method comparing
 293 the numerical SwanOne simulations with the measurements in the wave flume conducted without
 294 any structure. The results of the validation in this study are given in Figure 4 for both the wave
 295 generation zone (Figure 4 (a) and (c)) and the model zone (Figure 4 (b) and (d)).



296
 297 Figure 4. Comparison between: (a) the incident wave height obtained with the SwanOne and the
 298 measured significant wave height without a structure in generation zone, (b) the incident wave height
 299 obtained with the SwanOne and the measured significant wave height without a structure in model
 300 zone, (c) the incident mean period obtained with the SwanOne and the measured mean period
 301 without a structure in generation zone and (d) the incident mean period obtained with the SwanOne
 302 and the measured mean period without a structure in model zone.

303
 304 To quantify the goodness of fit in this study, the correlation coefficient (r), the coefficient of
 305 determination (R^2) and the relative bias ($bias$) were calculated. $0 \leq r \leq 1$ assesses the correlation, $0 \leq R^2$
 306 ≤ 1 estimates the proportion of the variance explained by the model and $-1 \leq bias \leq 1$ provides a
 307 dimensionless measure of the bias. Thus, the higher the r , the higher the R^2 and the closer the $bias$ to
 308 0, the better.

$$r = \frac{\sum_{i=1}^{N_o} (o_i - \bar{o})(e_i - \bar{e})}{\sqrt{\sum_{i=1}^{N_o} (o_i - \bar{o})^2 \sum_{i=1}^{N_o} (e_i - \bar{e})^2}} \quad (24)$$

$$R^2 = 1 - \frac{\frac{1}{N_o} \sum_{i=1}^{N_o} (o_i - e_i)^2}{\frac{1}{N_o} \sum_{i=1}^{N_o} (o_i - \bar{o})^2} \quad (25)$$

$$bias = \frac{1}{N_o} \sum_{i=1}^{N_o} \frac{(e_i - o_i)}{|o_i|} \quad (26)$$

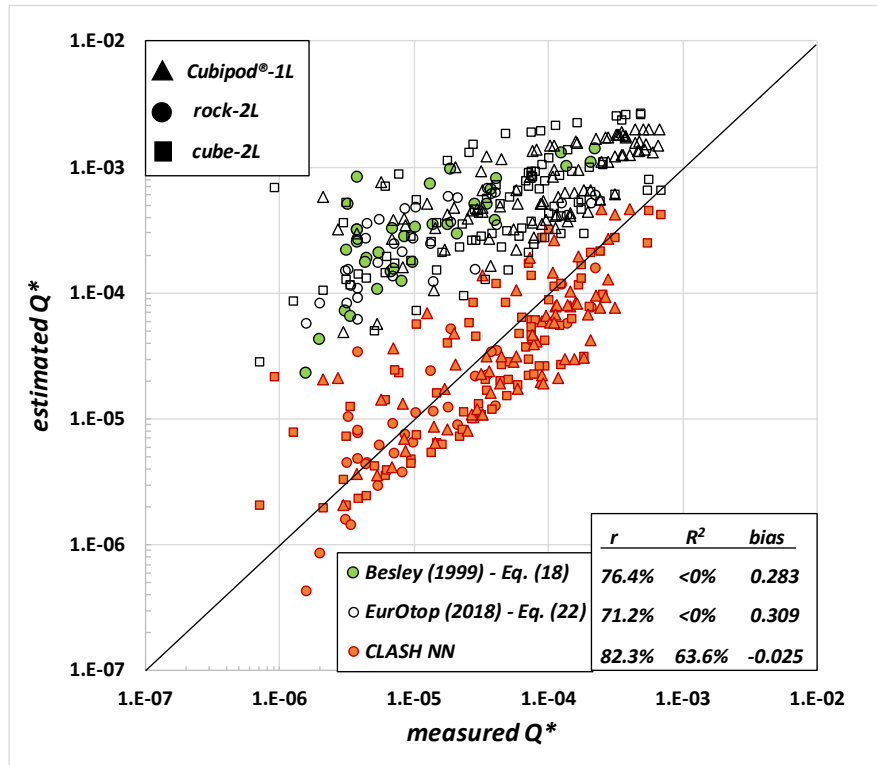
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310 where N_o is the number of observations, o_i and e_i are the observed and estimated values, and \bar{o} is the
 311 average observed value. Figure 4 shows that the agreement is very good for the fitted conditions in
 312 the wave generation zone ($R^2 > 95.4\%$). On the other hand, in the model zone good agreement is
 313 obtained for H_{m0} ($R^2 = 96.6\%$) whereas poor results are observed for T_{01} ($R^2 = 24.5\%$). Note that
 314 decreasing values of $bias$ were observed for H_{m0} in the model zone for increasing values of h_s : $bias =$
 315 0.057 for $h_s = 0.20\text{cm}$, $bias = 0.021$ for $h_s = 0.25\text{cm}$ and $bias = -0.018$ for $h_s = 0.30\text{cm}$. SwanOne clearly
 316 overestimates T_{01} measured at the model area.

317 During the design phase of a mound breakwater, the design wave conditions (H_{m0} and T_{01}) in the
 318 location where the structure will be built need to be estimated. Thus, in this study, both H_{m0} and T_{01}
 319 estimated by SwanOne were used.

320 4. Estimations of N_{ow} and V_{max} with methods given in the literature

321 In this section, the performance of the formulas to estimate N_{ow} and V_{max} presented in Section 2
 322 is analyzed using the experimental data described in Section 3. As mentioned in Section 2.2, the
 323 formulas given in the literature for estimating N_{ow} and V_{max} require knowing the mean individual
 324 wave overtopping volume, $\bar{V} = q T_{01} N_w / N_{ow}$. Therefore, q needs to be estimated using formulas in
 325 the literature when direct observations are not available. As shown in Table 2, estimators for q are
 326 suggested by the different authors of the methods to estimate N_{ow} and V_{max} on mound breakwaters.
 327 The goodness of fit of such estimators of q was assessed using the experimental data presented in
 328 Section 3. Figure 5 compares the observed and predicted Q^* using the estimators in Table 2.



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331

Figure 5. Comparison between measured and estimated $Q^* = q / (g H_{m0} T_{01})$.

332 The q estimator recommended by Besley (1999) - Eq. (18) was applied using $\gamma_f = 0.50$ for rock-2L.
 333 Since this author did not propose γ_f for Cubipod®-1L and cube-2L, it was not applied on those data.
 334 The q estimator proposed by EurOtop (2018) - Eq. (22) was used with $\gamma_f = 0.49, 0.40$ and 0.47 for
 335 Cubipod®-1L, rock-2L and cube-2L, respectively. CLASH NN was applied with γ_f recommended by
 336 Molines and Medina (2015b): $\gamma_f = 0.48, 0.49$ and 0.53 for Cubipod®-1L, rock-2L and cube-2L,
 337 respectively. q was estimated for 189 physical tests within the range of application of CLASH NN.

338 Note that the quantitative goodness-of-fit metrics in Figure 5 are given for $\ln(Q^*)$. As shown in Figure
 339 5, poor results ($R^2 < 0\%$) were obtained with the formulas given by Besley (1999) - Eq. (18) and
 340 EurOtop (2018) - Eq. (22). On the other hand, the predictor CLASH NN for Q^* suggested by Molines
 341 et al. (2019) and Nørgaard et al. (2014) provided good results with $R^2 = 63.6\%$.

342 4.1. Estimating N_{ow} with existing methods

343 In this section, the performance of the formulas to estimate N_{ow} presented in Section 2.2 is
 344 assessed. Figure 6 compares the observed N_{ow} with different estimators valid for mound breakwaters.
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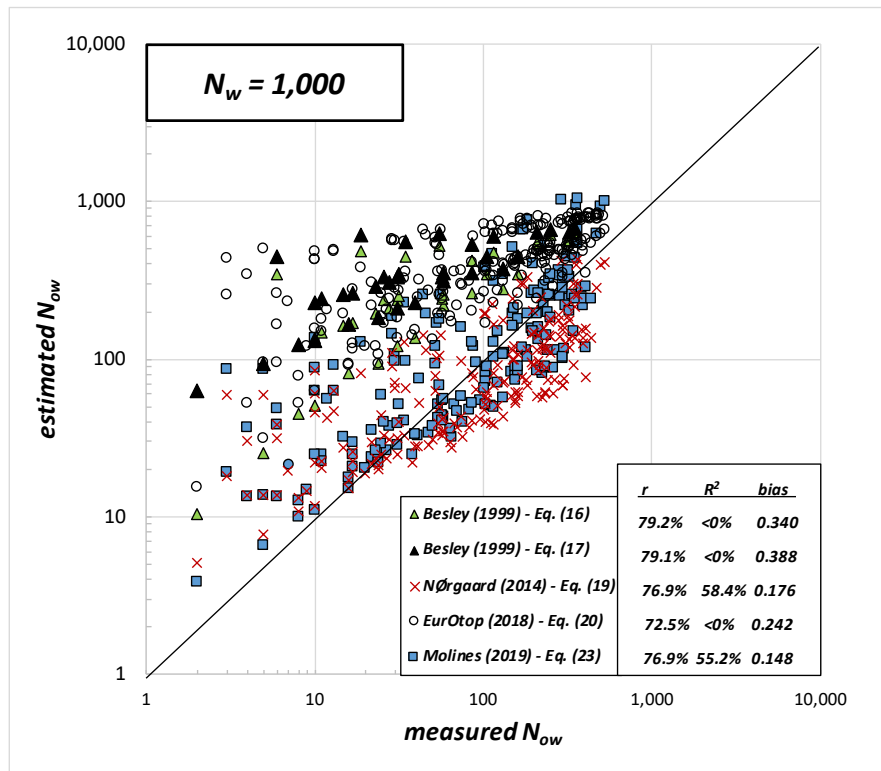


Figure 6. Comparison between measured and estimated N_{ow} with methods given in the literature.

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All formulas were applied no matter the application range. Note that estimated q suggested by the authors of these formulas was used in the N_{ow} predictors given by Besley (1999) - Eq. (17), Nørsgaard et al. (2014) - Eq. (19) and Molines et al. (2019) - Eq. (23). The N_{ow} estimator presented in Eq. (16) by Besley (1999) was applied using the roughness factor $\gamma_f = 0.50$ for rock-2L. Since this author did not recommend γ_f for Cubipod®-1L and cube-2L, it was not applied on those data. $K1 = 50.8$ was used with Eq. (16) for $\cot \alpha = 1.5$, obtained from the interpolation of the values given for $\cot \alpha = 2$ and $\cot \alpha = 1$. Eq. (20) was applied using γ_f given by EurOtop (2018): $\gamma_f = 0.49, 0.40$ and 0.47 for Cubipod®-1L, rock-2L and cube-2L, respectively.

The quantitative goodness-of-fit metrics are given for $\ln(N_{ow})$. Eqs. (19) and (23) proposed by Nørsgaard et al. (2014) and Molines et al. (2019), respectively, provided the best agreement with experimental data ($55.2\% < R^2 < 58.4\%$). Besley (1999) - Eqs. (16) and (17) and EurOtop (2018) - Eq. (20) overpredicted the values of N_{ow} and provided poor results ($R^2 < 0\%$). Note that all the compared methods from the literature overpredicted the values of $N_{ow} < 100$. Figure 6 shows that none of the existing N_{ow} estimators properly describe N_{ow} for the range of variables analyzed in this study. For this reason, a new N_{ow} estimator is developed in Section 5.

4.2. Estimating V_{max} with existing methods

In Section 2, several formulas to estimate the shape factor b were presented. Nevertheless, most of them are not valid for mound breakwaters. For this reason, only the formulas for mound breakwaters will be considered in the following comparison with the V_{max} measured in this study.

In this analysis, N_{ow} and q are estimated with the methods proposed by the authors (see Table 2). The estimators for N_{ow} and q have been previously assessed in this Section. The scale factor, A , was calculated using Eq. (3) for Nørsgaard et al. (2014) and EurOtop (2018), while Eq. (8) was applied for Molines et al. (2019). V_{max} was estimated by Eq. (15).

Figure 7 compares the measured and the estimated dimensional V_{max} (l/m) and the dimensionless $V_{max}^* = V_{max}/(g H_{m0} T_{01}^2)$ using the methods valid for mound breakwaters.

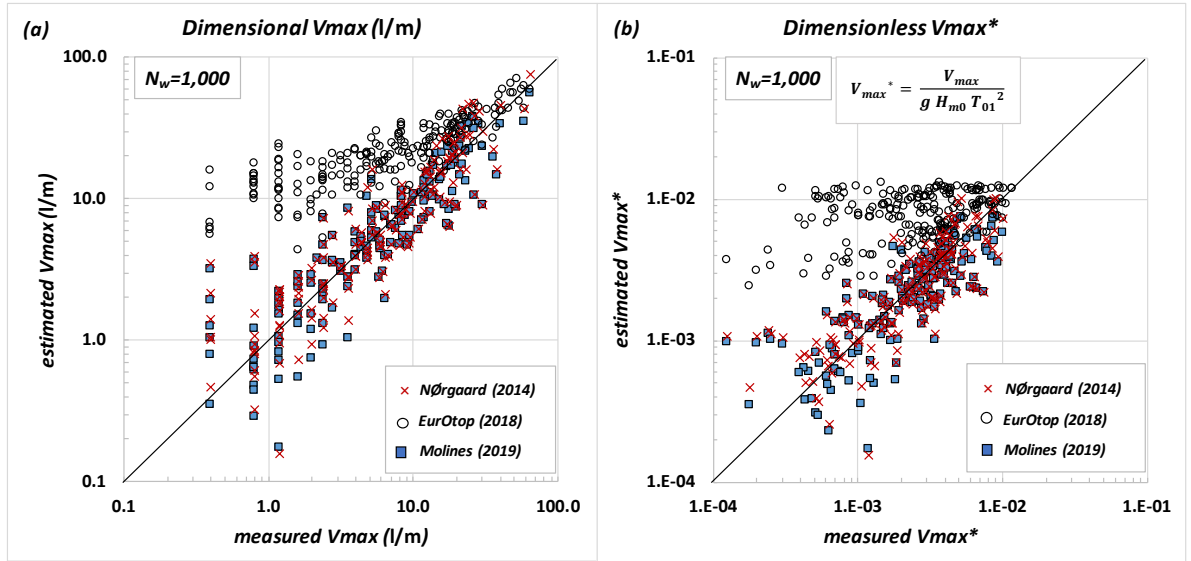


Figure 7. Comparison between measured and estimated (a) V_{max} (l/m) and (b) dimensionless V_{max}^* .

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V_{max}^* measured in this study agreed well with estimations given by Nørgaard et al. (2014) and Molines et al. (2019) with higher scatter for $V_{max} < 5$ l/m and $V_{max}^* < 2 \cdot 10^{-3}$. Note that Molines et al. (2019) was developed for mound breakwaters with crown wall in non-breaking wave conditions. Thus, depth-limited breaking may not have a significant effect on V_{max}^* . Table 4 presents the quantitative measurements of the goodness of fit as well as the number of variables and parameters of the methods shown in Figure 7. The quantitative goodness-of-fit metrics in Table 4 are given for $\ln(V_{max})$ and $\ln(V_{max}^*)$.

Author		V_{max} (l/m)	V_{max}^* (-)	# parameters	# variables
Nørgaard et al. (2014)	r	88.8%	79.0%	13	3
	R^2	78.4%	61.8%		
	$bias$	0.148	-0.002		
EurOtop (2018)	r	83.0%	38.2%	12	4
	R^2	< 0%	< 0%		
	$bias$	2.222	0.173		
Molines et al. (2019)	r	89.1%	79.8%	7	1
	R^2	78.8%	63.0%		
	$bias$	-0.023	-0.017		

385

Table 4. Goodness-of-fit metrics for the methods in the literature to estimate V_{max} .

386

5. Estimating of the number of overtopping events, N_{ow}

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5.1. A new formula to estimate N_{ow}

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As shown in Section 2, most of the existing estimators of $P_{ow} = N_{ow}/N_w$ are a function of a power of Q^* , as Eqs. (17), (19) and (23). Methods by Nørgaard et al. (2014) - Eq. (19) and Molines et al. (2019) - Eq. (23) provide good results within their range of application, but they do not properly estimate P_{ow} for very low or very large Q^* . When Q^* is very small, P_{ow} should tend to 0, and when Q^* is very large, P_{ow} should tend to 1. Therefore, an exponential model may be better than a power law of P_{ow} , because it is good for very low and very high values of Q^* .

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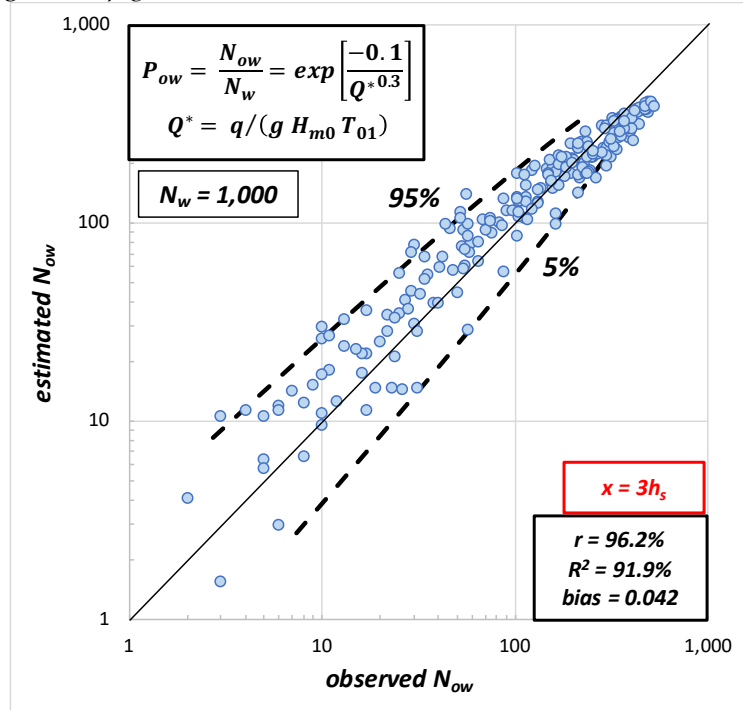
In this study, a clear correlation was found between P_{ow} and Q^* , as previously observed by Besley (1999), Nørgaard et al. (2014) and Molines et al. (2019). Therefore, an exponential function of Q^* was proposed and G_1 and G_2 in $P_{ow} = \exp(-G_1/Q^{*G_2})$ were calibrated based on the 219 tests maximizing R^2 of $\ln N_{ow}$. The three armor layers tested in this study are not distinguished in the analysis (Bruce et al., 2009).

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399 As exposed in Section 1, in depth-limited breaking wave conditions, the optimum point where
 400 wave characteristics are estimated is relevant for design and needs to be determined (Mares-Nasarre
 401 et al., 2020). Thus, G_1 and G_2 were calibrated considering wave characteristics at several distances
 402 from the structure toe in $Q^* = q/(g H_{m0} T_{01})$. H_{m0} and T_{01} were estimated with SwanOne at distances $x =$
 403 $0, h_s, 2h_s, 3h_s, 4h_s, 5h_s$ and $6h_s$ from the model toe. No significant differences were observed; $G_1 = -0.1$
 404 and $G_2 = 0.3$ were obtained for H_{m0} and T_{01} estimated between the model toe and at $6h_s$ from the model
 405 toe. In this study, wave characteristics estimated at a distance of $3h_s$ from the model toe are used,
 406 following Herrera et al. (2017) and Mares-Nasarre et al. (2020) recommendations. Note that this
 407 distance is approximately the same as $x = 5H_{m0}$, suggested by Melby (1999) and Goda (1985) to
 408 determine wave characteristics in breaking wave conditions. Figure 8 compares the experimental data
 409 and Eq. (27) with $R^2 = 91.9\%$, as well as the 90% error band.

$$P_{ow} = \frac{N_{ow}}{N_w} = \exp\left(\frac{-0.1}{Q^{*0.3}}\right) \quad (27)$$

410 where $P_{ow} = N_{ow}/N_w$ is the proportion of overtopping waves and Q^* is the dimensionless wave
 411 overtopping discharge, $Q^* = q/(g H_{m0} T_{01})$.



412 Figure 8. Comparison between observed number of overtopping events, N_{ow} , and estimated N_{ow}
 413 given by Eq. (27) using wave characteristics estimated at a distance of $3h_s$ from the model.
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 415

416 In this study, the methodology given in Herrera and Medina (2015) and applied in works such
 417 as Mares-Nasarre et al. (2019) is used to estimate the 90% error band. Thus, a Gaussian distribution
 418 of the error (ε) is assumed, with 0 mean and the variance given by

$$\sigma^2(\varepsilon) = 0.55 - 0.09 \ln N_{ow} \quad (28)$$

419 The 95% and 5% percentiles for the N_{ow} predicted by Eq. (27) can be calculated using Eq. (29).

$$\ln N_{ow}|_{5\%}^{95\%} = \ln N_{ow} \pm 1.64 \sqrt{0.55 - 0.09 \ln N_{ow}} \quad (29)$$

420 The range of application of Eq. (26) is $0.002 \leq P_{ow} \leq 0.53$ and $7.2 \cdot 10^{-7} \leq Q^* \leq 6.9 \cdot 10^{-4}$. Eq. (26) properly
 421 extrapolates the prediction of $P_{ow}=0$ when $Q^*=0$ and $P_{ow}=1$ when $Q^* \rightarrow \infty$.

422 5.2. Influence of bottom slope on N_{ow}

423 In Section 3, the experimental setup using two bottom slope configurations with $m = 1/50$ and m
 424 $= 1/25$ was described. No significant difference between bottom slopes $m = 1/50$ and $m = 1/25$ was
 425 observed in the scatter plot. A statistical analysis was performed to determine if N_{ow} are equally
 426 distributed for different bottom slopes. Since the data were not Gaussian distributed, a nonparametric
 427 Mann-Whitney test (Mann and Whitney, 1947) was applied. In this test, the null hypothesis (H_0)
 428 corresponded to N_{ow} not being affected by the bottom slope. Based on 103 tests with a bottom slope
 429 $m = 1/50$ and 116 tests with a bottom slope $m = 1/25$, H_0 was not rejected using a significance level $\alpha =$
 430 0.10. Thus, in this study the bottom slope does not show any significant influence on N_{ow} .

431 6. Estimating of the maximum individual wave overtopping volume, V_{max}

432 6.1. A new method to estimate V_{max} using the 2-parameter Weibull distribution

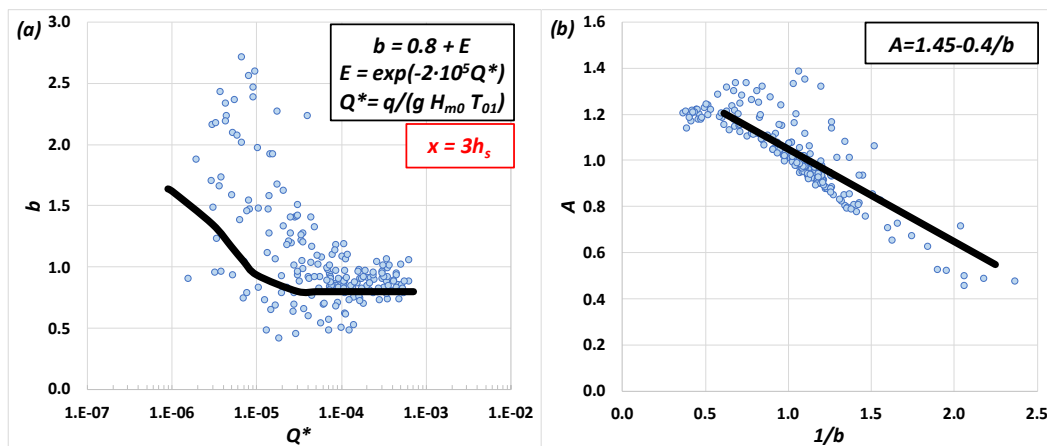
433 The maximum individual wave overtopping volume, V_{max} , is estimated using Eq. (15). Then, V_{max}
 434 depends on the number of overtopping events, N_{ow} , shape and scale factors of the Weibull
 435 distribution, A and b , and $\bar{V} = V_{total}/N_{ow}$. As previously mentioned, both A and b obtained for each test
 436 are fitted using a quadratic utility function applied to the whole individual wave overtopping volume
 437 dataset. Tests with very low values of N_{ow} ($N_{ow} < 5$) were not used in this analysis to prevent
 438 inconsistencies caused by a very low number of observations. Eqs. (30) and (31) were proposed to
 439 characterize A and b

$$b = K_{B1} + \exp(K_{B2} Q^*) \quad (30)$$

$$A = 1.45 - 0.4/b \quad (31)$$

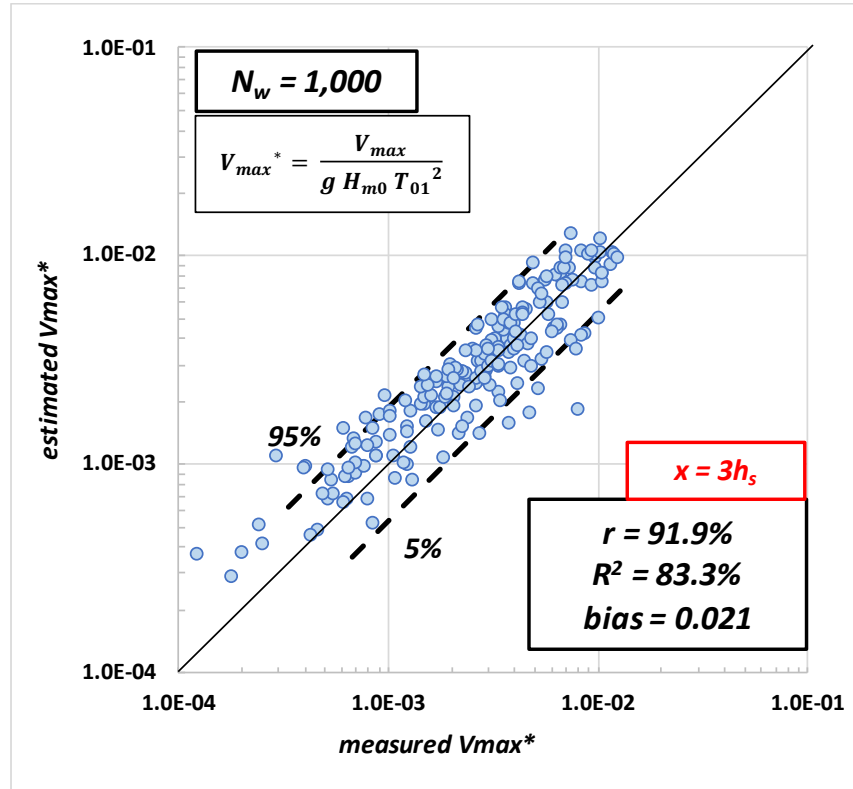
440 Similar to Section 5.1, K_{B1} and K_{B2} were calibrated considering wave characteristics at several
 441 distances from the structure toe in $Q^* = q/(g H_{m0} T_{01})$. H_{m0} and T_{01} were determined with SwanOne at
 442 distances $x = 0, h_s, 2h_s, 3h_s, 4h_s, 5h_s$ and $6h_s$ from the structure toe. The goodness-of-fit of $V_{max}^* = V_{max}/(g$
 443 $H_{m0} T_{01}^2)$ was assessed for every couple of coefficients calculated using Eqs. (15), (30) and (31) with
 444 the measured N_{ow} and \bar{V} . Best fit was obtained between $x = 2h_s$ and $x = 6h_s$, $K_{B1} = 0.8$ and $K_{B2} = -2 \cdot 10^5$
 445 were obtained with $r = 92\%$ and $R^2 = 83\%$. Wave characteristics were decided to be estimated at $x =$
 446 $3h_s$.

447 Figure 9a illustrates the relationship between Q^* and b and the least-squares fitting given by Eq.
 448 (30). Figure 9b relates A and $1/b$ and the least-squares fitting given by Eq. (31). The range of
 449 application of Eqs. (30) and (31) is $0.005 \leq P_{ow} \leq 0.53$ and $1.3 \cdot 10^{-6} \leq Q^* \leq 6.9 \cdot 10^{-4}$.
 450



451
 452 Figure 9. Relationship between explanatory variables and the least-squares fitting of Weibull
 453 distribution factors: (a) Weibull's shape factor, b , in Eq. (30) and (b) Weibull's scale factor, A , in Eq.
 454 (31).
 455

456 Figure 10 illustrates the performance of Eq. (15) to estimate $V_{max}^* = V_{max}/(g H_{m0} T_{01}^2)$ when using A
 457 and b obtained from Eqs. (30) and (31) and measured N_{ow} and \bar{V} . The agreement was good; $R^2 = 83.3\%$.



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460

Figure 10. Comparison between the measured and estimated dimensionless V_{max}^* by the 2-parameter Weibull distribution with shape and scale factors given by Eqs. (30) and (31).

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462

The variance of the error (ε) of $\ln V_{max}^*$ is $\sigma^2(\varepsilon) = 0.15$. Thus, the 95% and 5% percentiles for the predicted V_{max}^* by Eq. (15) are given by

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$$\ln V_{max}^* \Big|_{5\%}^{95\%} = \ln V_{max}^* \pm 0.63 \quad (32)$$

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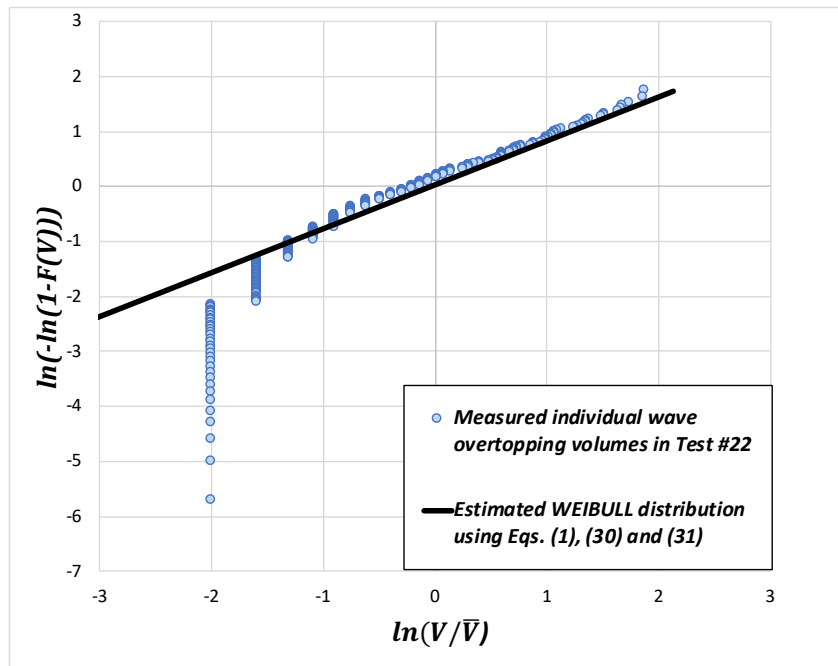
Figure 11 shows an example of the fit of the 2-parameter Weibull distribution to the measured individual wave overtopping volumes for a randomly-selected test (#22). Figure 11 is presented in a Weibull plot:

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$$\ln(-\ln(1 - F(V))) = b(\ln(V/\bar{V}) - \ln(A)) \quad (33)$$

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469 Figure 11. Comparison of measured and estimated individual wave overtopping volumes for
 470 Test #22 using Eqs. (1), (30) and (31).
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473 As explained in Section 2, Molines et al. (2019) proposed a 2-parameter Exponential distribution
 474 given by Eq. (9) to describe $F(V)$. In this study, the 2-parameter Exponential distribution was also
 475 fitted with good results.

476 6.2. Influence of bottom slope on the 2-parameter Weibull distribution

477 As presented in section 5.2, the influence of bottom slope on b and A estimated by Eqs. (30) and
 478 (31) is studied here. No relevant differences between the bottom slopes $m = 1/50$ and $m = 1/25$ were
 479 found. The Mann-Whitney test was applied for each parameter, as described in section 5.2. The null
 480 hypothesis (H_0) corresponded to b and A not being influenced by the bottom slope. Based on 97 tests
 481 with a bottom slope $m = 1/50$ and 106 tests with a bottom slope $m = 1/25$, H_0 was not rejected using a
 482 significance level $\alpha = 0.10$. Thus, in this case bottom slope does not show any significantly influence
 483 on V_{max} .

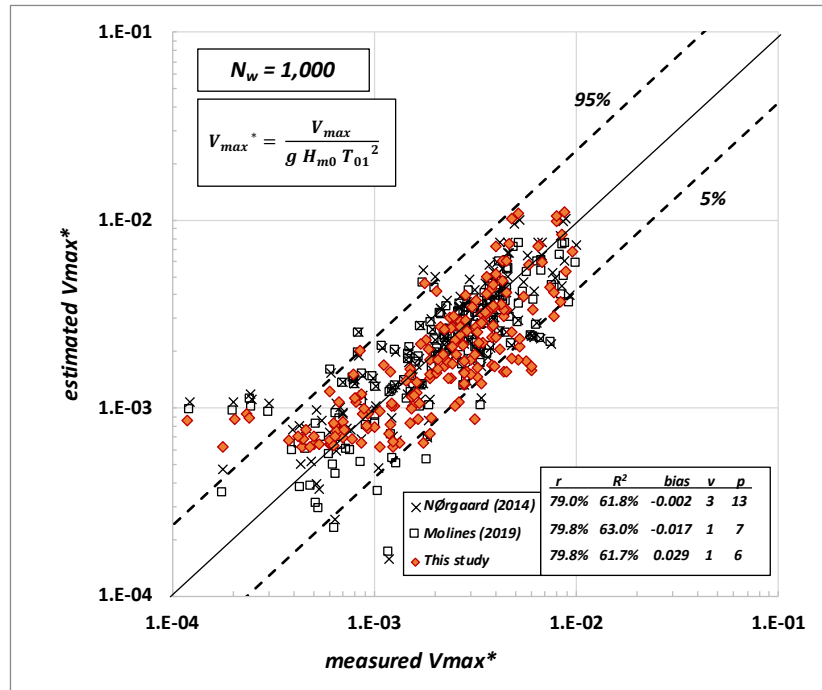
484 7. Estimation of V_{max} for mound breakwater designs

485 During the design phase of a mound breakwater, the design wave conditions (H_{m0} and T_{01}) and
 486 the structural geometry are given. The formulas given in the literature require knowing the mean
 487 individual wave overtopping volume, $\bar{V} = q T_{01} N_w / N_{ow}$, and the number of overtopping events, N_{ow} .
 488 Therefore, q and N_{ow} need to be estimated using formulas in the literature when direct observations
 489 (e.g. sections 5 and 6) are not available. Methods suggested in the literature to estimate q were
 490 assessed in Section 4. The best fit was given by CLASH NN with $R^2 = 63.6\%$. Thus, Q^* estimated with
 491 CLASH NN was used in the following. When using Eq. (27) developed in this study to calculate N_{ow}
 492 estimating q with CLASH NN, $r = 77.1\%$, $R^2 = 58.4\%$ and $bias = 0.026$. The agreement was worse than
 493 $R^2 = 91.9\%$ (see Figure 8) obtained when q is measured and not estimated. Note that the fitting of the
 494 new formula is equal to the one obtained with the method proposed by Nørgaard et al. (2014) (see
 495 Figure 6). However, the method proposed by Nørgaard et al. (2014) to estimate N_{ow} requires 3
 496 variables and 10 parameters, while the method developed in this study uses 1 variable and 2
 497 parameters.

498 Figure 12 illustrates the goodness of fit of Eq. (15) to estimate V_{max} based on the estimations of
 499 N_{ow} and \bar{V} when q is estimated using CLASH NN. Eqs. (30) and (31) were applied to estimate the

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Weibull parameters (A, b). $R^2 = 61.7\%$ was obtained for $V_{max}^* = V_{max}/(g H_{m0} T_{01}^2)$. Figure 12 also presents the estimations given by Nørgaard et al. (2014) and Molines et al. (2019) as well as the 90% error band.



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Figure 12. Comparison of measured dimensionless maximum individual wave overtopping volume, V_{max}^* , and the estimated V_{max}^* for the Weibull distribution model using Eqs. (15), (30) and (31) and the methods for mound breakwaters in the literature.

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As shown in Figure 12, the goodness of fit of the three compared methods is similar. However, the method proposed in the present study (1 variable (v) and 6 parameters (p)) is much simpler than the method proposed by Nørgaard et al. (2014) ($v = 3$ and $p = 13$) and simpler than the method proposed by Molines et al. (2019) ($v = 1$ and $p = 7$). Using the method proposed in this study, the ratio between estimated and measured V_{max}^* for design purposes falls within a factor of 2.0 (90% error band).

513 8. Conclusions

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Crest elevation of mound breakwaters is usually designed to limit the mean wave overtopping rate (q) or the maximum individual wave overtopping volume (V_{max}). Furthermore, rising sea levels caused by climate change and mounting social pressure to minimize the visual impact of coastal structures mean lower crest freeboards and increased overtopping hazards. Thus, coastal structure designs with relevant overtopping rates attacked by waves breaking on the sea bottom become relevant. Few studies have been conducted in depth-limited breaking wave conditions. This research is focused on mound breakwaters with significant overtopping rates ($0.002 \leq P_{ow} \leq 0.53$ and $7.2 \cdot 10^{-7} \leq Q^* = q / (g H_{m0} T_{01}) \leq 6.9 \cdot 10^{-4}$) and intermediate crest elevations ($0.33 \leq R_c / H_{m0} \leq 2.83$) with armor slope $\cot \alpha = 1.5$ in depth-limited breaking wave conditions ($0.2 \leq H_{m0} / h_s \leq 0.9$), considering two bottom slopes ($m = 1/50$ and $m = 1/25$).

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In this study, 105 physical tests with a bottom slope $m = 1/50$ and 114 tests with $m = 1/25$ were conducted at the LPC-UPV wave flume. Individual wave overtopping volumes were analyzed using Molines et al. (2019) methodology, based on a continuous record of accumulated overtopping volume. In order to estimate the incident wave conditions in the model zone, where breaking occurs, SwanOne model was used. The performance of SwanOne when estimating wave characteristics under depth-limited breaking wave conditions was assessed using tests without structure. The agreement was good for H_{m0} ($R^2=96.6\%$), but some bias was observed for the shallowest water depth ($bias = 0.057$). The agreement was poor for T_{01} ($R^2=23.3\%$); SwanOne overestimated the measured T_{01} in the model zone.

533 Estimators for P_{ow} given in the literature were assessed using the experimental data; it was
534 observed that most existing formulas overpredict P_{ow} for $P_{ow} < 10\%$. Most of the existing P_{ow} estimators
535 are a function of a power of Q^* , so they cannot fit the boundary limits ($P_{ow} \rightarrow 0$ when $Q^* \rightarrow 0$ and $P_{ow} \rightarrow 1$
536 when $Q^* \rightarrow \infty$). Hence, a new exponential estimator is given for P_{ow} valid for depth-limited breaking
537 wave conditions in Eq. (27) ($R^2 = 91.9\%$).

538 The quadratic utility function proposed by Molines et al. (2019) was applied in this study to all
539 the data to fit the 2-parameter Weibull distribution for individual wave overtopping volumes, $F(V)$.
540 Estimators were taken from Nørsgaard et al. (2014), EurOtop (2018) and Molines et al. (2019) for the
541 scale (A) and shape (b) factors of the Weibull distribution to compare with the measured data,
542 obtaining $0\% \leq R^2 \leq 63.0\%$ for the dimensionless maximum individual wave overtopping volume,
543 $V_{max}^* = V_{max}/(g H_{m0} T_{01}^2)$. Best results using methods given in the literature were obtained for Molines
544 et al. (2019) whose method was developed for mound breakwaters with crown wall in non-breaking
545 wave conditions. Thus, the influence of the depth-induced wave breaking or the presence of the crest
546 wall may not be significant.

547 New estimators for the factors A and b of the Weibull distribution were fitted using the
548 experimental data. The new Weibull (Eqs. (30) and (31)) distribution provide estimations of V_{max}^* with
549 $R^2 = 83.3\%$ and a number of variables and parameters lower than those of the methods in the
550 literature. In this study, no significant influence of bottom slope ($1/50 \leq m \leq 1/25$) was found on N_{ow}
551 and V_{max} .

552 During the design phase of a mound breakwater, the mean individual wave overtopping
553 volume ($\bar{V} = q T_{01} N_w/N_{ow}$) is required to estimate V_{max} . But q and N_{ow} are unknown, and they have to
554 be estimated using methods in the literature when direct observations are not available. Here,
555 CLASH Neural Network (CLASH NN) was used to estimate q with $R^2 = 63.6\%$. Using q estimated by
556 the CLASH NN and the new N_{ow} estimator given in Eq. (27), V_{max}^* was estimated with the 2-parameter
557 Weibull distribution proposed in this study. The prediction error of V_{max}^* dropped from $R^2 = 83.3\%$
558 when q and N_{ow} were measured in the laboratory to $R^2 = 61.7\%$ when q was estimated with CLASH
559 NN. The ratio between estimated and measured V_{max}^* falls within a factor of 2.0 (90% error band) for
560 design purposes.

561 The estimators and conclusions derived here are valid within the experimental ranges of this
562 study. Therefore, it is encouraged to check their validity out of these experimental ranges, paying
563 special attention to the significance of the depth-limited breakage of waves and the presence of a
564 crown wall.

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