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Additional Information

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3 SCALING SPRAY PENETRATION AT SUPERSONIC CONDITIONS  
4 THROUGH SHOCKWAVE ANALYSIS

5

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20 **Keywords**

21 Diesel spray, shockwave, penetration, visualization

## 22 **Abstract**

23 In the current paper, an investigation of the supersonic flow effect on shockwave  
24 generation and diesel spray penetration scaling has been performed. For this purpose,  
25 spray visualization tests have been carried out in a constant-pressure chamber at room  
26 temperature using shadowgraphy technique. Two working gases have been used: nitrogen,  
27 with similar thermodynamic characteristics to the engine environment, and sulfur  
28 hexafluoride, aimed at producing supersonic conditions at moderate injection pressure  
29 values. A total of 60 operating points, including different nozzle geometries, injection  
30 pressures and chamber densities have been studied. From the visualization study, two  
31 different kinds of shockwaves have been detected: normal or frontal, for moderate spray  
32 tip Mach (between 1 and 1.5); and oblique, when the Mach is higher than 1.5. The  
33 penetration results show that, for the same injection conditions in terms of injection  
34 pressure and chamber density, the spray propagation is equal for SF<sub>6</sub> and N<sub>2</sub> when the  
35 spray is on subsonic conditions, while penetration is higher for SF<sub>6</sub> when supersonic  
36 velocity is reached. This behavior has been related to the density gradient appearing  
37 across the shockwave. A new methodology to extrapolate supersonic penetration from  
38 the well-known subsonic penetration law has been proposed, showing good agreement  
39 with the experimental results.

## 40 **1. Introduction**

41 Spray penetration is one of the key physical parameters involved in diesel combustion  
42 process. Traditionally, spray penetration has been linked to the nozzle outlet momentum  
43 (which depends on the injection pressure) and the density ratio between the fuel and the  
44 gas inside the combustion chamber, with an exponent around -0.25 [1–3]. Furthermore,  
45 several authors have noticed that the dependency of this spray penetration shows two

46 different trends: initially, penetration increases linearly with time, while later on the spray  
47 tip velocity reduces, resulting in a square root scaling of penetration versus time [4,5].  
48 This transition time has been correlated with the transient behavior of the momentum at  
49 the outlet of the nozzle due to the injector needle dynamics [6].

50 In the recent years, fuel injection systems design has evolved towards higher injection  
51 pressures. On the one hand, higher turbulence levels at the nozzle outlet help to improve  
52 spray atomization and air entrainment [7–10]. On the other hand, faster spray penetration  
53 can help to achieve better air utilization and improve combustion speed at high loads [11].  
54 Linked to this injection pressure increase, several researchers have noticed that the spray  
55 tip can reach supersonic conditions. Kostas et al. [12] realized that supersonic conditions  
56 during the initial spray stages could affect spray penetration due to compressibility effects.  
57 Hillamo et al. [13] used backlighting visualization to detect shockwaves formation  
58 attached to the spray edge. Jia et al. [14] made experiments at ultra-high pressure (400  
59 MPa) and showed that two kinds of shockwaves can appear: spherical (or bow) and  
60 oblique. In a subsequent study, Jia et al. [15] compared the spray penetration at this  
61 supersonic conditions with the well-established spray penetration correlations,  
62 concluding that spray propagation was affected by the supersonic regime. Wang et al.  
63 [16] made a similar investigation, where also spray cone angle and atomization (in terms  
64 of Sauter Mean Diameter) began to vary when supersonic conditions were reached.  
65 Finally, Song et al. [17] studied the shockwave angles and the implications on air  
66 entrainment. Nevertheless, in any of the aforementioned investigations an analysis aimed  
67 at providing an analytical expression to estimate the supersonic spray penetration law has  
68 been performed.

69 In the current paper, a further investigation on supersonic spray propagation  
70 characteristics is made. For this purpose, fuel is injected on a pressurized sulfur

71 hexafluoride ( $\text{SF}_6$ ) environment, aimed at reaching very high local Mach numbers with a  
72 standard fuel injection system hardware. Spray and shockwave visualization is made  
73 using shadowgraphy technique on a constant pressure vessel. The results from  
74 visualization on  $\text{SF}_6$  are compared with results made on nitrogen atmosphere, where the  
75 chamber pressure has been adapted in order to achieve the same gas density but on  
76 subsonic regime. Later on, a methodology to estimate the supersonic spray penetration  
77 from the subsonic conditions data is proposed and validated, which is the main  
78 contribution from the current work.

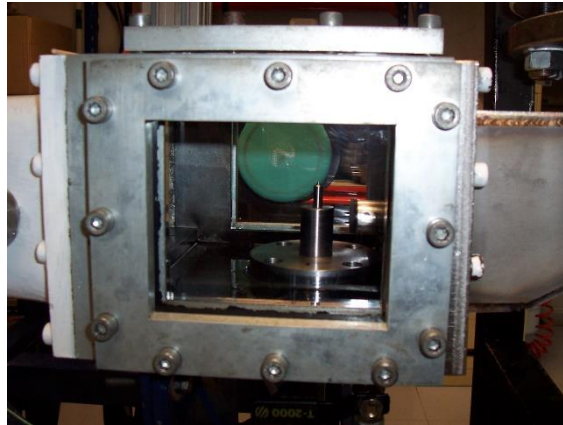
79 As far as the paper structure is concerned, five sections are developed. Section 2  
80 summarizes the main aspects of the experimental setup and used methodology. Section 3  
81 analyzes the shockwave visualization results. The methodology for spray penetration  
82 evaluation at supersonic conditions is proposed and evaluated in Section 4. Finally, the  
83 main conclusions extracted from the study are drawn in Section 5.

## 84 **2. Experimental setup**

85 Liquid spray penetration is characterized on a constant pressure chamber at room  
86 temperature conditions ( $25^\circ\text{C}$ ) [18]. A picture of this spray visualization chamber can be  
87 seen in Figure 1. The vessel has three rectangular optical accesses allowing to perform  
88 different optical diagnostics for spray characterization purposes. Initially, the chamber is  
89 filled with nitrogen, reproducing engine-like density conditions in a non-reactive  
90 environment. In particular, four levels of chamber pressure (1, 2, 3 and 4 MPa) are tested,  
91 corresponding to density values of 11.3, 22.6, 33.9 and  $45.2 \text{ kg/m}^3$ . Later on, the same  
92 density conditions are reproduced using sulfur hexafluoride ( $\text{SF}_6$ ). The main advantage  
93 of this gas consists of its higher molecular weight ( $146.06 \text{ g/mol}$  vs.  $28.014 \text{ g/mol}$  for  
94 nitrogen) and lower speed of sound ( $136.5 \text{ m/s}$  vs.  $351.8 \text{ m/s}$  for nitrogen at room

95 conditions). Thus, allowing local Mach in the spray tip is achieved at moderate injection  
96 pressures. To achieve the same densities, chamber pressures of 0.19, 0.38, 0.57 and 0.76  
97 MPa are selected for SF<sub>6</sub>.

98



99

100 *Figure 1. Detail of spray visualization test rig.*

101 A Bosch common-rail fuel injection system capable of reaching up to 160 MPa pressure  
102 is used for the experiments. The solenoid injector is equipped with five different single-  
103 hole axi-symmetric nozzles, with outlet diameters ranging from 0.119 to 0.206 mm. For  
104 each nozzle and density condition, three levels of injection pressure are evaluated (50, 80  
105 and 130 MPa) in order to reach both subsonic and supersonic conditions for the spray  
106 penetration evaluation. Thus, a total of 60 tests are performed for each filling gas. All  
107 tests have been performed using a European standard winter diesel fuel as working fluid.  
108 The main physical properties of this fuel, including density and viscosity, are reported in  
109 [19].

110 It has to be noted that the reduction in the discharge pressure achieved for the SF<sub>6</sub> tests  
111 could induce severe changes in the spray propagation characteristics if cavitation was  
112 induced inside the nozzles. However, for the set of nozzles used in the study (which are

113 convergent) a previous hydraulic characterization shows that no cavitation appears for  
114 the working pressures used in the study.

115 Regarding the optical setup, single-pass shadowgraphy technique is used to visualize both  
116 the spray and the shockwaves evolution. This technique is based on illuminating from  
117 one side of the spray with a Xe-arc lamp, and capturing the light with a camera on the  
118 opposite side. When there is a density change across the optical path (either produced by  
119 the spray droplets or by the shockwave), light beams deviate due to the different refractive  
120 index, producing a shadow on the camera sensor [20,21]. For this particular arrangement,  
121 a Photron high-speed camera set up at 40000 frames per second is employed. Each  
122 conditions is repeated a total of 8 times in order to evaluate the test repeatability. An  
123 example of the images obtained on the SF<sub>6</sub> tests, including the spray and the related  
124 shockwave, is seen in Figure 2. This images corresponds to an injection pressure of 80  
125 MPa, a gas density of 22.6 kg/m<sup>3</sup> and an elapsed time from the start of injection of 0.130  
126 milliseconds.



127

128 *Figure 2. Example of shadowgraphy visualization image for test in sulfur hexafluoride environment.*

129

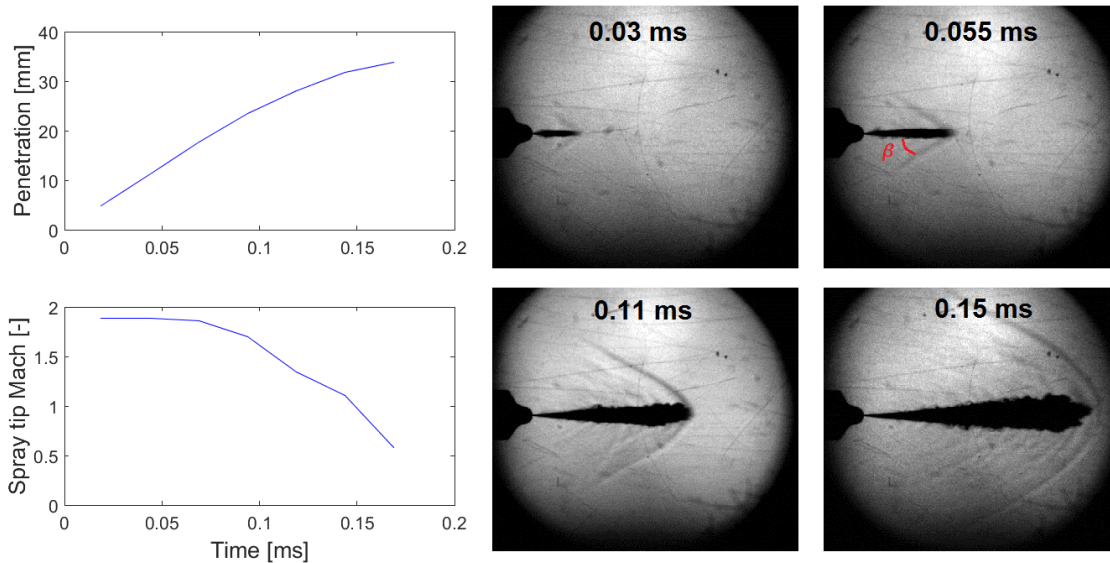
130 Later on, the shadowgraphy images are post-processed using a specific tool to determine  
131 the spray penetration. This penetration is defined as the maximum distance between the  
132 nozzle tip and the spray contour, previously detected using a threshold of 0.5% of the  
133 intensity dynamic range, according to the procedure described by Payri et al. [22].

### 134 **3. Spray visualization results**

135 Figure 3 shows a summary of the detected shockwave evolution in SF<sub>6</sub> from an injection  
136 pressure of 130 MPa and a chamber density of 22.6 kg/m<sup>3</sup>. In the left hand side of the  
137 image, the computed values of the spray tip penetration and spray tip Mach are plotted as  
138 a function of the time elapsed after the start of injection. The spray tip Mach is calculated  
139 from the time derivative of the spray penetration and the speed of sound in SF<sub>6</sub> at the  
140 corresponding thermodynamic conditions, which is estimated at 136.5 m/s.

141 In the right hand side of the figure, some images from specific moments along the spray  
142 evolution are depicted. Initially (up to approximately 0.1 ms, corresponding to a Mach  
143 number around 1.5), it is possible to observe the formation of a triangular (also called  
144 oblique) shockwave [23] attached to the spray tip, characterized by a certain shockwave  
145 semiangle  $\beta$ . This is the situation for the first two depicted images. From that moment  
146 onward, the shockwave transitions to a frontal or normal shockwave [23]. This can be  
147 clearly seen in the image at 0.11 ms, where the bow-shape tip of the wave appears.  
148 Considering all the tests, the transition between both kinds of shockwaves occurs in  
149 between one and two time steps of the visualization, corresponding to approximately  
150 0.025 to 0.05 ms. Finally, when the spray tip reaches subsonic conditions, the normal  
151 shockwave detaches from the spray tip and continues its evolution at the corresponding  
152 speed of sound.





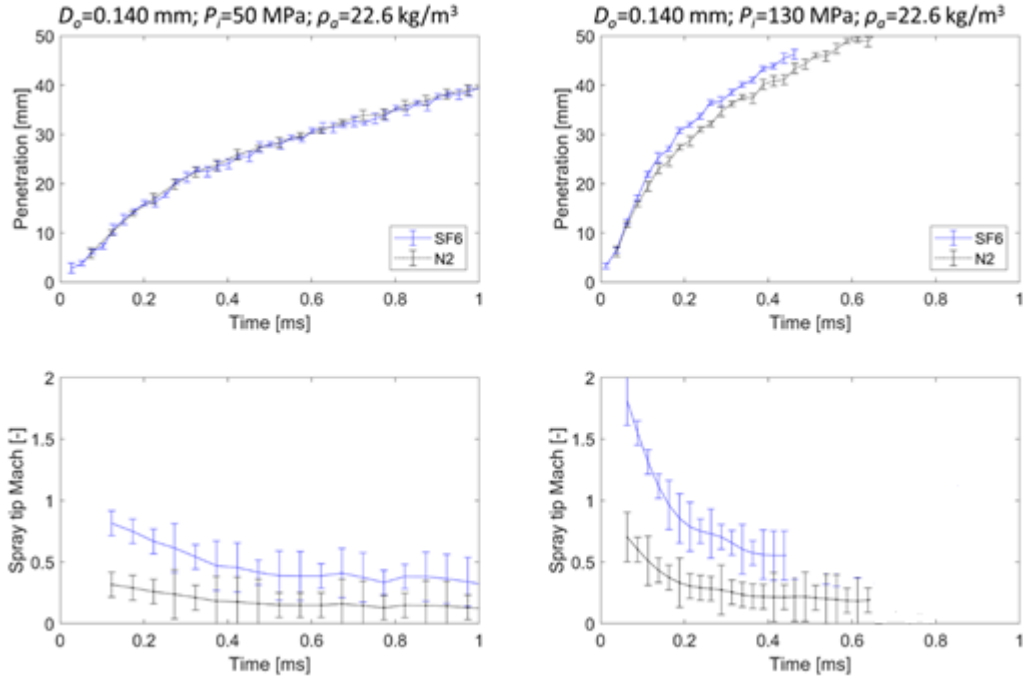
153

154 *Figure 3. Spray tip penetration, Mach and shockwave analysis for  $P_i=130$  MPa and density  $22.6$  kg/m<sup>3</sup>*  
 155 *from SF<sub>6</sub> tests.*

156 The aforementioned scenario repeats for the rest of the tests performed. Typically, when  
 157 the spray tip Mach is above 1.5, an oblique shockwave can be detected, while normal or  
 158 bow-shape shockwave is induced for Mach numbers between 1 and 1.5. Similar transition  
 159 from triangular to normal shockwaves was identified by Song et al. [17] when passing  
 160 from high to low spray velocity, but without a clear identification of the Mach threshold  
 161 for this transition.

162 Figure 4 shows the spray visualization results for the nozzle outlet diameter of 0.140 mm  
 163 and chamber density of  $22.6$  kg/m<sup>3</sup> comparing nitrogen and sulfur hexafluoride. In the left  
 164 hand side, these results are seen in terms of spray penetration and spray tip Mach for the  
 165 50 MPa injection pressure condition. On each curve, the vertical lines represent the  
 166 measurement standard deviation from all the repetitions taken. As it can be seen, despite  
 167 the lower speed of sound in SF<sub>6</sub>, the spray is always in subsonic regime (no shockwave is  
 168 generated) and the spray penetration achieved is practically equal for both tests. On the  
 169 contrary, in the right hand side of the figure the injection pressure is increased to 130 MPa

170 and supersonic conditions are reached for the SF<sub>6</sub> environment. In this case, it can be  
 171 observed how higher spray penetration compared to nitrogen environment is obtained.



172  
 173 Figure 4. Spray penetration and spray tip Mach results for 0.140 mm nozzle at injection pressures of 50  
 174 MPa (left) and 130 MPa (right).

#### 175 4. Spray penetration correction in supersonic conditions

176 As stated in section 1, one of the main aspects affecting the spray dynamics is the density  
 177 ratio between the fuel and the ambient. In the case of the current experiments, this ratio is  
 178 affected only by the gas density, since fuel composition and temperature are maintained  
 179 constant. When a shockwave appears during the spray development, a gas density gradient  
 180 across this shockwave is produced. This density gradient depends mostly on the spray tip  
 181 Mach number and the kind of shockwave created. In particular, for a frontal or bow-shape  
 182 wave the following relationship can be used assuming one-dimensional flow [23]:

$$\frac{\rho_{up}}{\rho_{ch}} = \frac{(\gamma + 1)M^2}{2 + (\gamma - 1)M^2} \quad (1)$$

183 Where  $\rho_{up}$  and  $\rho_{ch}$  are the densities upstream the shockwave (i.e. close to the spray tip) and  
184 in the rest of the chamber, respectively,  $\gamma$  is the specific heat ratio (equal to 1.0984 for SF<sub>6</sub>)  
185 and  $M$  is the spray tip Mach.

186 If instead the conditions are such that an oblique shockwave is created, the same  
187 relationship applies, but instead of using the experimental spray tip Mach, a corrected  
188 Mach,  $M_n$ , should be used, which can be calculated as [23]:

$$M_n = M \cdot \sin\beta \quad (2)$$

189 where  $\beta$  is the shockwave half angle, as described in Figure 3, which can be extracted from  
190 the post-processing images. Nevertheless, it has to be noted that for the conditions  
191 evaluated into the experimental campaign, high values of  $M$  are typically associated to low  
192 values of  $\beta$ , resulting in the fact that the change of density across the shockwave is  
193 generally below 5%.

194 According to equation (1), when supersonic conditions are reached, the density reduces  
195 locally in the vicinity of spray. From spray propagation theory, it is known that the density  
196 affects penetration with an exponent equal to -0.25 [3,24], so lower density would be  
197 translated into a larger spray penetration. Knowing this, it is possible to estimate the  
198 supersonic spray penetration at each time step from the subsonic one as:

$$S_{sup}(t) = S_{sup}(t - dt) + [S_{sub}(t) - S_{sub}(t - dt)] \left( \frac{\rho_{up}(t)}{\rho_{ch}} \right)^{-0.25} \quad (3)$$

199 where  $S_{sup}$  is the spray penetration at supersonic conditions,  $S_{sub}$  is the subsonic one,  $t$  is  
200 the time elapsed after the start of injection and  $dt$  is the time step used to discretize the  
201 spray penetration, which depends on the camera velocity for visualization tests. In this  
202 expression, the density ratio is calculated at each time step depending on the  
203 corresponding spray tip Mach, using expressions (1) and (2) according to the kind of

204 shockwave produced. As a simplification, for the current study it is assumed that for  
 205  $1 < M \leq 1.5$ , the shockwave is frontal, while for  $M > 1.5$  it is oblique, which well matches the  
 206 overall behavior observed in the experiments, arriving to expressions (4) and (5):

$$S_{sup}(t) = S_{sup}(t - dt) + [S_{sub}(t) - S_{sub}(t - dt)] \left( \frac{(\gamma + 1)(M(t) \cdot \sin\beta(t))^2}{2 + (\gamma - 1)(M(t) \cdot \sin\beta(t))^2} \right)^{-0.25} \quad (4)$$

*if*  $M(t) > 1.5$

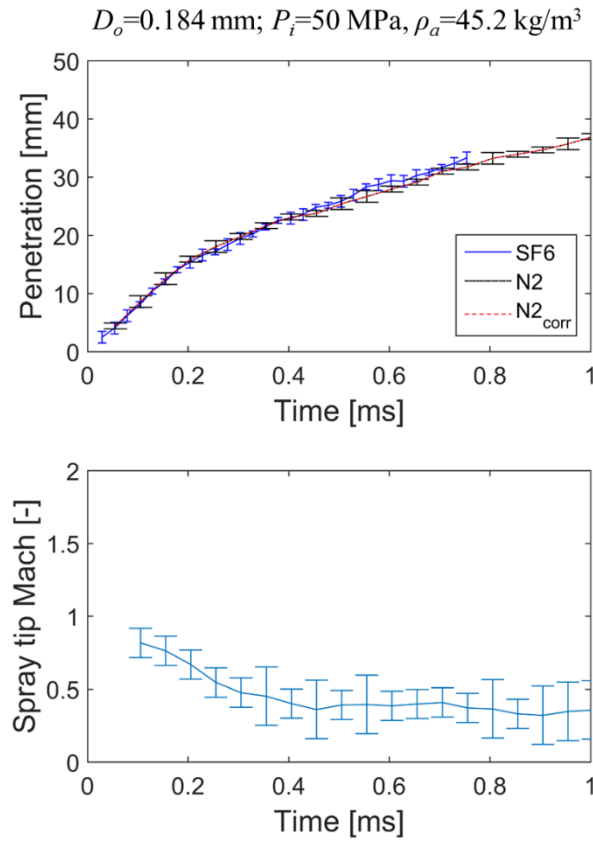
$$S_{sup}(t) = S_{sup}(t - dt) + [S_{sub}(t) - S_{sub}(t - dt)] \left( \frac{(\gamma + 1)M(t)^2}{2 + (\gamma - 1)M(t)^2} \right)^{-0.25} \quad (5)$$

*if*  $1 < M(t) \leq 1.5$

207 Taking into account these considerations, the next step is to evaluate if the differences  
 208 appreciated into the spray tip penetration between SF<sub>6</sub> and N<sub>2</sub> environment can be  
 209 explained in terms of this density gradient. This analysis is performed in Figures 5-7. On  
 210 each of these figures, the following data curves are depicted:

- 211 - Experimental spray penetration in N<sub>2</sub> (upper graph, dashed black line) and SF<sub>6</sub>  
 212 (upper graph, continuous blue line). The lines represent the average values, while  
 213 the error bars are the experimental uncertainty estimated from the standard  
 214 deviation along the eight repetitions.
- 215 - Theoretical spray tip Mach (bottom graph). It represents the Mach that the spray  
 216 tip would reach considering the tip velocity obtained from the N<sub>2</sub> measurements  
 217 penetration gradient, but considering the speed of sound characteristic of SF<sub>6</sub>.
- 218 - Corrected (or supersonic) N<sub>2</sub> penetration curve (upper graph, dashed red line),  
 219 calculated using equations (4) and (5). The aim of this correction is to see if the  
 220 methodology described previously can help to predict the supersonic penetration  
 221 from the subsonic one.

222 The first results presented in Figure 5 correspond to a subsonic case, reached with a low  
223 injection pressure and a high ambient density. As it could be expected, when the flow  
224 regime is subsonic, no density gradient appears, and all three curves are practically equal.

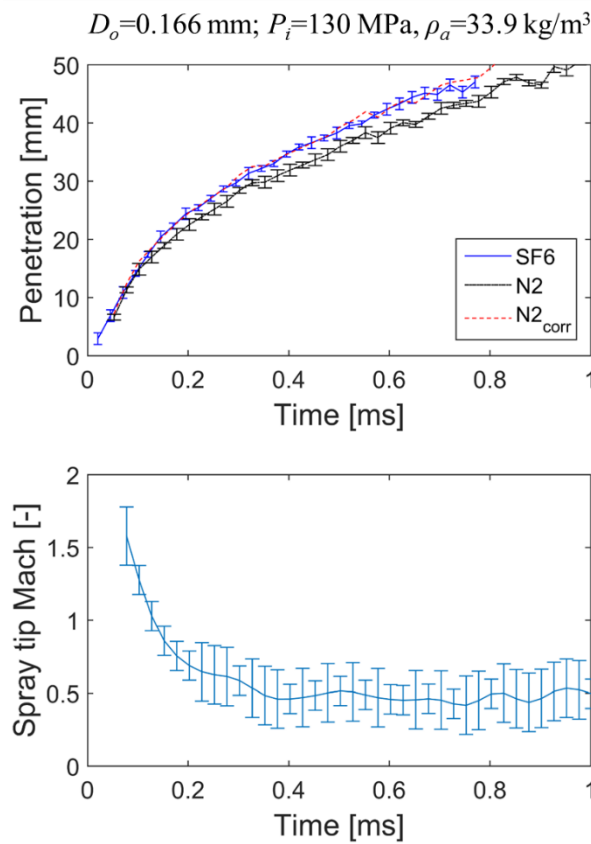


225

226 *Figure 5. Spray penetration evaluation for  $D_o=0.184 \text{ mm}; P_i=50 \text{ MPa}, \rho_a=45.2 \text{ kg/m}^3$ .*

227 Figure 6 represents a case where supersonic conditions are reached, mostly due to the  
228 increase of injection pressure. Nevertheless, the maximum  $M$  is around 1.5, which  
229 according to the previous analysis would imply that only normal shockwaves would be  
230 produced. From the experimental values, it can be clearly seen how the spray penetration  
231 is larger for the SF<sub>6</sub> tests, as already anticipated from the density ratio expressions. Once  
232 the correction is applied to the N<sub>2</sub> results, both curves present again a very good  
233 agreement. This can be seen as a confirmation that the most significant differences

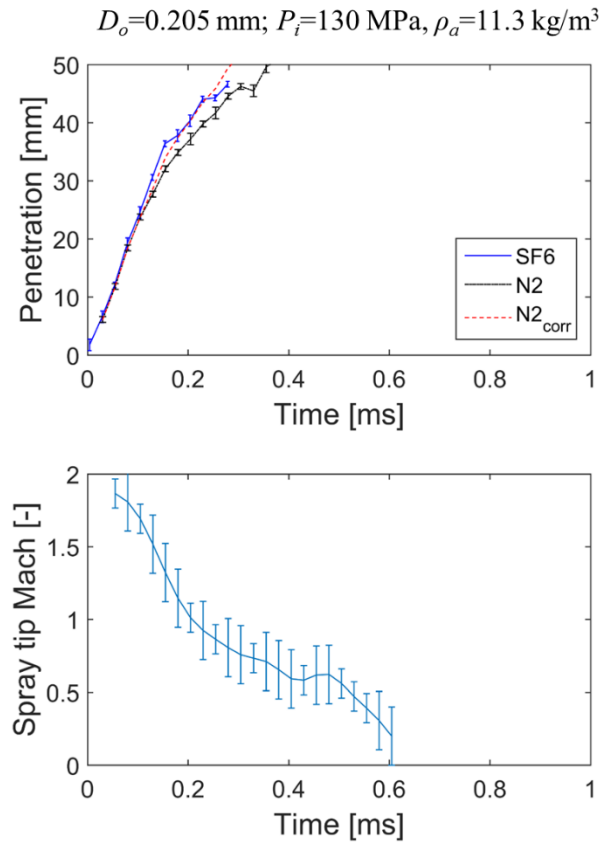
234 observed in terms of spray penetration are directly related to the density gradient produced  
235 by the shockwave.



236

237 *Figure 6. Spray penetration evaluation for  $D_o=0.166 \text{ mm}; P_i=130 \text{ MPa}, \rho_a=33.9 \text{ kg/m}^3$ .*

238 Finally, Figure 7 shows a case with a lower ambient density and a larger nozzle diameter,  
239 which produce faster spray penetration and higher Mach. In this case, it can be seen how  
240 both experimental curves are very similar up to  $\sim 0.1 \text{ ms}$  after the start of injection. This  
241 behavior is quite well reproduced also by the corrected spray penetration, since oblique  
242 shockwaves with lower density ratios would be predicted. Once the spray tip velocity  
243 reduces and the shockwave transitions from oblique to normal, the two experimental  
244 curves detach, behavior that is also well predicted according to the transition between  
245 equations (4) and (5).



246

247

*Figure 7. Spray penetration evaluation for  $D_o=0.205 \text{ mm}; P_i=130 \text{ MPa}, \rho_a=11.3 \text{ kg/m}^3$*

248

While in the current study the subsonic penetration has been directly taken from the N<sub>2</sub>

249

experiments, the advantage of this methodology is that there are well established models

250

and experimental correlations that can be used to accurately estimate this penetration.

251

Nevertheless, these predictions could significantly underestimate the spray penetration if

252

supersonic conditions are reached. According to the results just discussed, this

253

underprediction can be easily corrected by calculating the spray tip Mach and the

254

corresponding density ratio, as previously described.

255

Figure 8 shows the values of spray tip Mach and the subsequent spray penetration

256

correcting factor that could be expected in engine environment (with air as working gas)

257

for the fuel injector with 0.119 mm diameter, which is the most representative of a light-

258

duty application. This spray tip Mach is computed as a function of the combustion

259 chamber temperature, which affects the speed of sound, and the fuel injection pressure,  
 260 controlling the spray tip penetration. The temperature is changed in a range from 300 K  
 261 (approximately room temperature, as the conditions tested in the current work) to 1000K.  
 262 In the case of the injection pressure, the range has been extended up to 350 MPa, which  
 263 is seen as a reasonable future capability for diesel fuel injection systems considering that  
 264 there are already production samples operating at 300 MPa. The spray penetration law  
 265 used to obtain the spray tip Mach is obtained using the model proposed by Desantes et al.  
 266 [25], which derives the penetration from the spray momentum. This spray momentum has  
 267 been estimated as a function of the fuel injection pressure using the following expression:

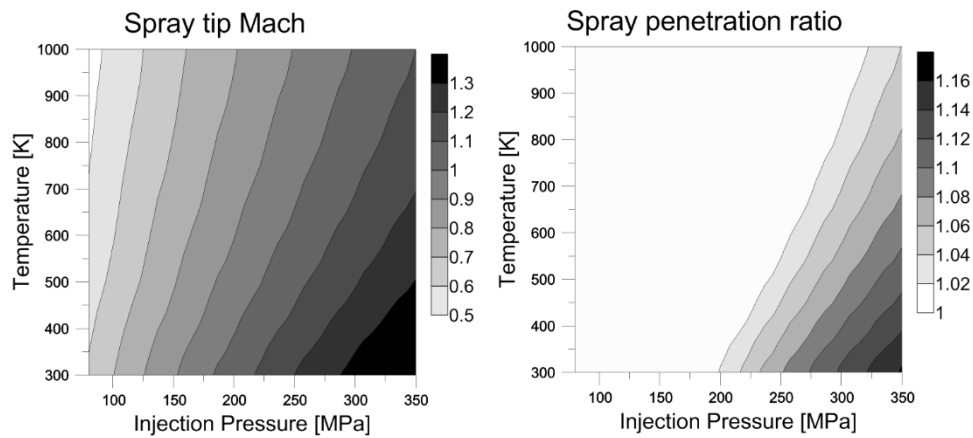
$$\dot{M}_o = \rho_f A_{eff} u_{eff}^2 = \rho_f C_a A_o C_v^2 u_{th}^2 = 2C_a A_o C_v^2 (P_i - P_{ch}) \quad (6)$$

268 Where  $\dot{M}_o$  is the spray momentum at the nozzle outlet,  $\rho_f$  is the fuel density,  $A_{eff}$  is the  
 269 nozzle effective outlet area,  $u_{eff}$  is the nozzle effective outlet velocity,  $C_a$  and  $C_v$  are the  
 270 nozzle area and velocity coefficients,  $A_o$  is the nozzle geometrical outlet area,  $u_{th}$  is the  
 271 nozzle outlet velocity according to Bernoulli's equation and  $P_i$  and  $P_{ch}$  are the injection  
 272 and chamber pressure, respectively [26]. Since no experimental information regarding  
 273 area and velocity coefficients is available for most of the injection pressure range  
 274 evaluated, since it is out of the fuel injector capabilities, constant values of 0.98 and 0.9  
 275 have been used for this estimation. Regarding the chamber pressure, also a constant value  
 276 of 5 MPa has been selected as a simplification. Finally, a constant spray cone opening  
 277 angle of 20° was also considered to compute the spray tip penetration as a simplification.

278 Looking at this figure, it can be seen how up to 200 MPa injection pressure the conditions  
 279 are always subsonic, so no correction is necessary. On the contrary, for higher injection  
 280 pressures supersonic conditions appear, and for 350 MPa spray penetration would always



281 be impacted regardless the chamber temperature used. The impact on spray penetration  
282 can be as high as 16% for the most critical conditions evaluated.



283

284 *Figure 8. Analysis of potential spray tip Mach and spray penetration correction in engine environment*  
285 *vs. temperature and injection pressure.*

286

287 However, it has to be noted that the information in Figure 8 has two main simplifications  
288 with respect to real engine-like operation. First, this calculation assumes constant  
289 thermodynamic conditions (thus density). In reality, pressure and temperature in the  
290 combustion chamber vary as a consequence of the piston compression. Nevertheless, the  
291 use of multiple injections with short durations in current diesel engines make that the  
292 variation of these thermodynamic conditions during the spray propagation time is very  
293 small, so the impact should be negligible, and the only uncertainty in this regard for  
294 engine operation would be the determination of the combustion chamber temperature  
295 needed to evaluate the density. The most significant simplification instead comes from  
296 the fact that the so called free spray assumption is taken. This means that spray  
297 propagation is not affected by the combustion chamber air motion nor by the interaction  
298 with the combustion chamber walls. While for heavy duty applications, which are  
299 characterized by large bores and quiescent combustion systems (i.e. with low swirl

300 motion) this assumption may be acceptable, the effects of spray-wall interactions and air  
301 motion should be considered for other applications.

302

### 303 **5. Conclusions and future works.**

304 In the current study, an investigation of diesel spray penetration characteristics on  
305 supersonic conditions has been performed. These supersonic conditions have been  
306 obtained by injecting into a sulfur hexafluoride environment, characterized by a lower  
307 speed of sound. The following conclusions have been drawn:

- 308 - From the SF<sub>6</sub> visualization tests, two different kinds of shockwaves have been  
309 detected: oblique, when the spray tip velocity is high, and normal, for lower values.  
310 The transition between both regimes can be found around a spray tip Mach  
311 approximately equal to 1.5.
- 312 - Spray tip penetration is generally higher in SF<sub>6</sub> environment than in N<sub>2</sub> once  
313 supersonic conditions are reached. This higher penetration is associated to the  
314 density ratio appearing across the shockwave, producing a lower density than  
315 anticipated in the spray tip vicinity.
- 316 - A methodology based on one-dimensional flow approximation for supersonic  
317 shockwaves and the typical exponent for the chamber density in spray penetration  
318 correlations is proposed to extrapolate supersonic spray penetration from subsonic  
319 values, showing good agreement with the experimental data.

320 From the current study, three main paths for future works arise. The first would be to  
321 perform the supersonic spray characterization at ultra-high injection pressure (over 200  
322 MPa) in order to confirm that the methodology proposed is capable to accurately predict  
323 these conditions. The second proposal for future work would be to explore close-coupled

324 multiple injections. In this case, the second injection would encounter an environment  
325 partially filled with vaporized diesel fuel, which would reduce the local speed of sound  
326 and potentially enhance the supersonic conditions at the same injection velocity. Finally,  
327 another potential field of interest would be to extend this work to an optical engine, where  
328 the effects of air motion and spray-wall interactions can be accounted for.

329

330

### 331 NOMENCLATURE

$A_{eff}$	Nozzle outlet effective area.
$A_o$	Nozzle outlet area.
$C_a$	Nozzle area coefficient.
$C_v$	Nozzle velocity coefficient.
$D_o$	Injector nozzle outlet diameter.
$dt$	Time step for the visualization tests.
$M$	Spray tip Mach.
$\dot{M}_o$	Momentum flux at the nozzle outlet.
$P_{ch}$	Discharge chamber pressure.
$P_i$	Injection pressure.
$S$	Spray tip penetration.
$S_{sub}$	Spray tip penetration in subsonic conditions.
$S_{sup}$	Spray tip penetration in supersonic conditions.
$t$	Time elapsed after the start of injection.
$u_{eff}$	Nozzle outlet effective velocity.
$u_{th}$	Nozzle outlet theoretical velocity according to Bernoulli's equation.

### 332 Greek symbols

$\rho_f$	Fuel density.
$\rho_a$	Air density.
$\gamma$	Specific heat ratio.
$\beta$	Oblique shockwave angle
$\rho_{ch}$	Discharge chamber density far from the shockwave.
$\rho_{up}$	Discharge chamber density upstream the shockwave.

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