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Additional Information

CFD analysis of the HVAD's hemodynamic performance and blood damage with insight into gap clearance

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Abstract – Mechanical circulatory support using ventricular assist devices has become commonplace in the treatment of patients suffering from advanced stages of heart failure. While blood damage generated by these devices has been evaluated in depth, their hemodynamic performance has been investigated much less. This work presents the analysis of the complete operating map of a left ventricular assist device, in terms of pressure head, power and efficiency. Further investigation into its hemocompatibility is included as well. To achieve these objectives, computational fluid dynamics simulations of a centrifugal blood pump with a wide-blade impeller were performed. Several conditions were considered by varying the rotational speed and volumetric flow rate. Regarding the device's hemocompatibility, blood damage was evaluated by means of the hemolysis index. By relating the hemocompatibility of the device to its hemodynamic performance, the results have demonstrated that highest hemolysis occurs at low flow rates, corresponding to operating conditions of low efficiency. Both performance and hemocompatibility are affected by the gap clearance. An innovative investigation into the influence of this design parameter has yielded decreased efficiencies and increased hemolysis as the gap clearance is reduced. As a further novelty, pump operating maps were non-dimensionalized to highlight the influence of Reynolds number, which allows their application to any working condition. The pump's operating range places it in the transitional regime between laminar and turbulent, leading to enhanced efficiency for the highest Reynolds number.

Keywords – Centrifugal blood pump, Operating map, Non-dimensional analysis, Gap clearance, Shear stress, Hemolysis

Declarations – Not applicable.

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1. INTRODUCTION

Cardiovascular diseases are one of the most frequent causes of death worldwide (World Health Organization 2014), with increasing prevalence in most developed countries (Mosterd and Hoes 2007). Despite the several available treatment options (Almenar et al. 2011), until recent decades patients afflicted by the most advanced stages of heart failure (HF) could not benefit from any medical therapy aside from transplants. However, availability of compatible heart donors is highly limited. Consequently, Mechanical Circulatory Support (MCS) using blood pumps has been developed as an alternative therapy that is being increasingly applied. First blood pumps consisted in extracorporeal circulation devices that provided temporary support, whereas more recent Ventricular Assist Devices (VADs) have made possible a longer-term and more comfortable assistance due to their smaller size (McKellar 2020).

VADs are mechanical blood pumps designed to assist or completely replace the function of left, right or both ventricles, in order to grant the required cardiac output (CO) when the heart itself is not able to provide it (Fraser et al. 2011). Left Ventricular Assist Devices (LVADs) aid in the pumping of the left ventricle (LV). Implanting blood pumps has become a standard therapy for patients suffering HF (Larose et al. 2010), as a bridge to transplant or to recovery (Stehlik et al. 2010) as well as a destination therapy (Kirklin et al. 2014), and nowadays it is the main treatment option when transplantation is not possible. Nevertheless, these devices are associated with complications that must be overcome, such as bleeding and stroke (Bluestein et al. 2010; Al-Quthami et al. 2012).

Two main aspects must be considered when designing and evaluating a VAD: hemodynamic performance and hemocompatibility. On the one hand, hemocompatibility is related to the blood damage caused by the device (Wiegmann et al. 2018) owing to non-physiological flow conditions such as device-induced turbulence (Avci et al. 2021). On the other hand, the performance of a hydraulic pump is evaluated by means of pressure head, power consumption and efficiency. These variables depend on the operating condition (rotational speed and flow rate), which justifies the importance of obtaining the entire operating map. The efficiency quantifies the hydraulic energy increment of the flow through the pump in percentage of the shaft power. Typical values of VAD's efficiency are in the range 20-30% (Fraser et al. 2011). Therefore, most of the consumed energy is not used to increase the pressure head of the flow. Increasing efficiency would imply a reduction in losses, as well as a prolongation of the battery life. Moreover, improving the hemodynamic performance to enhance the pump's efficiency constitutes a crucial task that will simultaneously improve the hemocompatibility of the device.

Numerous works have been published investigating different constant-flow devices based on axial and centrifugal rotary blood pumps. Nevertheless, while blood trauma is deeply investigated in every work regarding VADs, hardly any authors include an

42 evaluation of the pump's hemodynamic performance through its complete operating
43 map. Thamsen et al. (2015) analysed the blood damage generated within two LVADs:
44 HeartMate II (axial) and HeartWare (centrifugal); and obtained similar hemolysis for
45 both pumps, with the highest damage produced at the edges of blades and in the gap
46 regions respectively. Zhang et al. (2020) analysed both normal and hypertension
47 conditions in several pumps and obtained a 30% increase of hemolysis levels under
48 the latter. Furthermore, a potentially higher risk of blood trauma operating at low-flow
49 conditions is reported by several authors. Granegger et al. (2020) compared the
50 hemocompatibility of HeartWare for adult (high-flow) and pediatric (low-flow) operating
51 conditions, and discovered larger stagnation zones and residence times within the
52 pump for the pediatric case, resulting in higher hemolysis. Thamsen et al. (2020)
53 investigated the flow within the HeartMate III (centrifugal) under two flow conditions
54 representing the lower (diastole) and upper (systole) physiologic flow range. Schöps
55 et al. (2021) studied the blood trauma associated to an extracorporeal membrane
56 oxigenation (ECMO) device under low- and high-flow operating conditions. Both
57 obtained larger recirculation zones and more disturbed flow fields inside the pumps
58 and, therefore, increased hemolysis for the low-flow condition. Although different
59 operating conditions were considered in these works, the complete map was not
60 obtained and the performance variables were not evaluated. The blood flow rate
61 through the pump depends on patient's arterial pressure since it operates at a constant
62 speed. Thus, studying its complete operating map, in terms of both hemodynamic
63 performance and hemolysis, is of vital importance because it will work within a range
64 of operating conditions. Wang et al. (2019) and Wu et al. (2021) presented the map of
65 H-Q curves, which relates pressure head with operating conditions of a centrifugal
66 device. However, it was only used for validation purposes, and the results of both
67 studies focused on hemocompatibility. In this sense, the current work will contribute
68 by analysing the complete operating map of a LVAD, even in terms of non-dimensional
69 variables, and relating it to the device's hemocompatibility.

70 In this paper, blood flow within the Medtronic's continuous flow centrifugal pump
71 HeartWare VAD (HVAD) is analysed. This pump employs a hybrid levitation system
72 to position the impeller through the balance of magnetic and hydrodynamic forces. It
73 has a 4-wide-blade impeller whose large top and bottom surfaces are tapered to
74 produce hydrodynamic lift, while magnetic bearing is generated between center post's
75 coils and impeller's magnets (Foster 2018). The hydrodynamic lift is speed dependent,
76 leading to different gap clearances as the rotational speed changes. Hybrid and fully
77 magnetic levitation bearings have replaced mechanical bearings to avoid friction,
78 heating and dynamic sealing, which reduces blood damage risk (Wu et al. 2021).
79 Nonetheless, HVAD was recently withdrawn from the market owing to elevated
80 thrombogenicity. The aim of this work is to investigate in depth the HVAD's
81 hemodynamics and to clarify the causes of its thrombotic complications. The thrombus
82 begins to form with the chemical or mechanical activation of platelets, and then it grows
83 in zones where activated platelets are prone to deposit. Note that activation and

84 deposition might occur at different locations. Regions of the flow field characterized by
85 high velocity gradients are exposed to elevated non-physiological shear stresses
86 during certain exposure times, giving rise to potential risk of hemolysis and platelet
87 mechanical activation (Wiegmann et al. 2019). Recirculation and stagnation regions,
88 which involve large residence and exposure times, are prone to hemolysis, platelet
89 deposition and thrombus growth. Elevated values of plasma free hemoglobin, caused
90 by hemolysis, also promote the coagulation cascade (Bartoli et al. 2018). Considering
91 these relations, thrombosis will be qualitatively related to the velocity and hemolysis
92 fields. In this work, both hemodynamic and hemolytic performance will be related to
93 shear stresses produced within gaps, and the influence of gap clearance will be
94 discussed as a novel contribution to the knowledge about the HVAD. Furthermore, the
95 non-dimensional analysis of operating maps will help in understanding the device's
96 performance.

97 **2. MATERIALS AND METHODS**

98 *2.1 Computational modeling*

99 *2.1.1 Solver and governing equations*

100 Computational Fluid Dynamics simulations are performed to evaluate the flow within
101 the HVAD, using the software SimCenter STAR-CCM+ (Siemens). The blood is
102 modeled as a liquid with density $\rho = 1060 \text{ kg/m}^3$ and assumed to be a Newtonian
103 fluid, i.e. with constant viscosity $\mu = 3.5 \text{ mPa} \cdot \text{s}$, since its non-Newtonian behavior is
104 expected to be negligible at shear rates $\dot{\gamma} > 100 \text{ s}^{-1}$ such as those found in VADs
105 (Fraser et al. 2011; Wiegmann et al. 2018, 2019; Wang et al. 2019). There are zones
106 within the pump, however, where this assumption does not apply owing to lower shear
107 rates. Consequently, additional computations have been performed to ensure
108 applicability of this assumption, resulting in acceptable discrepancies: relative errors
109 were below 1% for all the performance variables, except at the low-flow condition
110 which resulted in errors of 6% and 9% for pressure head and efficiency respectively.

111 Given a working fluid, the fluid dynamics features within the device are entirely
112 defined by the rotational speed of the impeller (Ω) and volumetric flow rate (Q). Several
113 operating conditions are considered by varying these parameters in range of
114 [1800,4000] rpm and [1,10] L/min respectively. The flow field through the pump is not
115 affected by the mean pressure since the working fluid is incompressible. Therefore,
116 the reference stagnation pressure of $p_t = 760 \text{ mmHg}$ is imposed at the inlet. Mass flow
117 equal to $\dot{m} = \rho Q$ at the outlet and no-slip walls are imposed as further boundary
118 conditions.

119 The simulations are performed using both steady and transient approaches. For the
120 steady-state simulations, the impeller motion is imposed by means of a Moving
121 Reference Frame (MRF) approach. It consists in formulating the set of equations for

122 the rotating region in a reference frame that moves with the angular velocity of the
123 impeller, neglecting transient effects and allowing its solution in a steady-state
124 simulation. This is achieved by including an additional term of inertial forces in the
125 momentum equation. The MRF approach is employed together with a mixing plane
126 interface between rotating and static regions to avoid circumferential heterogeneity
127 due to the frozen impeller (Galindo et al. 2020). For the unsteady-state simulations,
128 the motion is reproduced using a sliding mesh approach and time is discretized
129 employing a second order approach. The time-step is chosen corresponding to 0.5
130 deg of rotation with 10 inner iterations per step. This time-step leads to convective
131 Courant numbers $CFL < 1$ in 96% of the fluid volume. Up to ten revolutions are needed
132 to achieve a cyclic solution. Time-averaged values of the performance variables are
133 obtained from the last revolution for the analysis.

134 Reynolds-Averaged Navier-Stokes (RANS) and Unsteady RANS (U-RANS)
135 equations are solved to obtain the mean flow solution, in steady and transient
136 approaches respectively. These equations are derived from the complete set of mass,
137 momentum and energy conservation equations by imposing the Reynolds
138 decomposition. The mass and momentum equations for the mean flow are solved
139 assuming incompressible flow, while the energy equation is not needed. These
140 equations can be found in Pope (2001). The momentum equation incorporates a
141 source term of inertial and body forces which corresponds to the additional term
142 included by the MRF approach, calculated as described by Torregrosa et al. (2019).

143 To achieve closure of RANS and U-RANS equations, the Reynolds stresses must
144 be estimated using an appropriate turbulence model. Several turbulence models have
145 been compared, resulting in differences below 3% on the performance variables. The
146 k- ω model with shear stress transport (SST), proposed by Menter (1993), is selected,
147 which is employed by most authors (Gross-Hardt et al. 2019) and is considered the
148 standard model for turbomachinery applications. This turbulence model solves two
149 additional transport equations and, hence, requires the imposition of two turbulence
150 boundary conditions. The turbulence intensity and turbulent length scale at the inlet
151 are set to be $I_t = 0.01$ and $L_t = 0.001$ m, respectively. The effect of these parameters
152 over the solution has been investigated and found to be negligible.

153 *2.1.2 Geometry, fluid domain and mesh*

154 A CAD model of the HVAD was obtained by 3D-scanning using a HDI Advance 3D
155 scanner, which employs structured-light technology and delivers high-resolution digital
156 scans with an accuracy of 50 μm . The resulting fluid domain is presented in Fig. 1 (a),
157 showing the wide-blade impeller. Its main dimensions are listed in Table 1. Since gap
158 clearance depends on rotational speed, a parametric study is done varying its value
159 within the range specified in Table 1. The top clearance is imposed by translating the
160 impeller along the shaft axis. The resulting values of bottom gap clearance are not
161 presented in Table 1 because it has a significant radial variation due to this surface's
162 taper angle, moreover its effect is found to be negligible compared to that of the top

163 gap clearance. The geometry with a top clearance of 40 μm is considered the main
 164 configuration, based on the order of magnitude used in the literature and shown in
 165 Table 2.

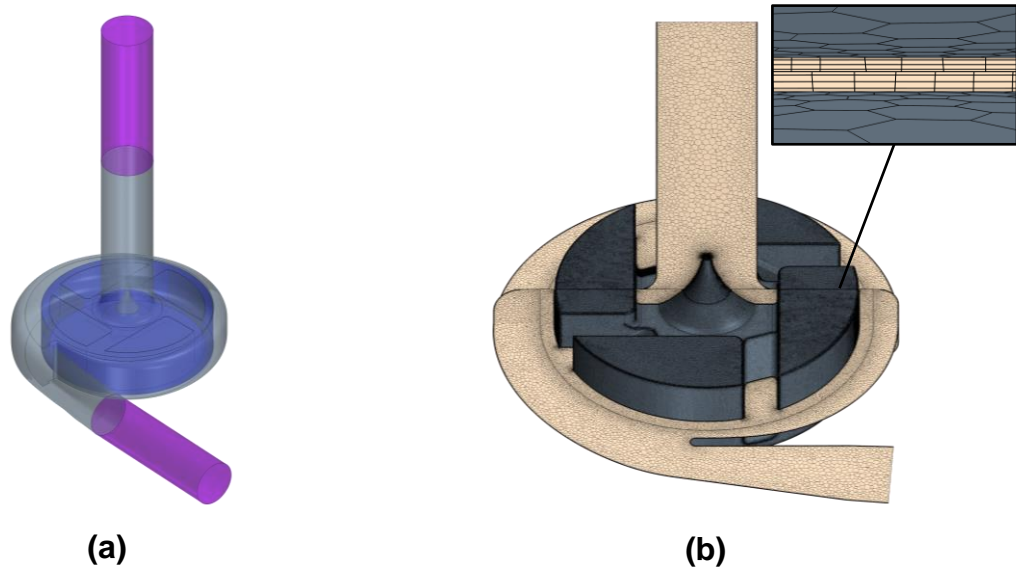


Fig. 1 Computational domain and mesh of the HVAD. **(a)** Fluid domain: static region (gray), rotating region (blue) and inlet/outlet cannulas (violet, not to scale); **(b)** Sketch of the mesh, including a zoom of the thin mesh in the gap region

Parameter	Symbol	Value	Units
Impeller diameter	D_{imp}	34.6	mm
Inlet diameter	D_{inlet}	12.75	mm
Outlet diameter	D_{outlet}	10	mm
Channel width	w	3.8	mm
Top gap clearance	c	[20,60]	μm

Table 1 Main dimensions of the HVAD and values of top gap clearance

Reference	Gap clearance
Chen et al. (2019)	40 μm
Thamsen et al. (2015)	40 μm
Granegger et al. (2020)	[20,23] μm
Foster (2018)	[40,70] μm

Table 2 Values of top gap clearance stated by other authors (Thamsen et al. 2015; Foster 2018; Chen et al. 2019; Granegger et al. 2020)

166 Polyhedral grids are generated to discretize the fluid volume. A polyhedral mesh is
 167 chosen since it ensures grid independence with fewer elements compared to
 168 tetrahedral meshes (Spiegel et al. 2011). The mesh includes a 10-element prism
 169 boundary layer along the walls, as well as local refinements in the rotating region and
 170 around the center post. A mesh independence study is performed for each
 171 configuration. Results of the mesh study for the main configuration are shown in Fig.
 172 2 and summarized in Table 3, and analogous results are obtained for the
 173 configurations with clearances of 20 and 60 μm . The mesh reference size is taken to
 174 be $8 \cdot 10^{-4}$ m with refinement sizes of around $1 \cdot 10^{-4}$ m in the rotating region, leading
 175 to a mesh of 9.5×10^6 volume cells (medium mesh), represented in Fig. 1 (b). Since
 176 a turbulence model (k- ω SST) is used, the quality of the boundary layer's mesh is
 177 evaluated by means of the wall $y^+ = \frac{u_{\tau} y}{\nu}$ (Pope 2001). Reynolds number is relatively
 178 low at every operating condition, hence y^+ must be less than 1 to ensure an accurate
 179 prediction of turbulence across the boundary layer. All grids have $y^+ < 1$ in most part
 180 of the walls, guaranteeing a correct modeling of the viscous sublayer. Specifically, the
 181 portion of wall areas with $y^+ > 1$ is 6% for the main configuration. Inlet and outlet
 182 cannulas with a longitude of 14-diameters are added in order to allocate the boundary
 183 conditions sufficiently far from the zone of interest. Note that this domain including long
 184 inlet and outlet ducts is not representative of the HVAD inserted in the heart apex, but
 185 similar to the experimental test bench configuration.

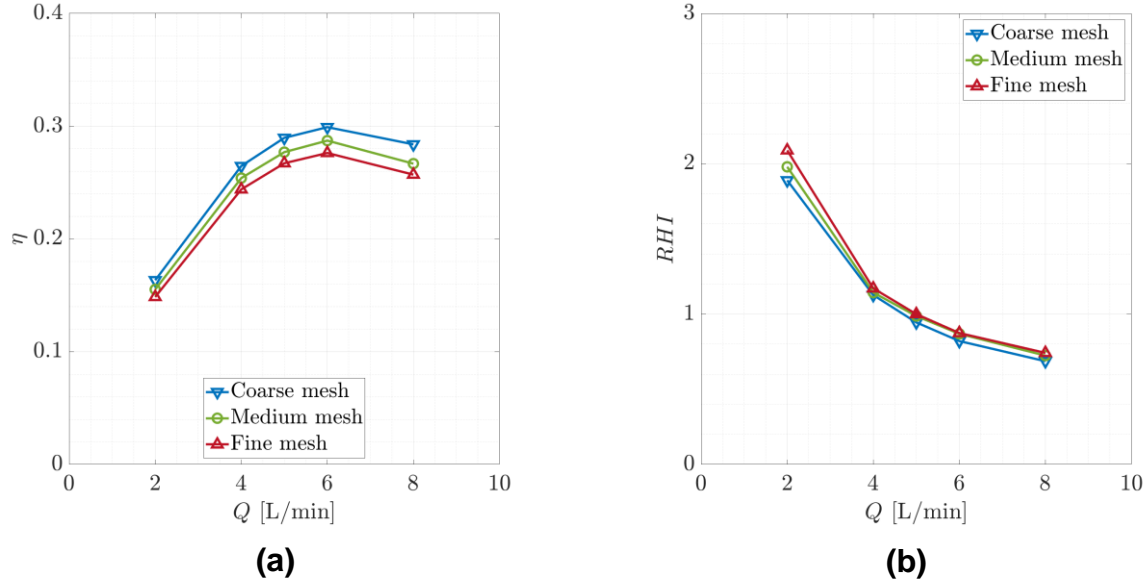


Fig. 2 Results of the mesh independence study for the HVAD's main configuration (40 μm -clearance gap) operating at $\Omega = 3000$ rpm: **(a)** Efficiency; **(b)** Relative hemolysis index, taking the hemolysis index obtained using the fine mesh and operating at $Q = 5$ L/min as the nominal hemolysis index

	Coarse mesh	Medium mesh	Fine mesh
Reference size	$1.6 \cdot 10^{-3}$ m	$8 \cdot 10^{-4}$ m	$4 \cdot 10^{-4}$ m
Number of cells	2.7×10^6	9.5×10^6	33×10^6
Mean relative error: Δp_t	4.8 %	2.2 %	—
Mean relative error: η	9.1 %	4.0 %	—
Mean relative error: HI	6.5 %	2.4 %	—

Table 3 Mesh settings and results for coarse, medium, and fine grids

186 2.1.3 Blood damage and hemolysis calculation

187 Despite the increasing use of VADs in HF patients, several device-induced
188 complications remain (Bluestein et al. 2010). As a device's size is reduced, gap widths
189 decrease and impeller speeds increase, and, consequently, flow-induced blood
190 damage becomes a more significant issue (Fraser et al. 2011). This damage is directly
191 related with the scalar shear stress (SSS) level, defined in Equation 1 where σ_{ij} are
192 the viscous stress components and u_i denotes the i^{th} component of velocity. Note that
193 SSS depends on both shear (σ_{ij}) and normal (σ_{ii}) components of stress. The Reynolds
194 stress components are not considered for SSS calculation since Reynolds stress
195 tensor is a statistical quantification of the averaged transport of fluctuating momentum,
196 and has no direct link to physical forces acting over blood cells (Ge et al. 2008).

$$SSS = \tau = \sqrt{\frac{1}{6} \sum_{i \neq j} (\sigma_{ii} - \sigma_{jj})^2 + \sum_{i \neq j} \sigma_{ij}^2} \quad 1$$

$$\sigma_{ij} = \mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

197 Hemolysis consists in the disintegration of red blood cells (RBCs) resulting in a
198 release of hemoglobin into the blood plasma. It is related to the flow-induced
199 mechanical damage to RBCs, and is normally more investigated than damage to other
200 blood components (Bluestein et al. 2010). Diverse methods have been proposed for
201 calculating the hemolysis index (HI). Most of them are power law models due to their
202 simplicity and applicability to a wide range of devices (Taskin et al. 2012). Equation 2
203 shows the expression of the power law, that relates HI with SSS (τ) and exposure time
204 (t). The HI is defined as the ratio of hemoglobin released into the blood plasma (Δhb)
205 to total blood hemoglobin concentration, assumed to be $HB = 10$ g/dL (Chen et al.
206 2019).

$$HI = \frac{\Delta hb}{HB} = C \cdot t^\alpha \cdot \tau^\beta \quad 2$$

207 Based on Equations 1 and 2, an accurate velocity field is needed to obtain a reliable
208 prediction of HI (Karimi et al. 2021). On the other hand, the non-linearity of Equation
209 2 implies that the HI at the domain outlet cannot be calculated as the sum of local
210 values of HI based on the scalar shear stress and the residence time of RBCs in each
211 grid cell (Wu et al. 2021). Thus, estimating the exposure time to a given shear stress

212 level becomes complicated since it is not equivalent to the residence time. Taskin et
 213 al. (2012) made an evaluation of different power law models, for both Eulerian and
 214 Lagrangian approaches, and stated that available methods could not predict an
 215 accurate value of HI, but they were useful in predicting relative hemolysis for
 216 comparative purposes. For the Eulerian approach, a passive scalar transport equation
 217 must be solved. Garon and Farinas (2004) and Farinas et al. (2006) demonstrated the
 218 conversion from the non-linear Equation 2 to an expression linearized in time that can
 219 be easily derived to formulate the transport equation. This conversion allows the
 220 evaluation of HI in the whole domain, avoiding the exposure time estimation.
 221 Moreover, an Eulerian approach is preferred since it considers all areas contributing
 222 to the hemolysis, whereas some regions could be omitted using Lagrangian methods
 223 based on streamlines. The scalar transport equation considered in this work is
 224 presented in Equation 3, where the passive scalar is $\Delta hb' = \Delta hb^{1/\alpha}$. Note that the
 225 diffusion term is excluded.

$$\frac{\partial(\Delta hb')}{\partial t} + u_j \cdot \frac{\partial(\Delta hb')}{\partial x_j} = (HB \cdot C \cdot \tau^\beta)^{\frac{1}{\alpha}} \quad 3$$

226 The right-hand side of Equation 3 represents the scalar source term. The values for
 227 the empirical constants are set to be $C = 1.8 \times 10^{-6} \text{s}^{-\alpha} \text{Pa}^{-\beta}$, $\alpha = 0.765$ and $\beta = 1.991$,
 228 obtained by regressing experimental data from Heuser and Opitz (1980) (Song et al.
 229 2003), since these values allow the application of the power law to wider ranges of
 230 shear stress (Taskin et al. 2012). This transport equation is solved, assuming $HI = 0$
 231 at the inlet. The device's HI is evaluated as the mass-flow-average of this parameter
 232 at the outlet, as defined in Equation 4 (Craven et al. 2019).

$$HI_{device} = \frac{\oint_{outlet} HI |\mathbf{u} \cdot d\mathbf{A}|}{\oint_{outlet} |\mathbf{u} \cdot d\mathbf{A}|} \quad 4$$

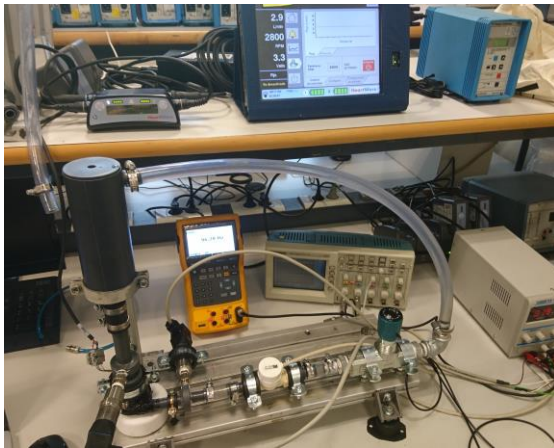
233 Since HI models do not predict an accurate absolute value, the relative hemolysis
 234 is evaluated by means of the relative hemolysis index (RHI), calculated as in Equation
 235 5 assuming $\{\Omega = 3000 \text{ rpm}, Q = 5 \text{ L/min}\}$ as the nominal operating condition.

$$RHI = \frac{HI}{HI_{nominal}} \quad 5$$

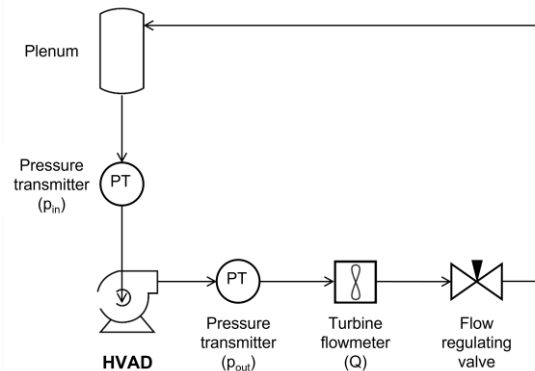
236 2.2 Experimental testing

237 The HVAD's operating map is obtained experimentally in order to validate the CFD
 238 results in terms of pressure head. The device is tested in a closed loop operating at
 239 different rotational speeds and adjusting the flow rate with a needle valve. The
 240 pressure head is measured using two high-quality pressure transmitters (WIKA PE
 241 81.61 S-20, [0,1600] bar, WIKA Instruments, Germany) allocated at around 5 cm from
 242 the device's inlet and outlet. A radial flow turbine flowmeter (RS PRO 257-133, [1.5,30]
 243 L/min, RS Components, UK) is connected downstream of the pump's outlet, at a
 244 distance of 15 cm, to measure the flow rate. A distilled water–glycerol mixture (40% of

245 glycerin) is used to simulate the blood viscosity of $\mu = 3.5 \text{ mPa} \cdot \text{s}$. The experimental
246 test bench is shown in Fig. 3 together with a sketch of the flow loop.



(a)



(b)

Fig. 3 Experimental set-up: **(a)** Photography of the test bench showing a plenum upstream of the pump, inlet and outlet pressure sensors, the flowmeter, a flow regulating valve and the HVAD's controller and monitor; **(b)** Sketch of the flow loop

247 3. RESULTS AND DISCUSSION

248 3.1 Validation of the model

249 Fig. 4 shows the pump's pressure head (Δp_t) against the volumetric flow rate for
250 three values of rotational speed, comparing the results from CFD steady-state
251 simulations to the experimental measurements. An acceptable degree of
252 discrepancies is achieved since the maximum relative error is less than 5% of the
253 calculated value. Note that the device's pressure head is validated solely, whereas
254 experimental validation for the mechanical shaft power and for the HI have not been
255 performed.

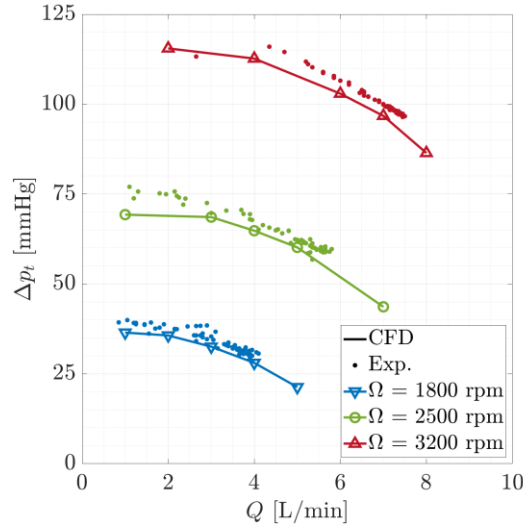


Fig. 4 Pressure head against volumetric flow rate for several values of rotational speed: curves obtained numerically compared to points measured experimentally

256 Additionally, the pump's operating maps obtained in steady and transient
 257 approaches are compared. The time-averaged values during the last revolution are
 258 considered for the transient approach.

259 The pressure head map is represented in Fig. 5 (a) for both steady- and unsteady-
 260 state simulations. Despite differences found at low and high flow rates with maximum
 261 relative errors of 6% and 10% respectively, the steady approach agrees reasonably
 262 well with transient results, principally at nominal operating conditions.

263 The efficiency (η) is defined as the hydraulic energy increment of the flow through
 264 the pump ($\Delta\dot{E}$) in percentage of the mechanical shaft power (P), as shown in Equation
 265 6 where T represents the torque acting on the impeller.

$$\eta = \frac{\Delta\dot{E}}{P} = \frac{Q\Delta p_t}{\Omega T} \quad 6$$

266 The pump's efficiency map is represented in Fig. 5 (b) for both approaches, showing
 267 discrepancies even at nominal operating conditions, with a 10% of maximum relative
 268 error between steady and transient predictions.

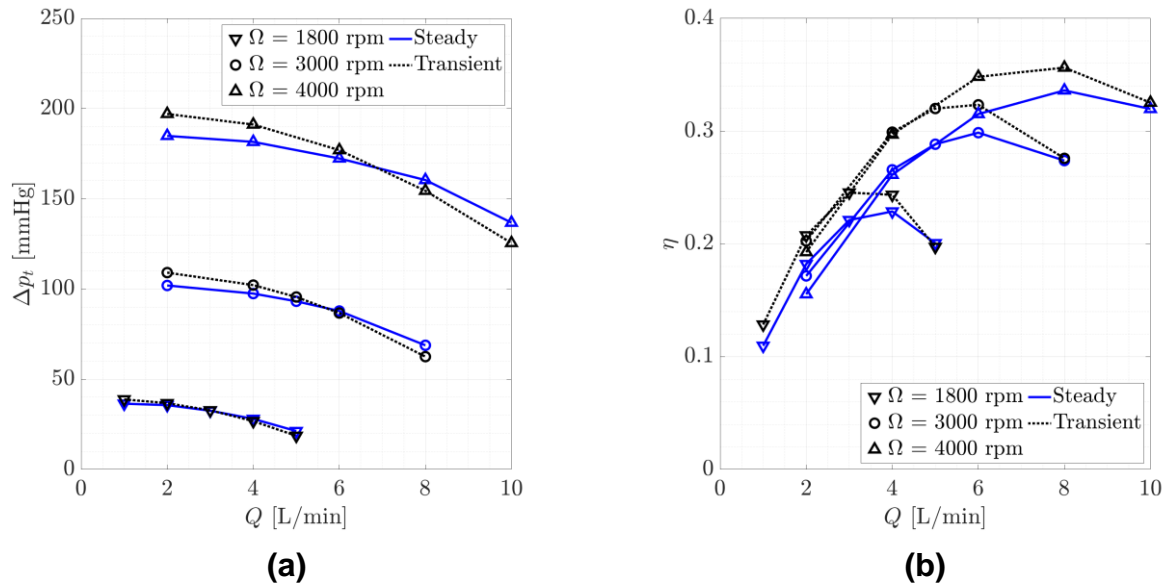


Fig. 5 Performance variables against volumetric flow rate for several values of rotational speed: comparison between steady and transient approaches. **(a)** Pressure head; **(b)** Efficiency

269 The device's RHI curve, operating at $\Omega = 3000$ rpm, is represented in Fig. 6 for both
 270 steady and transient approaches, and a constant relative error of 10% is obtained.

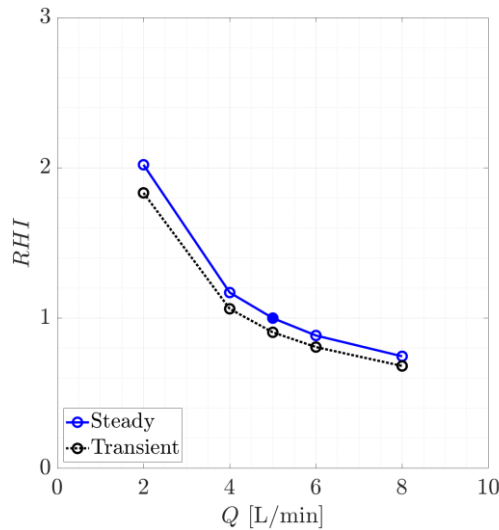


Fig. 6 Relative hemolysis index against volumetric flow rate, operating at $\Omega = 3000$ rpm: comparison between steady and transient approaches, taking the hemolysis index obtained using the steady approach and operating at $Q = 5$ L/min as the nominal hemolysis index

271 Based on the experimental validation, the CFD model predicts the HVAD's
 272 performance with enough accuracy. Despite discrepancies found between steady and
 273 transient approaches, the steady MRF approach is used for the evaluation of the

274 following sections, owing to its considerably lower computational cost. All the
275 simulations were performed on an Intel® Xeon® Gold 6248R CPU, using 48 parallel
276 processes. Computational times were 8.7 s per iteration and 163.2 s per time-step in
277 steady- and unsteady-state simulations respectively, resulting in 2 to 6 hours of
278 calculation for the steady approach, and around 330 hours of calculation for the
279 transient approach.

280 3.2 Pump's operating map

281 The hemodynamic performance of the HVAD is evaluated, for the main configuration
282 (40 μ m-clearance gap), by means of its operating maps. The pressure head and the
283 shaft power, against the volumetric flow rate, are represented in Fig. 7 for different
284 rotational speeds: $\Omega = 1800$ rpm and $\Omega = 4000$ rpm are the minimum and maximum
285 speeds respectively, while $\Omega \in [2400, 3200]$ rpm represents the pump's normal
286 operating range. The pressure head equals the difference between afterload (aortic
287 pressure) and preload (LV pressure). The mean arterial pressure (MAP) is the
288 temporal mean of the aortic pressure during one cardiac cycle. Hence, the appropriate
289 pressure head is determined by subtracting the LV pressure to the patient's MAP.
290 Although the LV pressure is fluctuating, it can be assumed to be constant and equal
291 to the LV end-diastolic pressure, $p_{LV} \approx 14$ mmHg (Jain et al. 2019), since the failing
292 heart is hardly pumping. The desired flow rate, for its part, is given by the required CO.
293 For a healthy adult weighing 70 kg, normal values of CO and MAP at rest are $CO \approx$
294 5 L/min and $MAP \in [70, 100]$ mmHg. Based on Fig. 7 (a), known as H-Q curves, these
295 values establish the normal operation range for rotational speed within $\Omega \in$
296 $[2400, 3200]$ rpm, as indicated previously. These H-Q curves follow the typical
297 tendency for centrifugal pumps.

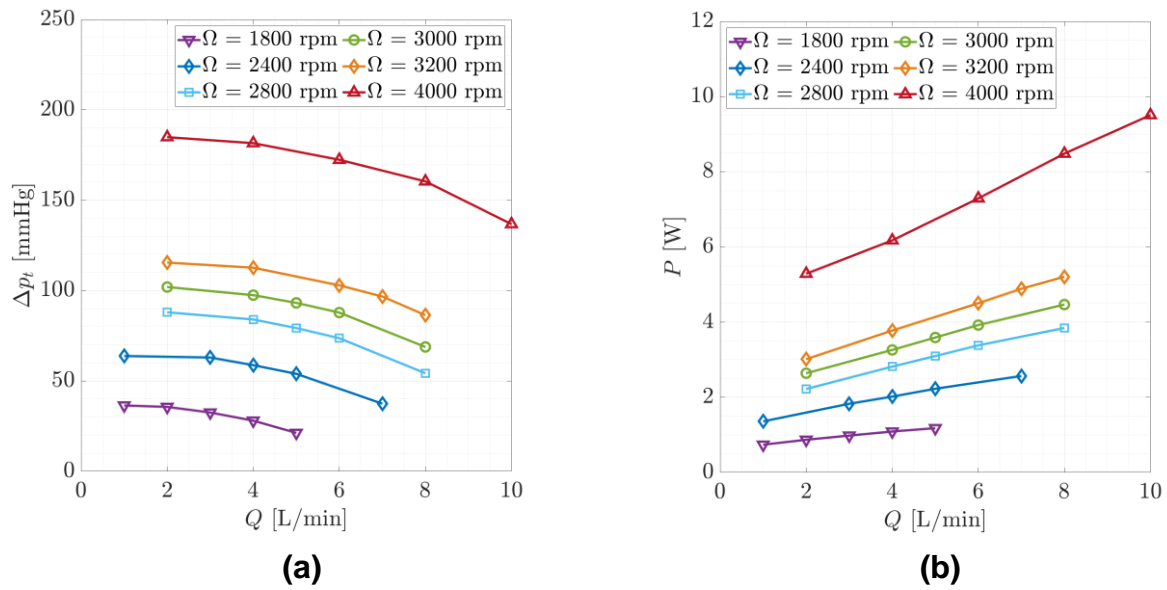


Fig. 7 Performance variables against volumetric flow rate for several values of rotational speed, for the configuration with 40 μ m-clearance gap, and blood with normal hematocrit ($\mu = 3.5$ mPa·s) as working fluid. **(a)** Pressure head; **(b)** Shaft power

298 An optimal volumetric flow rate exists for each rotational speed. This optimum
 299 involves a trade-off between sufficiently high increment of the flow energy, proportional
 300 to the pressure head, and low power consumption, and it is determined based on the
 301 efficiency map presented in Fig. 8.

302 Once implanted, the pump will operate at a constant rotational speed adjusted by
 303 the physician, whose value must be determined to produce the desired CO for normal
 304 values of MAP. Nevertheless, since arterial pressure can change, the flow rate will
 305 vary as the rotational speed is not readjusted in real time. Therefore, the device will
 306 eventually work at off-design conditions.

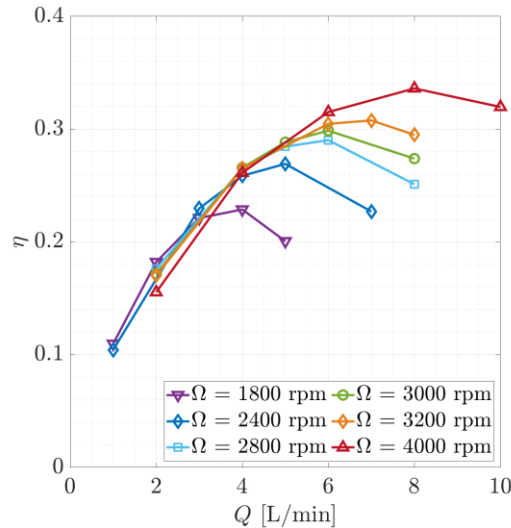


Fig. 8 Efficiency against volumetric flow rate for several values of rotational speed, for the configuration with 40 μ m-clearance gap, and blood with normal hematocrit ($\mu = 3.5$ mPa-s) as working fluid

307 To analyze flow patterns involving recirculating and stagnant flow, the velocity field
 308 must be represented in an appropriate reference frame, i.e. relative to moving walls in
 309 the rotating region and relative to static walls in the static region. Thus, the relative
 310 velocity in the rotating region is the result of subtracting the rotating motion of this
 311 region to the absolute velocity. The relative velocity field within the pump is presented
 312 in Fig. 9 (a,b) for two different operating conditions at $\Omega = 3000$ rpm: $Q = 5$ L/min
 313 (required CO for an adult) is a nominal condition near the optimal flow rate ($Q =$
 314 6 L/min), while $Q = 2$ L/min represents a point of excessively low flow rate.
 315 Recirculation zones are observed in the impeller channels for both operating
 316 conditions, indicating that an improved shape of the blades will lead to a more optimal
 317 design. An enlargement of those recirculation zones is detected for the low-flow case,
 318 as observed by other authors (Granegger et al. 2020; Thamsen et al. 2020; Schöps et
 319 al. 2021). As the flow rate increases, additional recirculation zones appear in the
 320 external wall of the impeller, at the adjacent blade of each channel. Note that the
 321 discontinuity in the velocity field at the interface between rotating and static regions is
 322 due to the change from moving to stationary reference frames (Karimi et al. 2021).

323 Flow patterns involving regions of recirculating and stagnant flow are, among other
 324 phenomena, responsible for thrombus formation within the device (Fraser et al. 2011).
 325 Zones with significant velocity gradients, as the impeller channels, are exposed to high
 326 shear stresses promoting platelet activation. Furthermore, within recirculating flow
 327 regions the blood's residence time is elevated promoting platelet deposition. Fig. 9 (c)
 328 shows an explanted HVAD where several thrombi have formed within the channels as
 329 well as around the impeller, where platelets were prone to deposit and form a clot due
 330 to flow recirculation seen in Fig. 9 (a,b). The photography in Fig. 9 (c) illustrates an
 331 advanced state of the thrombus formation reached after several months of operation,

332 whereas the relative velocity field observed in Fig. 9 (b) corresponds to the initial state
333 of the pump without clots. Note that, as activated platelets deposit in regions of low
334 velocity and low shear stress (recirculation regions) forming the clot, the blood flow
335 interacts with the growing thrombus leading to different fluid patterns. Therefore, the
336 clot formation would have initiated in the recirculation regions observed in Fig. 9 (b).
337 Then the clot, as it grows, would promote new recirculation regions where platelets
338 continue to deposit enlarging the clot until the final thrombotic state shown in Fig. 9 (c)
339 is reached. This HVAD was explanted from a patient who experienced complications
340 due to blood clotting inside the pump. The explant was done before the composition
341 of this article and the photography was taken just after the explant.

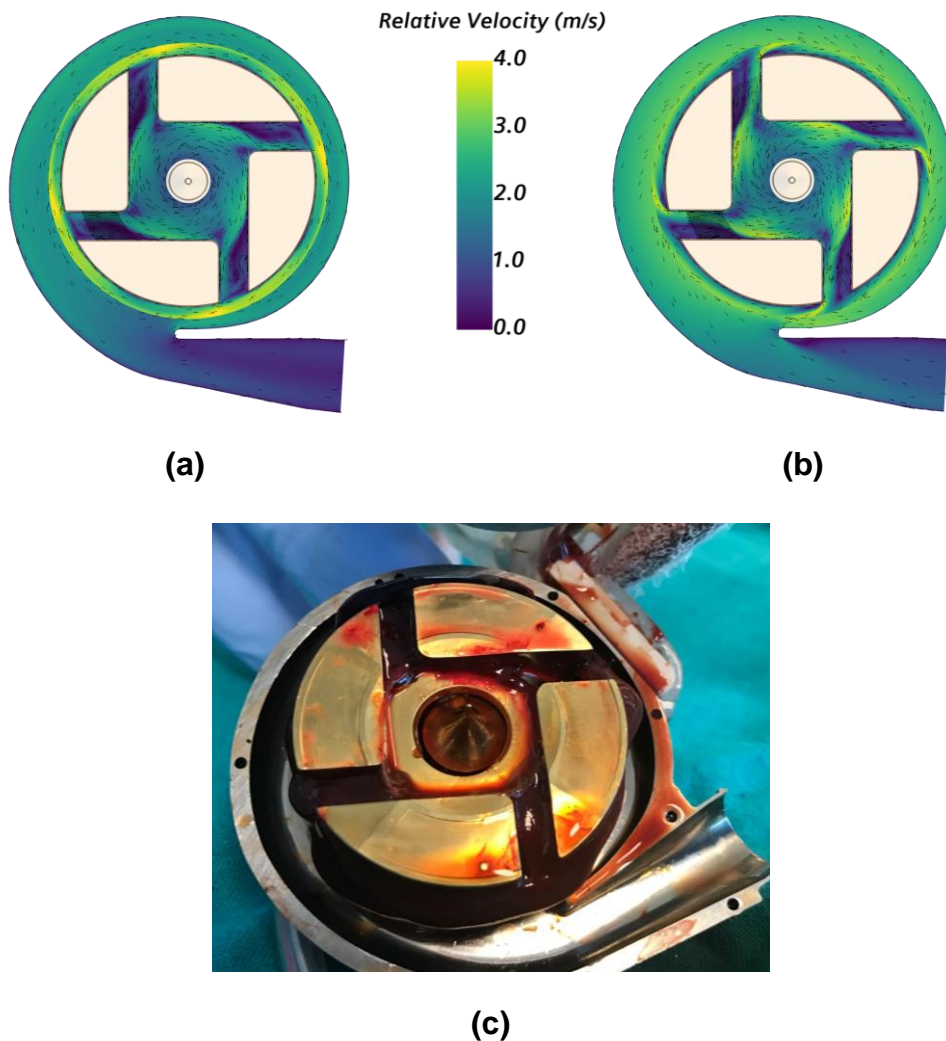


Fig. 9 Contours of the relative velocity field within the HVAD, for the configuration with 40 μ m-clearance gap operating at $\Omega = 3000$ rpm and for blood with normal hematocrit ($\mu = 3.5$ mPa·s) as working fluid: **(a)** Low-flow condition ($Q = 2$ L/min), **(b)** Nominal condition ($Q = 5$ L/min); **(c)** photography of an explanted HVAD showing the zones prone to thrombus formation, taken by the cardiac surgery team of Hospital La Fe

3.3 Gap clearance influence

343 One relevant design parameter of the HVAD is the gap clearance between impeller
344 and housing. The influence of this parameter is evaluated, as a novelty for this device,
345 since it is speed dependent and different gap clearances are used in the literature.
346 Fig. 10 depicts the performance curves at $\Omega = 3000$ rpm for the considered
347 configurations.

348 It could be expected that small gaps are preferable since they lead to higher
349 efficiencies because of the reduction in tip leakage (Fraser et al. 2011). Wiegmann et
350 al. (2018) investigated the effect of design parameters on efficiency and blood damage
351 for a centrifugal pump with conventional blades, and discussed the conflicting
352 requirements for gap clearance: while large gaps reduce the maximum shear stress
353 magnitude, small gaps induce less flow disturbances and provide higher efficiency.
354 Hence, small gaps are expected to be preferable to maximize the hydraulic efficiency
355 while large gaps lead to improved hemocompatibility. However, this device manifests
356 the opposite tendency for efficiency: it decreases as the gap clearance is reduced.
357 This decrease in efficiency can be justified by observing Fig. 10 (a,b): the pressure
358 head is not affected while the consumed power increases. The increase in power can
359 be explained based on Fig. 10 (d), where both pressure and shear power are
360 represented. The pressure and shear components of the power are related to the
361 torque exerted by the fluid on the impeller's walls, which consists of the moment due
362 to the static pressure on the surface and the moment resulting from the shear stresses
363 acting on it. While the shear power remains relatively constant, the magnitude of
364 pressure power increases with flow rate. The shear power experiences a significant
365 increase when halving the gap clearance since the secondary flow is generating higher
366 shear stresses in the gap regions. The effect of gap clearance on pressure power, for
367 its part, is negligible. Therefore, the shear component of the torque entails the main
368 contribution to the reduction of efficiency when reducing the gap clearance. Moreover,
369 the shear power is of the same order of magnitude than the pressure power, while it
370 is normally few orders of magnitude lower in conventional turbopumps.

371 The wide-blade impeller and the extremely narrow gaps of the HVAD are thought to
372 be originating this increase in shear torque, since these blades are notably thicker than
373 conventional blades used in turbomachinery and gap clearances are orders of
374 magnitude lower than those found in other pumps. While the efficiency decreases as
375 gap clearance increases due to tip leakage in conventional pumps, this device has
376 manifested the opposite effect. Potentially higher shear torque acts on the HVAD's
377 impeller owing to two design characteristics: the extremely low value of gap clearance
378 and the large blade tip areas involved in the gap region. This leads to reduced
379 efficiencies as the gap clearance decreases. Nonetheless, efficiency does not
380 increase monotonously with gap clearance, but will eventually decrease for larger
381 clearances owing to the increase in tip leakage that leads to a significant decrease in
382 pressure head. Nevertheless, larger values of gap clearance are not included in this

383 study because the HVAD's design does not allow for such large clearances due to the
 384 requirements for hydrodynamic lift.

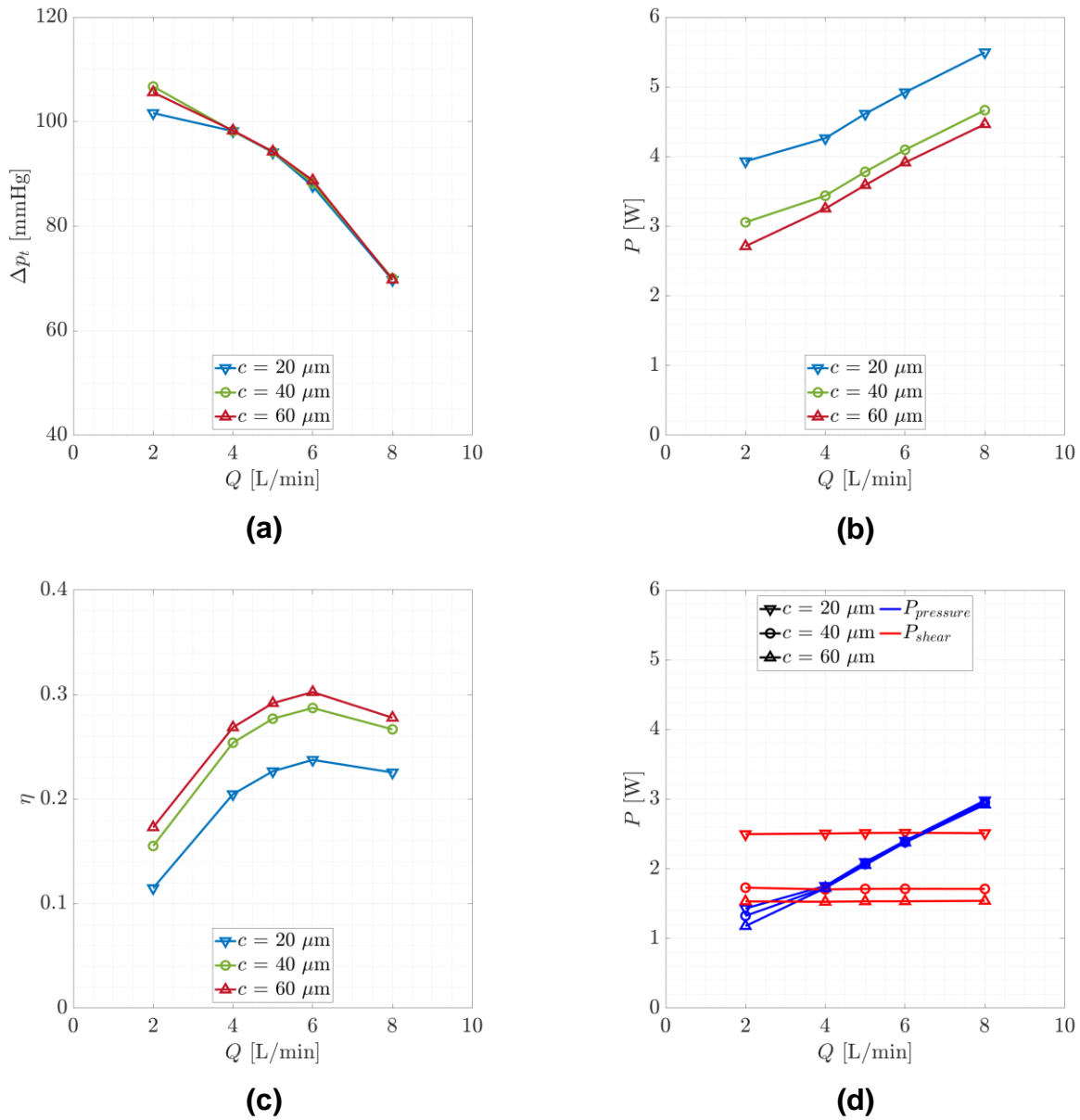


Fig. 10 Performance variables against volumetric flow rate, operating at $\Omega = 3000$ rpm, for blood with normal hematocrit ($\mu = 3.5$ mPa·s) as working fluid, and several values of gap clearance: **(a)** Pressure head; **(b)** Shaft power; **(c)** Efficiency; **(d)** Pressure and shear power

385 **3.4 Blood damage: hemolysis**

386 As described earlier, the hemolysis is investigated by means of the RHI defined in
 387 Equation 5. RHI against the volumetric flow rate is represented in Fig. 11 for several
 388 values of rotational speed (a) and gap clearance (b).

389 Fig. 11 (a) depicts that both increasing Ω and decreasing Q lead to an increase in
 390 HI, showing a similar trend as that obtained by Karimi et al. (2021). Increasing speeds
 391 produce more elevated shear stresses and, thus, higher hemolysis. Keeping constant
 392 the rotational speed, an asymptotic tendency is observed at high flow rates while an
 393 important increment of the risk of hemolysis is detected at low flow rates, as found by
 394 Thamsen et al. (2020), Granegger et al. (2020) and Schöps et al. (2021), due to longer
 395 washout times within the pump and larger zones of flow recirculation in the blade
 396 passages where the fluid is exposed to high velocity gradients and shear stresses.

397 According to the literature concerning other blood pumps, large gaps induce flow
 398 disturbances which increase shear stresses (Rezaienia et al. 2018), while narrow gaps
 399 induce cell screening in a way that, despite the higher maximum shear stress
 400 magnitude within the gap, less quantity of blood is exposed to those levels of shear
 401 stress since the secondary flow is reduced (Wiegmann et al. 2018). However, a global
 402 increase in the risk of hemolysis is detected in Fig. 11 (b) when reducing the gap
 403 clearance in the HVAD. Just as the decrease in efficiency, this increase in HI is a
 404 consequence of the increment of shear stresses over the large blade tip area involved
 405 in the gap region. Moreover, HVAD's gap clearances are extremely low, of the order
 406 of ten blood cells, each 6-8 μm in diameter, promoting elevated levels of hemolysis
 407 even for the larger gap considered in this work. For larger values of gap clearance, the
 408 HI will continue to decrease until the increment in secondary flow leads to a slight
 409 increase in hemolysis since more quantity of blood is exposed to elevated shear stress
 410 levels through gaps, but this increase is expected to be significantly lower than that
 411 observed for the values of gap clearance considered in this work and found in the
 412 HVAD.

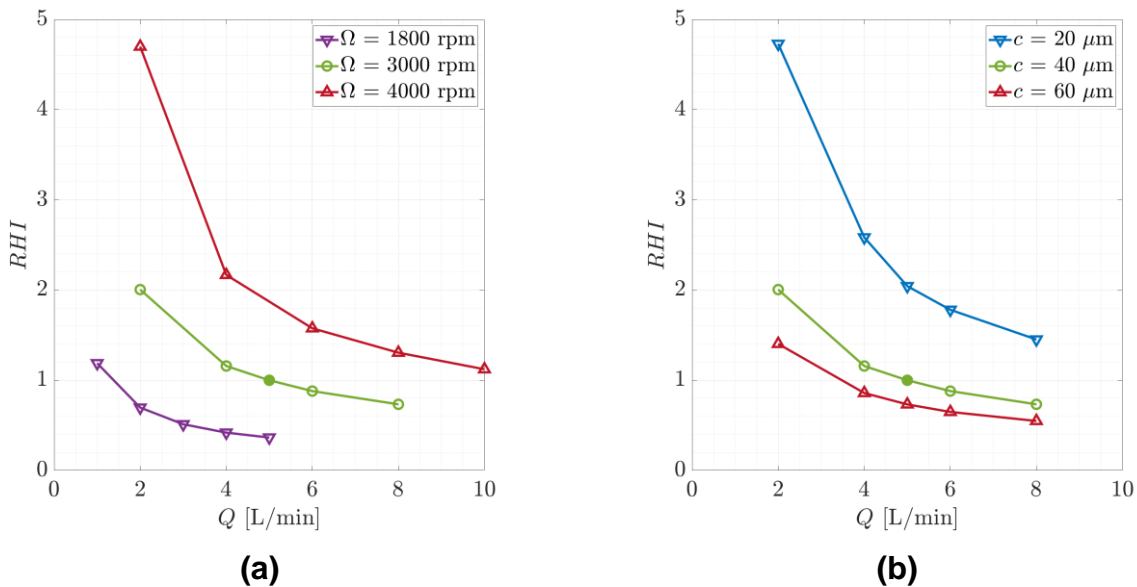


Fig. 11 Relative hemolysis index against volumetric flow rate for several values of rotational speed and gap clearance, and blood with normal hematocrit ($\mu = 3.5$ mPa·s) as working fluid: **(a)** Effect of rotational speed, for the configuration with 40 μm -clearance gap; **(b)** Effect of gap clearance, operating at $\Omega = 3000$ rpm

413 Fig. 12 depicts the HI field within the pump operating at $\Omega = 3000$ rpm at low-flow
 414 and nominal conditions. The former presents globally higher HI than the latter, as
 415 detected in the literature (Granegger et al. 2020; Thamsen et al. 2020; Schöps et al.
 416 2021) and discussed above. It can be seen how narrow gaps and recirculation zones
 417 induce a high degree of hemolysis, leading to increased HI downstream of the
 418 impeller.

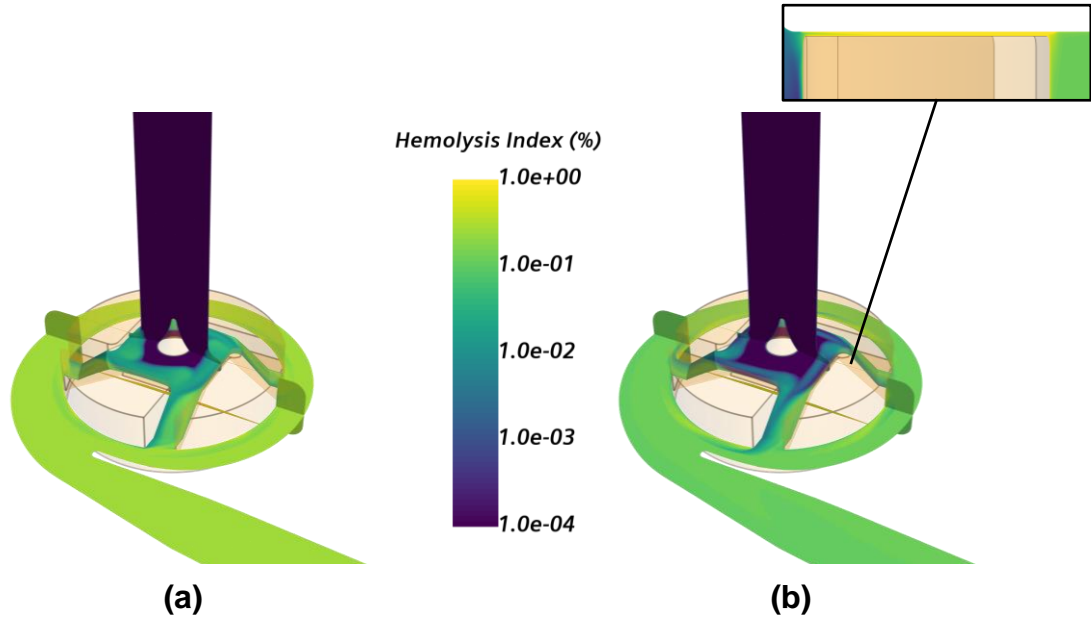


Fig. 12 Contours of the hemolysis index field within the HVAD, for the configuration with 40 μ m-clearance gap operating at $\Omega = 3000$ rpm, and blood with normal hematocrit ($\mu = 3.5$ mPa·s) as working fluid (gap zoom not to scale): **(a)** Low-flow condition ($Q = 2$ L/min); **(b)** Nominal condition ($Q = 5$ L/min)

419 3.5 Non-dimensional analysis

420 Previous figures can be extended to a non-dimensional form to ensure their
 421 applicability for multiple combinations of working conditions. Through a non-
 422 dimensional analysis it is observed that non-dimensional pressure head $\Delta p_t^* =$
 423 $\frac{\Delta p_t}{\rho \Omega^2 D_{imp}^2}$ and power $P^* = \frac{P}{\rho \Omega^3 D_{imp}^5}$ are a function of just the non-dimensional flow rate
 424 (flow number) $Q^* = \frac{Q}{\Omega D_{imp}^3}$ and the Reynolds number $Re = \frac{\rho \Omega D_{imp}^2}{\mu}$ (Heras 2011).

425 Note that increasing Reynolds numbers corresponds to increasing speeds or
 426 decreasing viscosities. The blood hematocrit indicates the quantity of RBCs present
 427 in the blood flow, which in turn affects the blood viscosity. The viscosity of $\mu = 3.5$ mPa·
 428 s corresponds to a hematocrit of 40% that is within the range of normal values of
 429 hematocrit for both men and women (Billett 1990). An increase in hematocrit produces
 430 an increment of the blood viscosity and, thus, a reduction of the Reynolds number.

431 Fig. 13 exhibits the non-dimensional operating maps, including two extreme
432 conditions:

- 433 ▪ Lowest Reynolds number: minimum speed ($\Omega = 1800$ rpm), extremely high
434 hematocrit of 60% ($\mu = 8$ mPa · s).
- 435 ▪ Highest Reynolds number: maximum speed ($\Omega = 4000$ rpm), extremely low
436 hematocrit of 20% ($\mu = 2$ mPa · s).

437 Whilst the effect of Reynolds number over the non-dimensional pressure head can
438 be assumed to be negligible for $Re \geq 10^5$, it is significant for the lowest Reynolds
439 numbers, as detected in Fig. 13 (a). This phenomenon is explained through the critical
440 Reynolds number marking the transition from laminar to turbulent flow, which is
441 conventionally taken to be of the order 10^5 for rotor-based Re (Heras 2011). Above
442 this value, Δp_t^* becomes Reynolds independent. Regarding the flow regime, Reynolds
443 numbers are relatively low for blood pumps as compared to other hydraulic pumps,
444 owing to their smaller size and the higher viscosity of the working fluid. Consequently,
445 the operating conditions typical of blood pumps place them in the transitional regime
446 between laminar and low Re turbulent regimes, where some degree of dependency
447 on Reynolds number still exists (Smith et al. 2004).

448 The main contribution of the non-dimensional analysis is related to the operating
449 map of P^* , shown in Fig. 13 (b). The HVAD's controller estimates the flow rate through
450 the pump based on the supplied power. To do this, it consults several maps of $P(Q, \Omega)$,
451 each corresponding to a different value of blood hematocrit. The non-dimensional map
452 of $P^*(Q^*, Re)$ allows this estimation, for every value of hematocrit, without needing
453 more than one map.

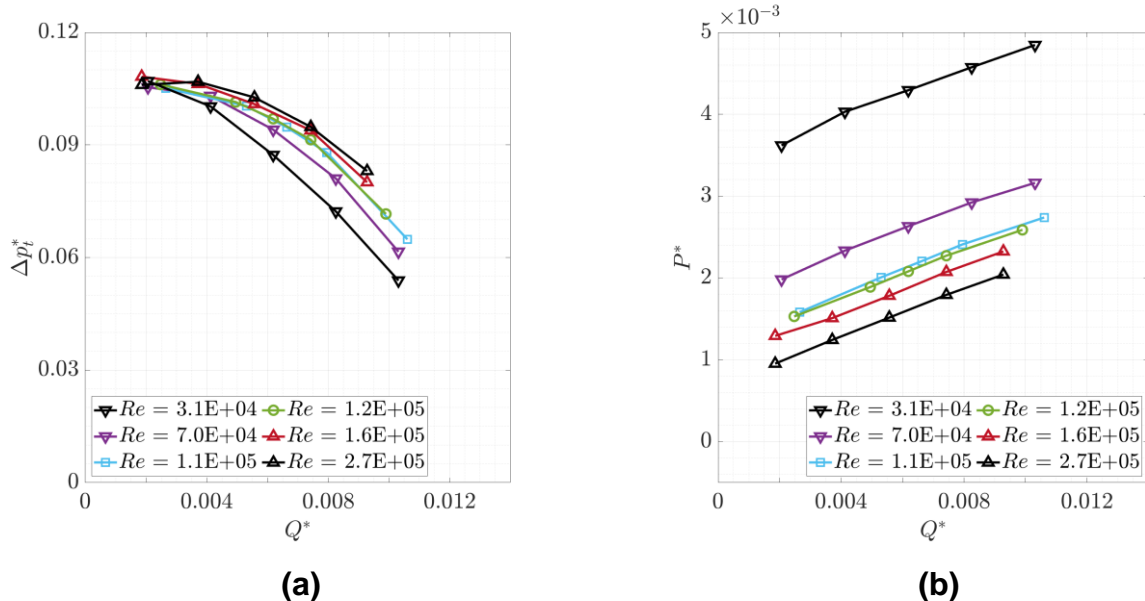


Fig. 13 Non-dimensional performance variables against non-dimensional flow rate for several values of Reynolds number, for the configuration with 40 μ m-clearance gap: **(a)** Pressure head; **(b)** Shaft power. The different values of Reynolds number correspond to normal blood hematocrit ($\mu = 3.5$ mPa·s) operating at four rotational speeds for the colored curves, and extremely low and high blood hematocrit ($\mu = 2$ mPa·s and $\mu = 8$ mPa·s) operating at the pump's maximum and minimum rotational speed, respectively, for the black curves

454 An important influence of Reynolds number is observed in Fig. 13 (b) for the non-
 455 dimensional power within the whole range, due to shear effects occurring in the gap
 456 region. Therefore, the different combinations of operating conditions do not collapse
 457 to the same efficiency curve in Fig. 14, where the optimal operating condition (Ω, Q)
 458 corresponds to $Q^* \approx 0.008$. A maximum pump efficiency of 10-40%, depending on
 459 Reynolds number, is obtained, which implies that most of the consumed energy is lost
 460 in form of vorticity and friction. This means that the mechanical energy supplied by the
 461 shaft is not just transformed into useful hydraulic energy provided to the flow, but also
 462 into lost energy consisting of friction losses due to the fluid's resistance to motion,
 463 related with the fluid's viscosity, and energy transferred to the recirculating flow and
 464 other turbulent features in form of vorticity. Friction losses explain the decrease in
 465 efficiency when augmenting the blood viscosity.

466 Although low viscosities would be preferable to maximize efficiency, low values of
 467 hematocrit involve complications regarding bleedings and insufficient clotting (Boneu
 468 and Fernandez 1987). Note, however, that this blood parameter depends on each
 469 patient and cannot be chosen.

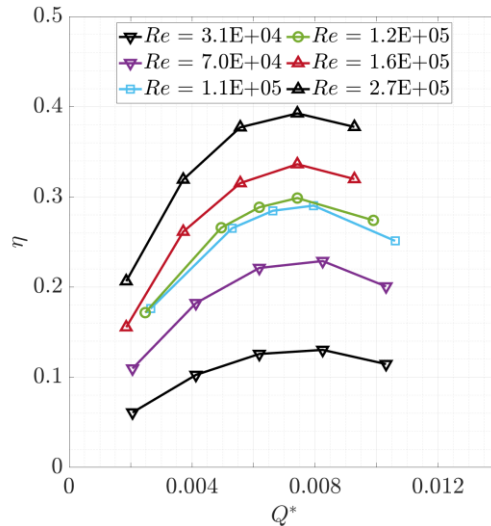


Fig. 14 Efficiency against non-dimensional flow rate for several values of Reynolds number, for the configuration with 40 μ m-clearance gap. The different values of Reynolds number correspond to normal blood hematocrit ($\mu = 3.5$ mPa·s) operating at four rotational speeds for the colored curves, and extremely low and high blood hematocrit ($\mu = 2$ mPa·s and $\mu = 8$ mPa·s) operating at the pump's maximum and minimum rotational speed, respectively, for the black curves

470 **4. CONCLUSIONS**

471 This work has presented the complete operating maps for the HVAD in terms of
 472 pressure head, power and efficiency. The pressure head map (H-Q curves)
 473 reproduces the typical tendency for centrifugal pumps. The power map depicts an
 474 almost linear dependency on flow rate. An efficiency of around 30% is obtained at
 475 nominal conditions.

476 By relating the hemocompatibility of the HVAD to its hemodynamic performance, it
 477 has been demonstrated that the optimal combination of operating conditions
 478 corresponds to that of reduced blood damage. The influence of flow conditions on
 479 hemolysis has been investigated in a different manner, by means of the RHI map,
 480 leading to the same conclusion than other authors: the low-flow condition induces
 481 potentially higher risk of hemolysis due to longer residence times and larger zones of
 482 recirculating flow (Granegger et al. 2020; Thamsen et al. 2020; Schöps et al. 2021).

483 The study regarding the influence of gap clearance in the HVAD has manifested
 484 tendencies different than those expected for conventional turbopumps. Shear stresses
 485 within the narrow gap region act over large blade tip areas. Thus, potentially high shear
 486 torque acts on the wide-blade impeller, leading to reduced efficiencies and raised
 487 hemolysis as the gap clearance decreases. Furthermore, gap clearances in HVAD
 488 take values that correspond to the size of less than ten blood cells and, thus, leading
 489 to significantly high risk of RBC damage. As a contribution of this work, these narrow

490 gaps, needed to produce hydrodynamic lift, are found to be significantly lower than the
491 optimum value that would lead to maximum efficiency and minimum hemolysis.
492 Hence, a fully magnetic levitation system would be preferable in comparison to the
493 hybrid (hydrodynamic and magnetic) system used in the HVAD, since it would allow
494 larger gap clearances.

495 In conclusion, the HVAD's complications that bring about its removal from the
496 market can be attributed to its narrow gaps and to the shape of its wide-blade impeller,
497 that lead to elevated shear stresses within gaps and flow recirculation zones in the
498 blade passages of the impeller respectively, promoting both hemolysis and
499 thrombosis. These observations will help in the design of a new device with enhanced
500 hemocompatibility and improved hemodynamic performance.

501 Furthermore, the non-dimensional representation of operating maps has
502 demonstrated that there is a significant influence of Reynolds number on performance
503 variables within the normal range of operating conditions of these devices, especially
504 on power and efficiency. This Reynolds dependency owes to the narrow gaps of the
505 device, that produce elevated shear stresses over a large area of the wide-blade
506 impeller. Moreover, the HVAD's controller estimates the flow rate based on the
507 consumed power using several operating maps for different values of blood
508 hematocrit. This estimation can be achieved using only one non-dimensional map
509 containing performance data in terms of flow rate and power as a function of Reynolds
510 number. Therefore, the non-dimensional representation of maps can imply a reduction
511 in the data stored for the flow rate estimation performed by the controller.

512 The computational approach employed in this work has limitations. On the one hand,
513 no experimental validation has been conducted for shaft power and hemolysis. On the
514 other hand, several limitations regarding the numerical set-up must be mentioned.
515 Firstly, the blood has been modeled as a Newtonian fluid, instead of considering its
516 variable viscosity, but this is a valid assumption since shear rates within the device are
517 greater than 100 s^{-1} in most of the domain. Secondly, most studies of this work have
518 been solved through steady-state simulations, dismissing transient effects.
519 Nevertheless, some operating conditions have been calculated in unsteady-state
520 simulations and acceptable differences have been found with respect to steady
521 results, proving the validity of the MRF approach. Note that an important reduction of
522 calculation time is achieved using this approach. Thirdly, the blood flow through the
523 extremely narrow gaps of this device does not behave as a continuum. Since the
524 hemolysis model employed in this work assumes that blood is a continuum, non-
525 continuum effects within gaps have not been considered. Finally, some phenomena
526 have not been evaluated, such as the effect of remaining native pulsatility and the
527 consideration of realistic boundary conditions that model the LV volume at the inlet or
528 the anastomosis to the aorta at the outlet.

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