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The shape influence on the overall single scattering properties of a sample in random orientation

R.Vilaplana^{a,*}, R.Luna^a, D.Guirado^b,

^a*Centro de Tecnologías Físicas, Universidad Politécnica de Valencia, Escuela Politécnica Superior de Alcoy, Alcoy, Alicante, Spain 03801*

^b*Instituto Astrofísico de Andalucía, Granada, Spain*

Abstract

The particle shape influences the curves of the scattering efficiency factor (Q_{sca}) as a function of the size parameter (X), and consequently on the overall single scattering properties of a sample of particles in random orientation. In order to show how the influence of the particle shape works, a model consisting of aggregates of different numbers of spheres has been used to fit laboratory measurements of fly ashes. The results for other shapes, such as rectangular prisms with different axial proportions, particles made of joined cubes, and particles with different fluffiness, are also shown. From all these calculations, it is concluded that the size averaged scattering matrix element resembles Rayleigh features, for the size distribution stopping at $1.0 \mu\text{m}$, when either the number of spheres or cubes of the aggregates is increased, the shape becomes flatter or the fluffiness degree is increased.

Keywords: Single scattering, Scattering efficiency, nonspherical particles

*Corresponding author

Email addresses: rovilap@fis.upv.es (R.Vilaplana), ralunam@fis.upv.es (R.Luna), dani@iaa.es (D.Guirado)

1. Introduction

The measurements carried out in the scattering laboratories [1] and the astronomical observations [2] provide with information on the overall scattering properties of a sample formed by small particles in random orientation. Single scattering properties of a distribution of particles in random orientation depend on properties of the grains, such as refractive index, size, shape, and the degree of fluffiness of the particles. Individually identifying or “untangling” the way in which each parameter is affecting the overall scattering properties is difficult since these parameters have a collective influence on the scattering pattern. In order to better understand how the single scattering properties observed are affected by these single parameters, a lot of research has been carried out [3-7]; however, some questions remain open. The goal of this study is to show how the particles shape influences the single scattering matrix elements of a sample of small particles in random orientation. In order to do this, we show the curves of Q_{sca} vs. X for different shapes and the overall single scattering properties of several samples of particles in random orientation. Although, our main goal is not to exactly fit the laboratory measurements, we present an attempt to fit a set of scattering laboratory measurements [1] by modelling it as a distribution of different aggregates of spheres. This comparison becomes a good example to point out how Q_{sca} vs. X , and consequently the overall single scattering properties, is affected by the particles shape. The results for other particle shapes are also showed to clarify the shape influence on the overall single scattering properties.

24 **2. Model**

25 The scattering matrix of a sample of particles of different sizes and the
26 same shape is calculated as in Eq. (1):

$$F_{ij}(\lambda, \theta) = \int_{r_1}^{r_2} F^{ij}(\lambda, \theta, r)n(r)dr \quad (1)$$

27 where $n(r)$ is the size distribution as a function of the radius, r_1 and r_2
28 correspond to the smallest and largest particles in the distribution respec-
29 tively, and $F^{ij}(\lambda, \theta, r)$ is one of the elements of the scattering matrix for a
30 single particle of radius r , at a certain wavelength λ for a scattering angle θ .

31 We have used the DDA (Discrete Dipole Approximation) for all calcu-
32 lations [8] because it has the potential to reproduce any particle shape, al-
33 though it is not suitable for performing calculations for particles much larger
34 than the wavelength of the incident light, on the current computers, in a rea-
35 sonable time (days). For all our calculations, the size distribution was chosen
36 a power law with an index of -1.8, and 35 equally spaced radii between 0.1
37 and 1.0 μm . The calculations were averaged over 2000 orientations, to mimic
38 the random orientation, and the number of dipoles was chosen so that the
39 accuracy condition $|\text{mkd}| < 0.5$ was fulfilled [9].

40 **3. Calculations for aggregate of spheres compared to measure-** 41 **ments**

42 The first shapes considered were aggregates of spheres. In this case, we
43 have compared our calculations with the single scattering laboratory mea-
44 surements of fly ashes [1]. The sample used in these measurements resembles

45 aggregates of spheres. In Fig. 1, we show an image of the eight aggregates
46 considered in our calculations. These aggregates are made of 5, 7, 7 in a line,
47 9, 14, 19, 25 and 36 spheres, so the four on the top are made of a smaller
48 number of spheres than those on the bottom. The value of the refractive
49 index was chosen $1.5+0.001i$ and the wavelength $0.633 \mu\text{m}$, as given in [1].

50 In Fig. 2, we show the comparison of the laboratory measurements with
51 our results, size averaged, for the eight aggregates of spheres showed in the
52 Fig. 1, plus a single sphere. In Fig. 2, we can see the overlapped images
53 of the size averaged results for each of the aggregates of spheres (blue and
54 red lines) and for a single sphere (dashed-dot-dot black line). It comes out
55 from Fig. 2 that the contribution of the aggregates of less number of spheres
56 (≤ 9) is necessary to approach the laboratory measurements (see in Fig. 2 the
57 blue and red lines). We can also notice a tendency to resemble the Rayleigh
58 features of the scattering matrix elements as functions of the scattering angle
59 when the aggregates are made of more spheres (see in Fig. 2 the dashed-dot
60 red lines). On the other hand, we note that the real size distribution, as
61 given in the reference [1], has constituents with X larger than 10. In Fig.
62 3, we show comparison of the laboratory measurements with our results size
63 and shape averaged considering the eight aggregates of spheres showed in
64 Fig. 1 plus a single sphere equally weighted. Without the aggregates with
65 less number of spheres (see blue lines in Fig. 2) the size and shape average
66 can not even approach the measurements.

67 In Fig. 3, we see a not perfect fit of the results of DDA to the measure-
68 ments, the calculations stopping at $r_2 = 1.0 \mu\text{m}$. In particular, the deviation
69 of the calculated values from the measurements points to a Rayleigh-like be-

70 haviour. From Fig. 2, we infer that the more spheres the aggregates are made
71 of, the more the calculated values resemble Rayleigh features of the scatter-
72 ing matrix elements as functions of the scattering angle. This is suggesting
73 us an explanation for the unperfected fitting: when aggregates are made of a
74 large number of spheres, the curve of Q_{sca} as a function of X changes so that
75 we are skipping some of its main features by cutting our size distribution at
76 $r_2 = 1.0 \mu\text{m}$. In order to prove this, we present on Fig. 4 the Q_{sca} curves
77 for the four aggregates with a number of spheres ≤ 9 till $X = 10$, along with
78 the Q_{sca} curve of the single sphere, calculated till $X = 15$. A progressive
79 displacement to the right and rising of the Q_{sca} curves is observed when the
80 number of spheres of the aggregates is increased. Due to this displacement,
81 some of the features of Q_{sca} that correspond to $r > 1.0 \mu\text{m}$ are lost in our cal-
82 culations, and this effect becomes more important as the number of spheres
83 of the aggregates increases. The result is a Rayleigh-like behaviour, because
84 only the first oscillation of the curve of Q_{sca} is been considered in the size
85 distribution.

86 4. Calculations for rectangular prisms

87 The next shapes considered were rectangular prisms with different axial
88 proportions namely 5:5:5, 5:4:4, 5:3:3, 5:2:2, 5:1:1, 5:4:1, 5:3:1, 5:2:1, 5:5:1,
89 5:5:2, 5:5:3, and 5:5:4. The values of the refractive index and the wavelength
90 for these calculations were $1.62+0.09i$ and $0.6 \mu\text{m}$, respectively. In Fig. 5,
91 we show four images, which show the overlapped curves of the Q_{sca} vs. X ,
92 obtained for different combinations of these prisms compared with the case of
93 a single sphere. Image a) of Fig. 5 shows the results for the prisms with the

94 extreme axial proportions 5:5:5, 5:1:1 and 5:5:1. A progressive displacement
95 to the right and rising of the Q_{sca} values compared with the curve of a
96 sphere is observed. The closest result to the sphere is for the cube and
97 the highest displacement and rising is for the prism with axial proportion
98 5:5:1. The other three images in Fig.5 show the results when varying the
99 axial proportions. Image b) of Fig. 5 shows the results for the transition
100 between the axial proportion from 5:1:1 to 5:5:5, image c) is the transition
101 from 5:5:5 to 5:5:1 and image d) from 5:5:1 to 5:1:1. All these images show
102 how is affected Q_{sca} vs. X depends on the axial proportion of the prisms.
103 The flattest shapes (5:5:1, 5:4:1 and 5:3:1) give the largest displacements and
104 risings.

105 In previous paper, we have compared the size-averaged scattering matrix
106 elements as functions of the scattering angle for some rectangular prisms with
107 axial proportions (see Fig.1 in reference [6]). The size-averaged scattering
108 matrix elements showed in this figure shows a Rayleigh-like behaviour for the
109 flattest rectangular prisms (5:5:1, 5:3:1 and 5:4:1). The reason is that only
110 the first oscillations of the Q_{sca} curves are considered in the size-average
111 of the scattering matrix elements (see the values of Q_{sca} curves marked with
112 arrows on the image a) and the values of Q_{sca} for the axial proportions 5:5:1,
113 5:3:1 and 5:4:1 in image c) of Fig. 5. A similar result with the flattest shapes
114 was obtained in other studies with platelike and needlelike particles [10] and
115 with spheroids of different axis ratios [11].

116 5. Calculations for aggregate of cubes

117 Other shapes considered were made by joined cubes. In Fig. 6, we show
118 an image with these shapes and the labels used. The shape called Test-h has
119 a hole in the center so then only six cubes forms it. The B9 and B3 have
120 the largest number of cubes, the B4 is the flattest one and the B1 and B7
121 differ only by a less cube in the bottom right corner of the B7. The values of
122 the refractive index and the wavelength were again $1.62+0.09i$ and $0.6 \mu\text{m}$,
123 respectively.

124 In Fig. 7, we show four plots with the overlapped curves of Q_{sca} vs. X for
125 different combinations of the shapes showed in Fig. 6, compared to the result
126 for a single sphere. Image a) of Fig. 7 show the Q_{sca} vs. X curves for all the
127 shapes in Fig. 6. The highest displacement and rise of the first maximum
128 of the Q_{sca} curves is observed for the particle B3 (one of the shapes made of
129 the largest number of cubes). The effect on the Q_{sca} curves of increasing the
130 number of cubes is showed more clearly in image b) of Fig. 7 for shapes B8,
131 B2, B6, B3 and B9. The flattest shapes B4 and B5 also show high values and
132 displacement of the first maximum of Q_{sca} although the number of cubes is
133 smaller than in shapes B1 and B7. This is showed in image c) of Fig. 7.
134 Finally, image d) of Fig. 7 shows the differences between the Q_{sca} values for
135 a single the cube, Test, Test-h and B1 particles. Surprisingly, the shape B9
136 which is made of 100 cubes gives a low value of the first maximum however
137 as we will show immediately after, this shape also is strongly influencing the
138 ripples over the Q_{sca} curves which seems to be totally different depending of
139 the shape.

140 In previous papers, we have compared the size-averaged scattering matrix

141 elements as functions of the scattering angle for some of these shapes with
 142 the same values of the refractive index and wavelength (see Fig. 5 in refer-
 143 ence [5] and Fig. 7 in reference [12]). The size-averaged scattering matrix
 144 elements showed in these figures seems to have values quite close to each
 145 other. However, in Fig. 8, we show the size-averaged scattering matrix ele-
 146 ments and the corresponding Q_{sca} vs. X curves for shapes B8, B3 and B9,
 147 which have an increasing number of cubes. This figure shows again that, the
 148 more cubes the particles are made of, the more the calculated values resem-
 149 ble Rayleigh features of the scattering matrix elements as functions of the
 150 scattering angle. This effect can be clearly observed in previous calculations
 151 with a shape of made of 256 cubes and an equal-size configuration but having
 152 spherical monomers instead of cubes [6]. In Fig.8 we have also marked with
 153 an arrows two pairs of points on Q_{sca} curves in which the high-frequency
 154 ripples observed are in a equivalent state. In other words, due to the influ-
 155 ence of the shape on the the low-frequency maxima and minima and on the
 156 high-frequency ripples some of the features of Q_{sca} curves that correspond
 157 to $r > 1.0 \mu\text{m}$ are lost, becoming this effect more important as the number
 158 of cubes of the aggregates increases. The result is a Rayleigh-like behaviour,
 159 because only the first oscillation of the curve of Q_{sca} is been considered in
 160 the size distribution.

161 6. Calculations for fluffy particles

162 The last particles considered have a fluffiness degree generated by uni-
 163 formly randomly removing dipoles in different percentages. The particle
 164 refractive index and the wavelength were again $1.62+0.09i$ and $0.6 \mu\text{m}$, re-

165 spectively. In Fig. 9, we show the Q_{sca} vs. X curves and the size-averaged
166 scattering matrix elements as functions of the scattering angle (excluding
167 F_{44}/F_{11}) for a compact sphere and this shape with fluffiness degrees of 15%,
168 25% and 50%. A displacement to the right and rising of the first maximum
169 of Q_{sca} is observed as the fluffiness degree increases; which implies that the
170 fluffier the sphere is, the more the scattering matrix elements as functions of
171 the scattering angle resemble Rayleigh features.

172 Our simple way to generate the porosity degree in the particles allow as
173 to obtain the effective refractive index by using the mixing rules of Effective
174 Medium Approximation [13], considering the inclusion refractive index as
175 $1.0+0.0i$. The result is that the real and imaginary part of the effective
176 refractive index decrease simultaneously and progressively as the fluffiness
177 degree is increased. On the other hand, in chapter 9 of reference [14] we can
178 see the effect of varying the real and imaginary part of the refractive index
179 on the scattering efficiencies curves for a sphere. A displacement to higher
180 values of X of all the maxima of the Q_{sca} curves is produced by a decrement
181 of the real part of the refractive index; however, a decrease of the imaginary
182 part of the refractive index produces a rise of Q_{sca} maxima without changing
183 their positions. Consequently, the displacement to the right and a rise of
184 the Q_{sca} values we observed in our calculations with fluffy sphere could be
185 understood in these terms.

186 To check what it happens with fluffy shapes different of a sphere, in Fig.
187 10 we show Q_{sca} vs. X curves and the corresponding size-averaged scattering
188 matrix elements as functions of the scattering angle (excluding the F_{44}/F_{11})
189 for shapes Test, Test-h and Test with different fluffiness degrees (14%, 25%

190 and 50%). The percentage of 14% corresponds to the quantity of matter
191 eliminated in the center of the shape Test-h. It is again observed that the
192 fluffier the shape is, the displacement and rise of the first maximum of Q_{sca}
193 become, and consequently, the scattering matrix elements as functions of the
194 scattering angle more resemble Rayleigh features.

195 7. Calculations for an equal size-configuration of cubes and spheres

196 Finally, we have carried out calculations with an equal size-configuration
197 of cubes and spheres for a very high refractive index. In Fig. 11, we show
198 the Q_{sca} vs. X curve and the corresponding scattering matrix elements as
199 functions of the scattering angle (excluding the F_{44}/F_{11}) for the shape B3 of
200 Fig. 6, which is made of 40 cubes and an equal-size configuration but having
201 spherical monomers instead of cubes; compared with a single sphere for a
202 refractive index of $2.0+0.4i$. Although the size distribution stops at $1.0 \mu\text{m}$,
203 a tendency to reach the geometric optic regime is observed for the two shapes
204 (equal-size configuration of cubes and spheres). In Fig. 11, the maximum
205 of the lineal polarization increases, does not show negative branch, and is
206 displaced to a scattering angle smaller than 90° .

207 8. Discussion and conclusions

208 We have carried out calculations with different shapes to shed light on
209 how the particle shape influences the single scattering matrix elements of a
210 sample of small particles in random orientation, concluding that the scat-
211 tering efficiency is an essential parameter to understand the influence of the
212 shape.

213 The Q_{sca} vs. X curve of a sphere is characterized by a succession of ma-
214 jor low-frequency maxima and minima with superimposed high-frequency
215 ripples. The low-frequency maxima and minima have been traditionally ex-
216 plained as the “interference structure” and the high-frequency ripples as a
217 consequence of the total reflection of the ray inside of the particle [15]. Tak-
218 ing into account these explanations is clear that not only the refractive index
219 of the particles will have a strong effect on the Q_{sca} vs. X curves but also the
220 shape. One of the main conclusions, we have reached with our calculations
221 is that fixed all the parameters of the model, the shape is influencing on
222 the Q_{sca} vs. X curve in such a way that the size-averaged scattering matrix
223 elements resemble Rayleigh features in three cases: a) the number of spheres
224 or cubes is increased in the particle, b) the flatter the particle is, c) the fluffi-
225 ness degree of the particles is increased. In other words, in the three cases
226 mentioned, a displacement to the right and rising of the first low-frequency
227 maximum is observed as a consequence of stopping the size distributions at
228 $r_2=1.0 \mu\text{m}$. Thus, it is easy to understand the Rayleigh-like behaviour ob-
229 served, because only the first oscillations or part of the Q_{sca} vs. X curves
230 is been considered in the size distribution, which corresponds to small size
231 parameters.

232 Our calculations also show how the high-frequency ripples are strongly
233 influenced by the shape. This effect is clearer observed for the refractive
234 index of $1.5+0.001i$ with the aggregates of spheres than for the refractive
235 index of $1.62+0.09i$ used for the rest of shapes. Changes in the absorption
236 or in the real part of the refractive index produce not only an effect on the
237 interference structure but also on the ripples structure. It takes great values

238 of absorption to eliminate the ripples and still greater (about 0.1) to eliminate
239 the interference structure. On the contrary, large values of the real part of
240 the refractive index increase the ripple structure becoming less pronounced
241 the interference structure (see chapter 9 in reference [14]). All these results
242 indicate that if the real and imaginary parts of the refractive index become
243 large at the same time, the effect on the scattering matrix elements for a fixed
244 shape will be the opposite of resembling the Rayleigh features; even stopping
245 the size distribution at $1.0 \mu\text{m}$. We have checked this with calculation using
246 the shape B3 and an equal-size configuration but having spherical monomers
247 instead of cubes and a value of the refractive index of $2.0+4.0i$. The optic
248 geometric regime is reached despite we cut the size distribution at $1.0 \mu\text{m}$.

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251 code available.

252 **References**

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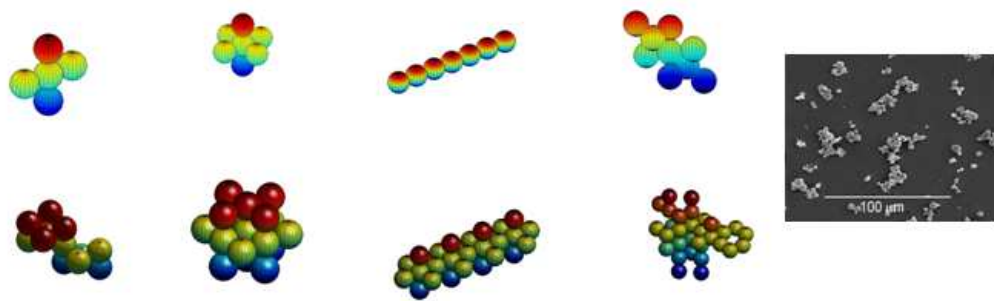


Figure 1: Eight aggregates made of 5, 7, 7 in a line, 9, 14, 19, 25 and 36 spheres.

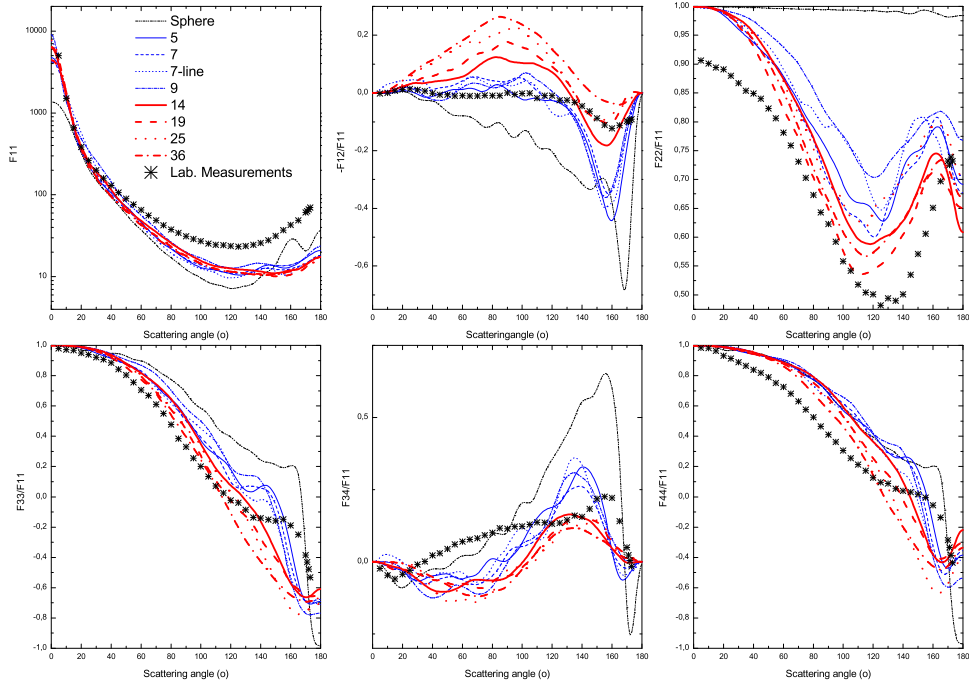


Figure 2: Comparison of laboratory measurements of the single scattering matrix elements of fly ashes with our size averages from 0.1 to 1.0 μm for each of the aggregates of spheres (blue and red lines) and for a single sphere (dashed-dot-dot black line). A refractive index of $m = 1.5 + 0.001i$ was used for the calculations along with a wavelength of 0.633 μm . The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and 1.0 μm .

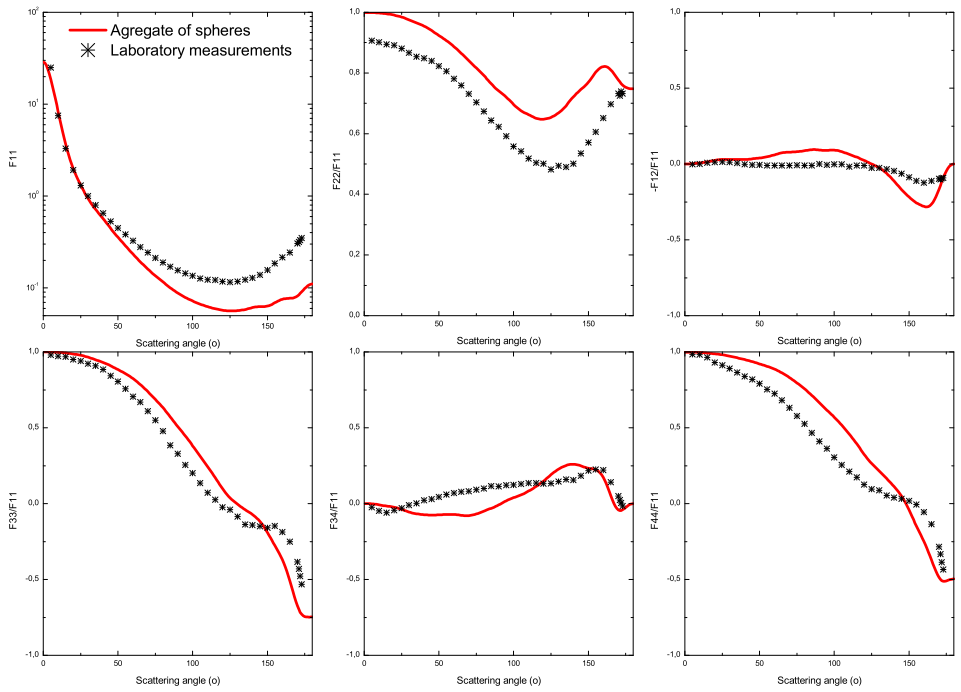


Figure 3: A comparison of laboratory measurements of fly ashes with our size and shape averages from 0.1 to $1.0 \mu\text{m}$, considering the eight aggregates of Fig. 1 and a single sphere, equally weighted.

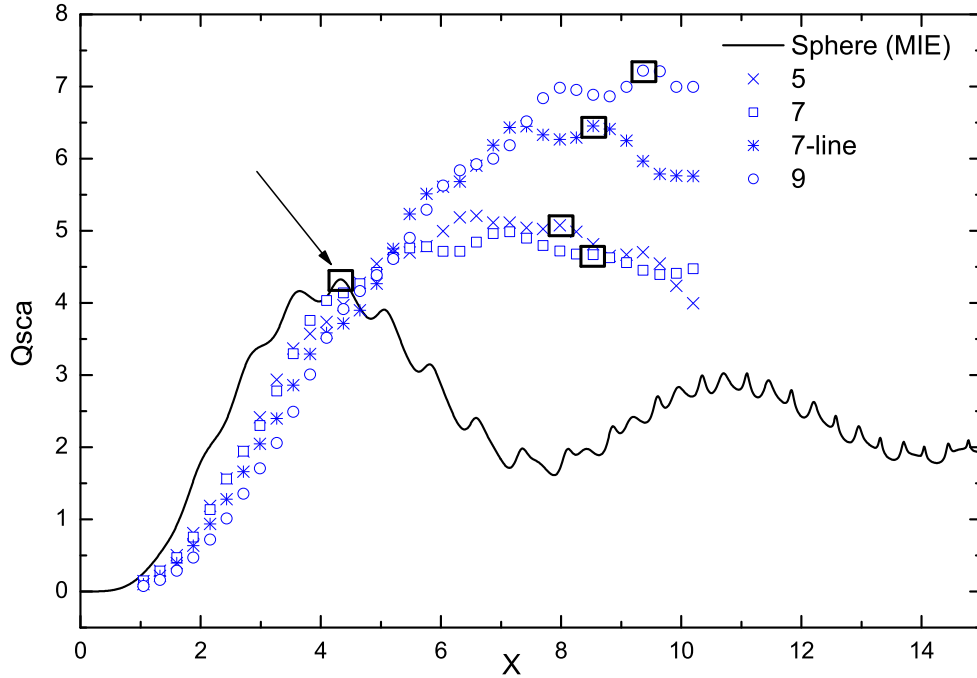


Figure 4: Q_{sca} versus X for the aggregates of Fig. 1 with a number of spheres ≤ 9 (5: \times , 7: \square , 7-line: $*$ and 9: \circ) and a single sphere (Mie: solid line). The squares on the Q_{sca} curves of the aggregates with 5, 7, 7-line and 9 spheres are considered in the "same" state of oscillation as that marked by an arrow on the Q_{sca} curve of the single sphere (solid line). A refractive index of $m = 1.5 + 0.001i$ was used for the calculations along with a wavelength of $0.633 \mu\text{m}$.

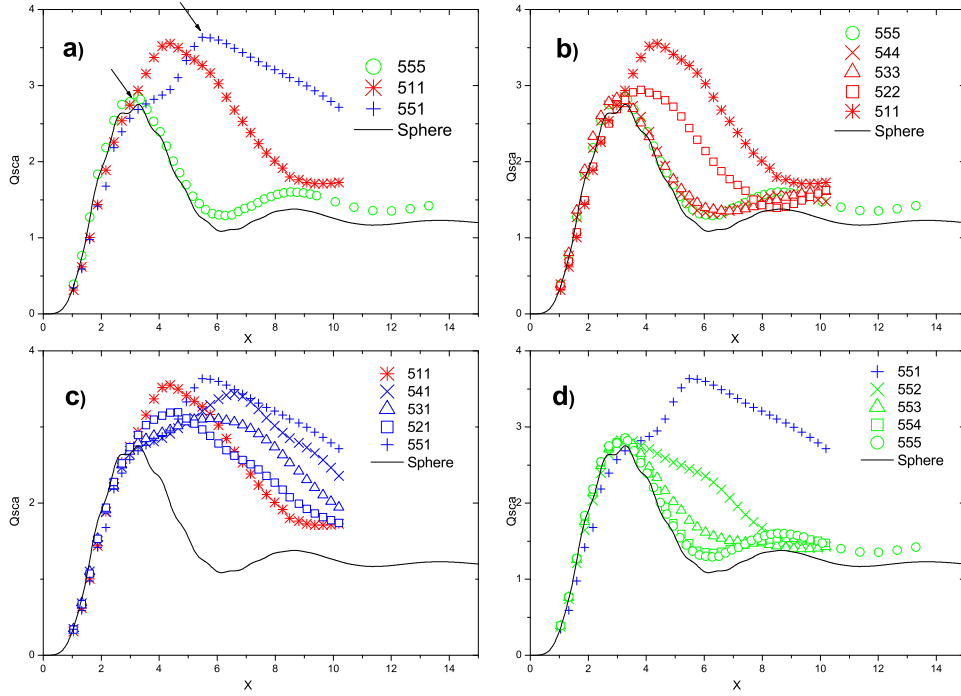


Figure 5: Q_{sca} vs. X curves for rectangular prisms with different axial proportions (5:5:5, 5:4:4, 5:3:3, 5:2:2, 5:1:1, 5:4:1, 5:3:1, 5:2:1, 5:5:1, 5:5:2, 5:5:3 and 5:5:4) compared with the results for a sphere. A refractive index of $m = 1.62+0.09i$ was used for the calculations along with a wavelength of $0.6 \mu\text{m}$. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.

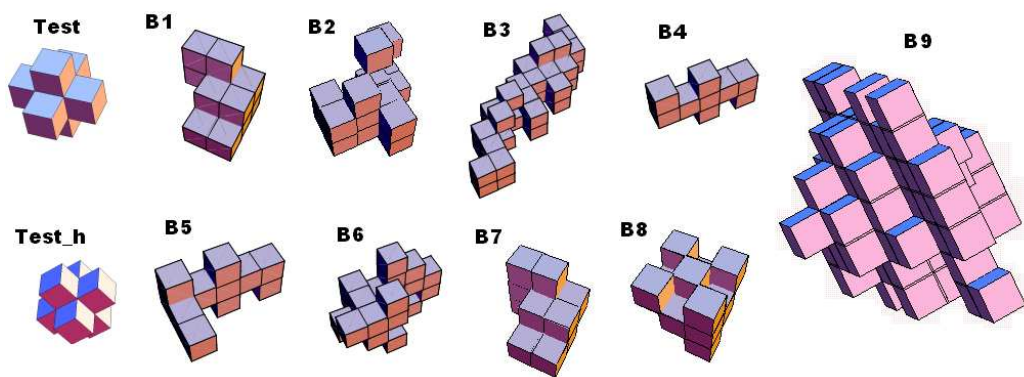


Figure 6: Ten shapes formed by joined cubes: Test, Test_h, B1, B2, B3, B4, B5, B6, B7, B8 and B9 and the labels used to refer them.

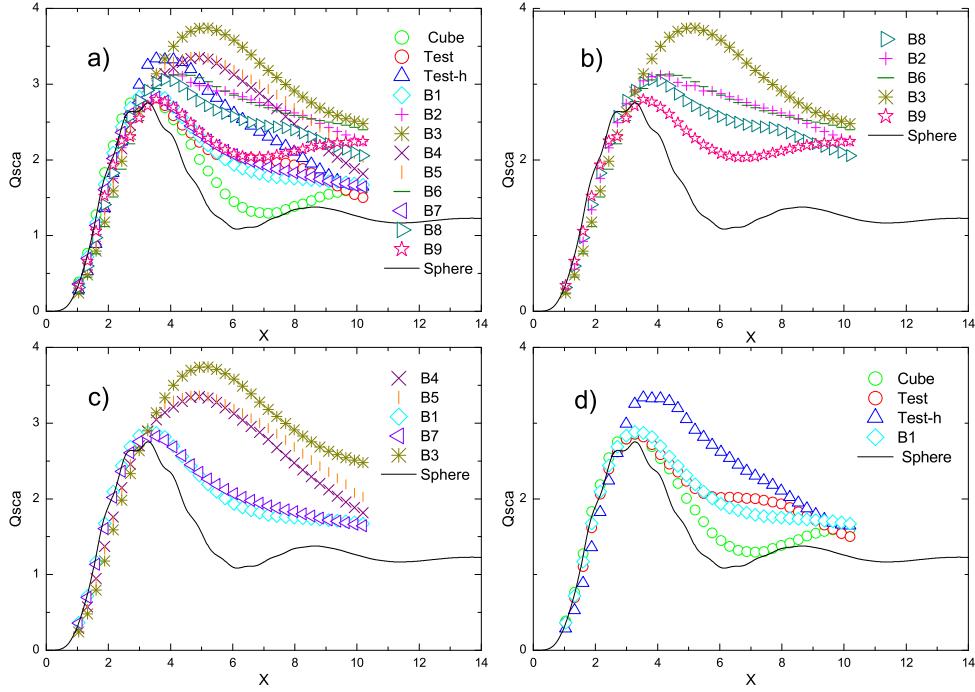


Figure 7: The overlapped curves of the Q_{sca} vs. X for different combinations of the shapes in Fig. 6: a) All shapes plus a cube; b) B8, B2, B6, B3 and B9; c) B4, B5, B1, B7 and B3; d) Test, Test-h, B1 plus a cube; and a single sphere. A refractive index of $m = 1.62 + 0.09i$ was used for the calculations along with a wavelength of $0.6 \mu\text{m}$. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.

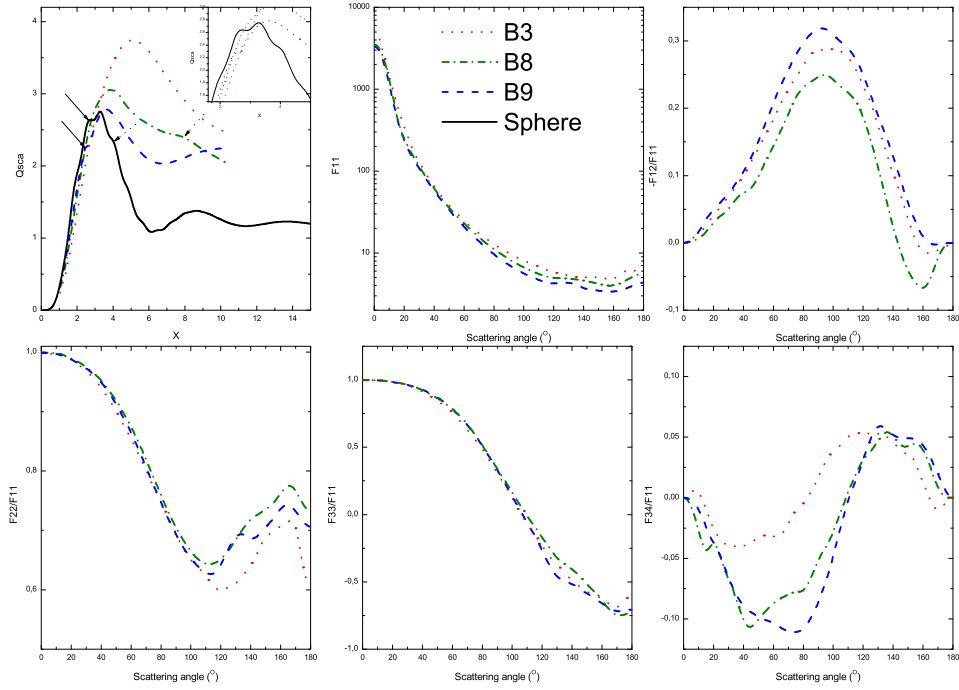


Figure 8: Q_{sca} vs. X curves and corresponding size-averaged scattering matrix elements as functions of the scattering angle (excluding F_{44}/F_{11}) for shapes B3, B8, and B9. A refractive index of $m = 1.62 + 0.09i$ was used for the calculations along with a wavelength of $0.6 \mu\text{m}$. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.

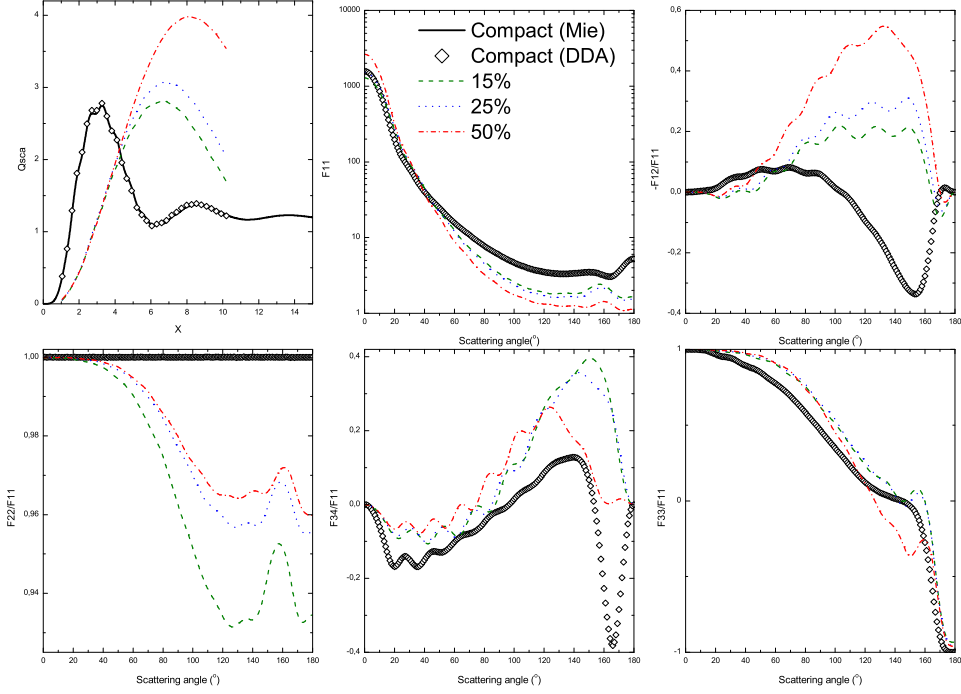


Figure 9: Q_{sca} vs. X and the corresponding size-averaged scattering matrix elements (excluding the F_{44}/F_{11}) as functions of the scattering angle for a single sphere and the sphere with the fluffiness degrees of 15%, 25% and 50%. A refractive index of $m = 1.62 + 0.09i$ was used for the calculations along with a wavelength of $0.6 \mu\text{m}$. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.

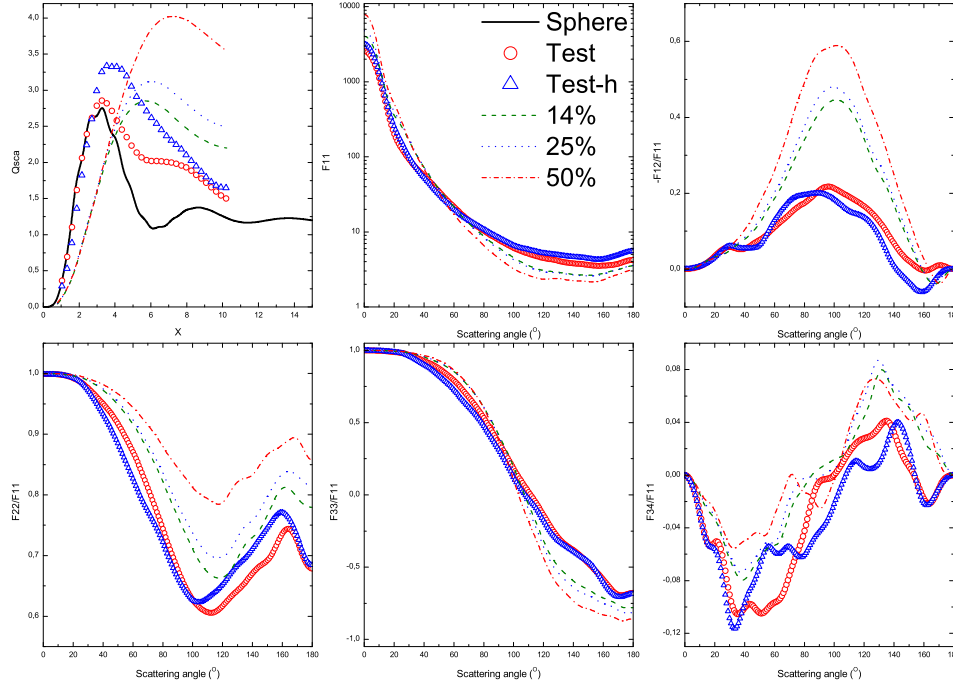


Figure 10: Q_{sca} vs. X and the corresponding size-averaged scattering matrix elements (excluding the F_{44}/F_{11}) as functions of the scattering angle for shapes Test, Test-h and Test with the fluffiness degrees of 14%, 25% and 50%. A refractive index of $m = 1.62 + 0.09i$ was used for the calculations along with a wavelength of $0.6 \mu\text{m}$. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.

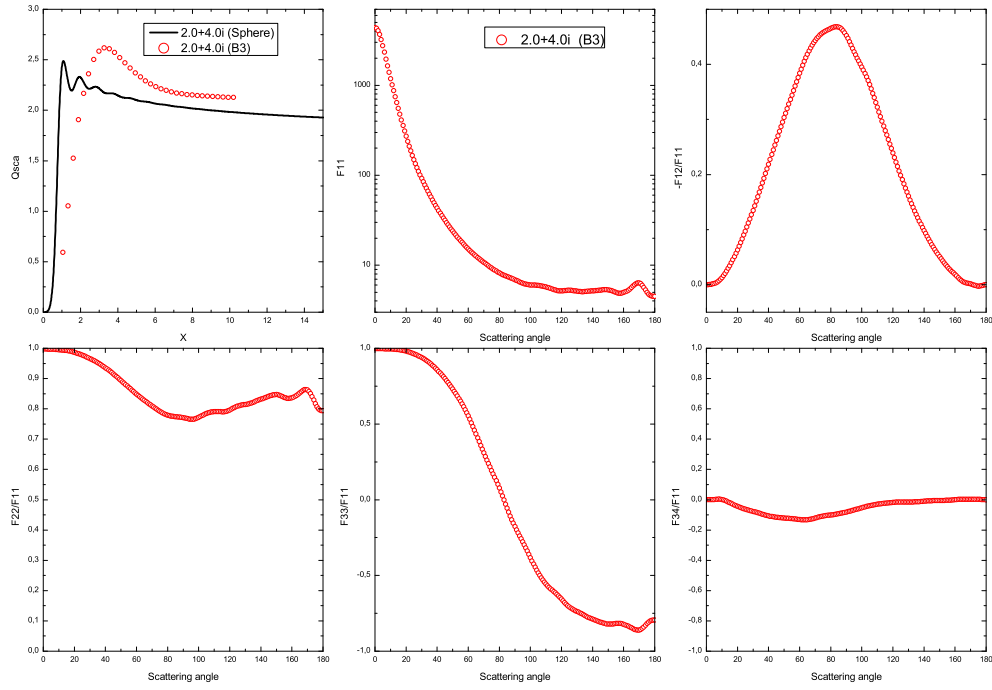


Figure 11: Q_{sca} vs. X and the corresponding size-averaged scattering matrix elements (excluding the F_{44}/F_{11}) as functions of the scattering angle for shape B3. A refractive index of $m = 2.0+4.0i$ was used for the calculations. The elements of the scattering matrix were size-averaged over a power law distribution with an index of -1.8 and 35 equally spaced radii between 0.1 and $1.0 \mu\text{m}$.