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Comparison of SERPENT and SCALE methodology for LWRs transport calculations and additionally uncertainty analysis for cross-section perturbation with SAMPLER module

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Abstract. In nuclear safety research, the quality of the results of simulation codes is widely determined by the reactor design and safe operation, and the description of neutron transport in the reactor core is a feature of particular importance. Moreover, for the long effort that is made, there remain uncertainties in simulation results due to the neutronic data and input specification that need a huge effort to be eliminated. A realistic estimation of these uncertainties is required for finding out the reliability of the results. This explains the increasing demand in recent years for calculations in the nuclear fields with best-estimate codes that proved confidence bounds of simulation results. All this has lead to the Benchmark for Uncertainty Analysis in Modelling (UAM) for Design, Operation and Safety Analysis of LWRs of the NEA. The UAM-Benchmark coupling multi-physics and multiscale analysis using as a basis complete sets of input specifications of boiling water reactors (BWR) and pressurized water reactors (PWR). In this study, the results of the transport calculations carried out using the SCALE-6.2 program (TRITON/NEWT and TRITON/KENO modules) as well as Monte Carlo SERPENT code, are presented. Additionally, they have been made uncertainties calculation for a PWR 15×15 and a BWR 7×7 fuel elements, in two different configurations (with and without control rod), and two different states, Hot Full Power (HFP) and Hot Zero Power (HZP), using the TSUNAMI module, which uses the Generalized Perturbation Theory (GPT), and SAMPLER, which uses stochastic sampling techniques for cross-sections perturbations. The results obtained and validated are compared with references results and similar studies presented in the exercise I-2 (Lattice Physics) of UAM-Benchmark.

1 Introduction

This work takes part in the framework of the Organization for Economic Cooperation and Development/Nuclear Energy Agency (OECD/NEA) Benchmark, for Uncertainty Analysis in Best-Estimate Modelling (UAM), for Design, Operation and Safety Analysis of LWRs. The groundwork was launched in 2005 with the objective to prepare a benchmark work program with steps (exercises) that would be needed to define the uncertainty and modelling task for the development of uncertainty analysis methodologies, for multi-physics and multi-scale simulation. The final goal will create a roadmap along with schedule and organization, for the development and validation of method and codes required for uncertainty and safety analysis in LWR design [1].

Reference systems and scenarios for coupled code analysis are defined to study the uncertainty effects for all stages of the system calculations. Measured data from plant operation and experimental reference data are available for the chosen scenarios. The full chain of uncertainty propagation from basic data, engineering uncertainties, across different scales (multi-scale), and physics phenomena (multi-physics) is tested on some benchmark exercises for which experimental data are available and for which the power plant details have been released. The general frame of the OECD UAM LWR Benchmark consists of three phases with different exercises for each phase: Phase 1 (neutronics phase), Phase 2 (core phase) and Phase 3 (system phase). The focus of Phase 1 is on propagating uncertainties in stand-alone neutronics calculations and consists of the following three exercises:

 Exercise I-1: "Cell Physic" focused on the derivation of the multigroup microscopic cross-section libraries and associated uncertainties;

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- Exercise I-2: "Lattice Physics" focused on the derivation of the few-group macroscopic cross-section libraries and associated uncertainties;
- Exercise I-3: "Core Physics" focused on the core steady state stand-alone neutronics calculations and associated uncertainties.

The present paper deals with Cell Physic and Lattice Physics exercises, determining uncertainties associated with basic nuclear data, method and modelling approximation used in lattice physics codes.

This document is structured as follows: the Introduction in Section 1. Section 2 is focused on the description of the codes. The description of the models is shown in Section 3. The results for different codes are presented in Section 4. Finally, the conclusions are shown in Section 5.

2 Codes description

2.1 Transport calculation

In this work, two- and three-dimensional lattice codes (deterministic and stochastic) were selected to perform transport calculations: SCALE-6.2 with TRITON/NEWT and TRITON/KENO modules and SERPENT-2.1.22 code.

The SCALE code system [2] is a collection of computational modules whose execution can be linked by various "sequences" to solve a wide variety of applications.

TRITON (Transport Rigor Implemented with Timedependent Operation for Neutronics depletion) is a multipurpose SCALE control module for transport and depletion calculations for reactor physics applications. TRITON is used to provide automated, problem-dependent cross-sections processing followed by multigroup neutron transport calculation for one-, two- or threedimensional configuration [3].

NEWT (New ESC-based Weighting Transport code) is a two-dimensional (2D) discrete ordinates transport code developed at Oak Ridge National Laboratory. It is based on the Extended Step Characteristic (ESC) approach for spatial discretization in an arbitrary mesh structure. This discretization scheme makes NEWT an extremely powerful and versatile tool for deterministic calculation in real-world non-orthogonal problem domains. The NEWT computer code has been developed to run on SCALE. Thus, NEWT uses AMPX-formatted cross-sections processed by other SCALE modules [4].

The implemented methodology of these coupled modules of SCALE allows carrying out transport calculation with the computation of energy collapsed and homogenized macroscopic cross-sections. In TRITON, NEWT is used to calculate weighted burnup-dependent cross-sections that are employed to provide localized fluxes used for multiple depletion regions. Additionally, TRITON uses a two-pass cross-section update approach to perform fuel assembly burnup calculations and generates a database of cross-sections and other burnup-dependent physics data that can be used for full-core analysis [5]. KENO is a functional module in the SCALE system and a Monte Carlo criticality program used to calculate the k_{eff} of three-dimensional (3D) system [6]. It uses the SCALE Generalized Geometry Package (SGGP), which offers a powerful geometric representation. KENO was one of the oldest criticality safety analysis tools in SCALE. The primary purpose of its employment in this work is to determine k_{eff} calculations and compare KENO results with TRITON/NEWT and SERPENT-2 calculation.

SERPENT is a three-dimensional continuous-energy code, based on the Monte Carlo method, for reactor physics burnup calculation [7,8]. The SERPENT project started in 2004 at the VTT Technical Research Centre of Finland. The first version of the code was available to universities and research institutes from 2008 and currently it is still under development. The suggested applications of SER-PENT include, among other applications, the spatial homogenization and constant group generation for deterministic reactor calculations and the validation of deterministic lattice transport codes.

2.2 Sensitivity and uncertainty calculation

Sensitivity analysis and propagation of uncertainties of cross-sections have been carried out using TSUNAMI and SAMPLER modules.

TSUNAMI-2D (Tools for Sensitivity and Uncertainty Analysis Methodology Implementation in Two Dimension) is a SCALE sequence for calculating sensitivity coefficients and response uncertainties to nuclear systems analyses for criticality safety applications. TSUNAMI uses the Generalized Perturbation Theory (GTP) that performs similarity analysis and consolidates experimental and computational results through data adjustment. The uncertainties, resulting from uncertainties in the basic data, are estimated using energy-dependent cross-section-covariance matrices [9]. The SAMS module is used to determine the sensitivities of calculated value of k_{eff} and other system responses to the nuclear data used in the calculations as a function of a nuclide, reaction type, and energy. This sensitivity of k_{eff} to the number density is equivalent to the sensitivity of k_{eff} to the total cross-section, integrated over energy. Because the total cross-section sensitivity coefficient tests much of the data used to compute all other sensitivity coefficients, it is considered an adequate test for verification. For each sensitivity coefficient examined by direct perturbation, the k_{eff} of the system is computed first with the nominal values of the input quantities, and then with a selected nominal input value increased by a certain percentage, and then with the nominal value decreased by the same percentage. The direct perturbation sensitivity coefficient of k_{eff} to some input value α is computed as:

$$S_{k,\alpha} = \frac{\alpha}{k} \times \frac{dk}{d\alpha} = \frac{\alpha}{k} \times \frac{k_{\alpha^+} - k_{\alpha^-}}{\alpha^+ - \alpha^-}, \qquad (1)$$

where α^+ and α^- represent the increased and decreased values, respectively, of the input quantity α , and $k_{\alpha+}$ and $k_{\alpha-}$ represent the corresponding values of k_{eff} . are propagated to direct perturbation sensitivity coefficients by standard error propagation techniques as:

$$\sigma_{S} = \sqrt{\left(\frac{\left(\sigma_{k^{+}}^{2} + \sigma_{k^{-}}^{2}\right)}{\left(k^{+} - k^{-}\right)} + \frac{\sigma_{k}^{2}}{k^{2}}\right)} \times \left(\frac{k^{+} - k^{-}}{k}\right)^{2} \times \frac{\alpha}{\alpha^{+} - \alpha^{-}}.$$
(2)

It is important in sensitivity calculations to ensure that the k_{eff} value of the forward and adjoint solutions closely agree, and typically, the transport calculation of concern is the adjoint calculation. More details of the GPT methodology are provided in the SAMS manual [10].

SAMPLER is a module for statistical uncertainty analysis of any SCALE sequences. The SAMPLER methodology samples probability density functions (pdf) defined by information in the SCALE multigroup covariance library by XUSA program and produces a random sample for the input computational data vector (CDV) that contains all nuclear cross-sections used in a transport calculation. After making random perturbations in input data, SAMPLER responses uncertainties are computed by statistical analysis of output responses distribution [11].

The perturbed data vector can be used in any SCALE functional module to perform a single forward solution that computes all the desired perturbed responses. The process is repeated for the specified number of random input samples, and the resulting distribution of output responses from SCALE can be analysed with standard statistical analysis tools to obtain the standard deviations and correlation coefficients for all responses. The typical approach is to assume that the generic multigroup (MG) data pdf is a multivariate normal distribution, which is completely defined by the expected values and covariance matrices for the data. An XSUSA statistical sample consists of a full set of perturbed, infinitely dilute MG data for all groups, reactions, and materials. The SCALE generic multigroup covariance data are given as relative values of the infinitely dilute cross-sections, so a random perturbation sample for cross-sections $\sigma_{x,q}(\infty)$ corresponds to $\Delta \sigma_{x,q}(\infty) / \sigma_{x,q}(\infty)$. XSUSA converts these values to a set of multiplicative perturbation factors Q_x . $_{q}$ that are applied to the reference data to obtain the altered values:

where

$$Q_{x,g} = 1 + \frac{\Delta \sigma_{x,g}(\mathbf{\infty})}{\sigma_{x,g}(\mathbf{\infty})}.$$
(4)

(3)

Subsequently, the multiplicative perturbation factors for all data are pre-processed and stored in a data file for subsequent SCALE calculations [12].

 $\sigma'_{x,q} = Q_{x,q}\sigma_{x,g},$

To obtain the uncertainty and correlation coefficient, all parameters are randomly perturbed for each calculation, and the uncertainties and correlations are determined.

Statistical uncertainties in the computed values of k_{eff} Mathematically, the uncertainty in an individual output parameter k is determined as:

$$\Delta k^{\exp}(i) = \hat{\mu}_i = \sqrt{\frac{1}{n-1} \sum_{a=1}^n \left(\left(k^{MC}_{calc}(i) \right)_a - \overline{k^{MC}_{calc}(i)} \right)^2}, \quad (5)$$

where $\Delta k^{\exp}(i)$ is the uncertainty in system *i* due to uncertainties in the input parameters.

 $(k_{calc}^{MC}(i))_a$ is the *a*th Monte Carlo (MC) sample of system *i*, where all uncertain input parameters have been randomly varied within the specified distribution.

The covariance between two systems, i and j, is determined as shown in equation (6).

$$\hat{\mathcal{L}}_{ij} = \sqrt{\frac{1}{n-1} \sum_{a=1}^{n} \left(\left(k_{calc}^{MC}(i) \right)_{a} - \overline{k_{calc}^{MC}(i)} \right) \left(\left(k_{calc}^{MC}(j) \right)_{a} - \overline{k_{calc}^{MC}(j)} \right)}.$$
(6)

The correlation coefficient between systems i and j can be determined from equations (5) and (6) as:

$$c_{ij} = \frac{\Sigma_{ij}}{\hat{\mu}_i \hat{\mu}_j}.$$
(7)

3 Model description

Two main LWR types have been selected for this study, based on previous benchmark experience and available data:

- Pressurized Water Reactor (PWR) Three Mile Island 1 (TMI-1);
- Boiling Water Reactor (BWR) Peach Bottom 2 (PB-2).

Both models have been analyzed at Hot Full Power (HFP) and Hot Zero Power (HZP) conditions.

Additionally, the two models have been designed with and without control rod.

The different fuel pin cell geometry and reference configuration are schematized in Figures 1 and 2.

SCALE calculations use the Extended Step Characteristic (ESC) approach. The entire problem domain is mapped regarding a set of finite cells. Cells sharing a given side share the value of the angular flux on that side. Once the angular flux has been determined for all sides of the cell for the given direction, it is possible to use a neutron balance to compute the average value of the angular flux within the cell. The process is then repeated for all direction. Numerical quadrature can then be used to determine the average scalar flux in each cell in the problem domain and can be used to determine fission and scattering reaction rates and to update the value of average cell source. In this way, the spatial discretization in SCALE allows to obtain satisfactory results.

The propagation of the cross-sections uncertainties across lattice physics is the main purpose of this exercise. To achieve that, in UAM-Benchmark instructions there are defined two assembly design for the models studied (a PWR

month month		value
	Unit cell pitch, [mm]	18.75
187 187	Fuel pellet diameter, [mm]	12.1158
	Fuel pellet material	UO
fuel pin	Fuel density, [g/cm ³]	10.42
	Fuel enrichment, [w/o]	2.93
	Cladding outside diameter, [mm]	14.3002
	Cladding thickness, [mm]	0.9398
	Cladding material	Zircaloy-2
\sim	Cladding density, [g/cm ³]	6.55
) – pitch of the unit cell	Gap material	He
- pitch of the unit cell	Moderator material	H ₂ O
Parameter / Reactor condition	HZP HFP	
Parameter / Reactor condition Fuel temperature, [K]	HZP HFP 552.833 900	
Parameter / Reactor condition Fuel temperature, [K] Cladding temperature, [K]	HZP HFP 552.833 900 552.833 600	
Parameter / Reactor condition Fuel temperature, [K] Cladding temperature, [K] Moderator (coolant) temperature, [K]	HZP HFP 552.833 900 552.833 600 552.833 557	
Parameter / Reactor condition Fuel temperature, [K] Cladding temperature, [K] Moderator (coolant) temperature, [K] Moderator (coolant) density, [kg/m ³]	HZP HFP 552.833 900 552.833 600 552.833 557 753.978 460.72	
Parameter / Reactor condition Fuel temperature, [K] Cladding temperature, [K] Moderator (coolant) temperature, [K] Moderator (coolant) density, [kg/m ³] Reactor power. [MW/t]	HZP HFP 552.833 900 552.833 600 552.833 557 753.978 460.72 3.293 3.293	

Fig. 1. The configuration of PB-2 BWR unit cell.

moderator cladding	Para	meter		Value
····/	Unit	cell pitch, [mm]		14.427
gap	Fuel	pellet diameter, [1	mm]	9.391
fuel nin	Fuel	pellet material		UO ₂
	Fuel	density, [g/cm ³]		10.283
>	Fuel	enrichment, w/o		4.85
	Clade	ding outside diam	eter, [mm]	10.928
	Clade	ding thickness, [m	ım]	0.673
\smile	Clade	ding material		Zircaloy-4
L	Clade	ding density, [g/cr	m ³]	6.55
p – pitch of the unit cell	Gap material		He	
	Mode	Moderator material		H ₂ O
Parameter / Reactor condition		HZP	HFP	
Fuel temperature, [K]		551	900	
Cladding temperature, [K]		551	600	
Moderator (coolant) temperature,	[K]	551	562	
Moderator (coolant) density, [kg/n	n ³]	766	748.4	
Reactor power ,[MWt]		2.772	2 772	

Fig. 2. The configuration of TMI-1 PWR unit cell.

and a BWR) [13]. These assemblies are shown in Figures 3 and 4 while Table 1 shows both fuel assemblies data.

As a result of implementing both fuel assemblies in TRITON/NEWT code for transport calculation, the layout outputs for both types of assemblies have been obtained, as shown in Figures 5 and 6.

There are different methodologies used in our calculation of this work. These methods cover the deterministic approach (TRITON/NEWT) and Monte Carlo methodology (TRITON/KENO and SERPENT), as well as the Generalized Perturbation Theory (TSUNAMI) to stochastic sampling techniques (SAMPLER).

4 Results

In this section, the results of transport calculation with TRITON/NEWT, SERPENT-2 and TRITON/KENO are



Fig. 3. PB-2 BWR assembly design.



Fig. 4. TMI-1 PWR assembly design.

Table 1. Fuel assembly data for the test cases of Exercises I-2.

Parameter	BWR	PWR
FA geometry	7×7	15×15
FA pitch (mm)	152.4	218.11
Fuel rods per assembly	49	208
Number of guide tubes per FA	_	16
Number of instrumentation	_	1
tubes per FA		
Number of GD pins per FA	4	4
Guide tube outside diameter (mm)	_	13.462
Guide tube inside diameter (mm)	_	12.649
Instrumentation tube outside	_	12.522
diameter (mm)		
Instrumentation tube inside	_	11.201
diameter (mm)		



Fig. 5. BWR and PWR without control rod 2D assembly.



Fig. 6. BWR and PWR with control rod 2D assembly.

5

 Table 2. Parameter list to compare - Exercise I-2.

Output	Description	Units
k_assembly	Eigenvalue/multiplication factor for two-group assembly	_
fuel_maca_ $1/2$	Macroscopic absorption cross-section for both groups	$1/\mathrm{cm}$
${\rm fuel_macf_1/2}$	Macroscopic fission cross-section for both groups	$1/\mathrm{cm}$
$diff_1/2$	Diffusion coefficient for both groups	cm
$flux_1/2$	Neutron flux for both groups	$1/{ m cm}^2{ m s}$

presented, as well as the sensitivity analysis and propagation of uncertainties results that were performed using TSUNAMI and SAMPLER modules. All calculations were carried out for four assemblies models (like shown in Figs. 5 and 6) and two different states (HFP and HZP), with a total of eight configurations for this test case in UAM-Benchmark, Exercise I-2.

4.1 TRITON/NEWT and SERPENT-2 results

Results for the k_{eff} and cross-section values are summarized in this section. The aim is to compare TRITON/NEWT and SERPENT-2 results with the average of Benchmark participation values. In the TRITON/NEWT modules, the 238-group nuclear data library collapse was used. Otherwise, the computation with the SERPENT-2.1.22 code was carried out with two libraries, JEFF-3.1 and ENDF/B-VII, to compare results with SCALE (ENDF/B-VII.1 library).

The output values compared in this paper are listed in Table 2. Comparison of TRITON/NEWT and SERPENT-2 are carried out for the multiplication factor (k_{eff}) , the Absorption cross-section, the Fission cross-section, the Diffusion cross-section and the Flux. All cross-section values are presented for both groups, Fast and Thermal.

In Tables 3 and 4, the first column shows the output values compared in this paper, the second one are the reference values found in UAM-Benchmark results. The third and fourth columns represent TRITON/NEWT calculations and its error with UAM-Benchmark references. Subsequently, the results of SERPENT-2 calculations and their comparison with TRITON/NEWT and UAM-Benchmark values are presented.

The reference values adopted in these tables have been calculated as an average of all submitted results of all benchmark participants, referring to the last submission of 2013.

It should be considered that each participant uses their own code, and this can introduce errors due to the different methodology of each code, but the objective of the benchmark program takes into account these discrepancies.

Therefore, in this work, we calculated the average results of all participants and it is intended to calculate the error between our simulation against SCALE and SER-PENT codes.

It is evident from Table 3 that there is good agreement between TRITON/NEWT and reference values, especially for the unrodded configurations. There is a slight recurring discrepancy in the flux value (in both groups, one and two), but it is assumed to be because of different normalization methods of implemented codes in this benchmark exercise.

In Table 3, comparing the SERPENT-2.1.22 results with TRITON/NEWT and UAM-Benchmark, it can be seen a good agreement in both models (BWR and PWR) as has been observed for SCALE results. The slight differences, especially in flux and diffusion coefficient results, can be due to the different methodologies (Monte Carlo) implemented in SERPENT code. In fact, in order to obtain most accurate results, the option B1 was adopted in SERPENT calculation while this option is not available in the SCALE beta version used in this work.

Additionally, we have compared SERPENT calculations using both JEFF-3.1 and ENDF/B-VII libraries, and the JEFF-3.1 results are more close to TRITON/NEWT values despite TRITON/NEWT uses the ENDF/B-VII library.

Even though the slight differences comparing SER-PENT with validated and reference results, it could be a good transport calculation code, ongoing testing and validation.

With the aim to show a clearest exposition of the results, the standard deviations in SERPENT, which are lower than 30 pcm for all cases presented, are avoided in Tables 3 and 4.

Table 4 shows good agreement between TRITON/ NEWT and reference values like it was presented in Table 3. Comparing SERPENT-2 results with TRITON/ NEWT and UAM-Benchmark, the results are similar, as shown in PWR configuration in Table 3.

Even though a different approach is used in the SERPENT code, there is an acceptable agreement between SERPENT-2 and TRITON/NEWT and Benchmark values. In this case, JEFF-3.1 library results are a little bit better comparing with ENDF/B-VII library.

As a conclusion, it is important to give relevance to the short computational times in transport calculation for SERPENT code that was found faster than TRITON/ NEWT code. In fact, in SERPENT calculation they have simulated 50,000 particles and 350 cycles in which the first 50 cycles were discarded because of its low statistical weight.

For all exercises, it was used a cluster composed of four blocks with 18 servers equipped with two processors Intel Xeon E5-4620 8c/16T and with a RAM of 64 GB DDR3, and $2 \times \text{interfaces } 10 \text{ GbE}.$

I-2.
exercises
configuration
Benchmark c
of PWR
values
cross-section
of
Comparison
Table 3.

Output	Benchmark	SCALE (TRITON/ NEWT)	Error (%)	SERPENT-2 (JEFF)	Error (%)	SERPENT-2 vs SCALE (%)	SERPENT-2 (ENDFB)	Error (%)	SERPENT-2 vs SCALE (%)
		(
r with_IIFF_uill	DULLEU 1 200 E - DD	1 90415-00	06.0	1 9075 - 00	040	010	1 9075 - 00	0.00	0 69
	1.030E+00	1.03412+00	00		0.10	0.40		0.02	20.0
tuel_maca_1	1.040E-02	1.077E-02	3.58	1.016E-02	2.32	6.04	1.016E-02	2.30	0.02
$fuel_maca_2$	1.090 E-01	1.098E-01	0.78	1.109E-01	1.76	0.97	1.111E-01	1.96	1.16
$fuel_macf_1$	$3.530 \text{E}{-}03$	3.613 E-03	2.35	$3.459 \text{E}{-}03$	2.00	4.44	3.459 E- 03	2.01	4.45
$fuel_macf_2$	7.720E-02	7.848E-02	1.66	7.910E-02	2.46	0.78	7.926E-02	2.66	0.98
diff 1	$1.423\mathrm{E}{+00}$	$1.430\mathrm{E}{+00}$	0.47	$1.352\mathrm{E}{+}00$	4.97	5.71	$1.353\mathrm{E}{+00}$	4.94	5.68
diff^-2	3.640E-01	3.623 E-01	0.46	4.080E-01	12.10	11.20	4.119E-01	13.17	12.04
flux l	8.648E-01	8.658E-01	0.12	8.808E-01	1.86	1.70	8.811E-01	1.88	1.73
$\frac{1}{100}$	1.352E-01	1.342E-01	0.77	1.192E-01	ll.85	12.57	1.189E-01	12.04	12.81
PWR HZP unr	odded								
$k_{assembly}$	$1.413 \mathrm{E}{+00}$	1.411E + 00	0.17	$1.404\mathrm{E}{+}00$	0.63	0.46	$1.404\mathrm{E}{+00}$	0.61	0.45
$fuel_maca_1$	1.050E-02	1.060 ± 0.02	0.95	9.976E-03	4.99	6.26	9.978E-03	4.97	6.23
fuel_maca_2	1.110E-01	1.111E-01	0.09	1.127E-01	1.54	1.42	1.131E-01	1.87	1.74
fuel macf 1	3.550E-03	$3.624 \text{E}{-}03$	2.08	3.466E-03	2.36	4.55	$3.465 \text{E}{-}03$	2.39	4.59
fuel macf 2	7.870E-02	7.933 E-02	0.80	8.041E-02	2.17	1.34	8.067E-02	2.51	1.67
diff^{-} 1	1.371E+00	$1.416E \pm 00$	3.29	$1.335\mathrm{E}{+}00$	2.59	6.04	$1.336\mathrm{E}{+00}$	2.58	6.03
diff_2	3.480E-01	3.542 E-01	1.77	4.025 E- 01	15.66	12.01	4.063E-01	16.74	12.83
flux_1	8.572E-01	$8.629 E_{-01}$	0.67	8.792 E-01	2.56	1.85	8.795 E-01	2.60	1.89
$flux_2$	1.428E-01	1.371E-01	4.01	1.209E-01	15.36	13.41	1.205E-01	15.62	13.76
PWR HFP rod	ded								
$k_assembly$	$1.065\mathrm{E}{+00}$	$1.025\mathrm{E}{+00}$	3.77	$1.011\mathrm{E}{+00}$	5.04	1.34	$1.012 \mathrm{E}{+00}$	4.98	1.27
$fuel_maca_1$	1.290E-02	1.356E-02	5.10	1.360E-02	5.41	0.29	$1.360 \text{E}{-}02$	5.45	0.33
$fuel_maca_2$	1.350E-01	1.403 E-01	3.89	1.421E-01	5.25	1.30	1.424E-01	5.45	1.48
$fuel_macf_1$	$3.420 \text{E}{-}03$	3.478E-03	1.71	$3.427 \text{E}{-}03$	0.19	1.51	3.427E-03	0.19	1.51
$\mathrm{fuel}_\mathrm{macf}_2$	7.890E-02	8.049 E-02	2.02	8.121E-02	2.93	0.89	8.138E-02	3.14	1.09
diff_1	$1.430\mathrm{E}{+00}$	$1.393\mathrm{E}{+}00$	2.58	$1.373\mathrm{E}{+}00$	3.98	1.45	$1.373\mathrm{E}{+00}$	3.98	1.45
$\operatorname{diff}_{-}2$	3.610E-01	3.659 E-01	1.37	4.207E-01	16.54	13.02	4.242 E-01	17.51	13.73
$flux_l$	9.017E-01	9.110E-01	1.03	9.147E-01	1.44	0.41	9.148E-01	1.45	0.42
flux_2	9.827E-02	8.901E-02	9.42	$8.529 \text{E}_{-}02$	13.21	4.37	8.518E-02	13.32	4.50
PWR_HZP_rod	ded								
$\rm k_assembly$	$1.095\mathrm{E}{+}00$	$1.040\mathrm{E}{+00}$	4.99	$1.037\mathrm{E}{+}00$	5.27	0.30	$1.037\mathrm{E}{+00}$	5.26	0.29
$fuel_maca_1$	1.280 E-02	1.342 E-02	4.86	1.313E-02	2.61	2.19	1.314E-02	2.63	2.18
$fuel_maca_2$	1.360E-01	1.415E-01	4.05	1.453E-01	6.82	2.60	1.458E-01	7.17	2.91

Table 3. (conti	nued).								
Output	Benchmark	SCALE (TRITON/ I NEWT) (Error (%)	SERPENT-2 (JEFF)	Error (%)	SERPENT-2 vs SCALE (%)	SERPENT-2 (ENDFB)	Error (%)	SERPENT-2 vs SCALE (%)
fuel macf 1	3.430 E-03	3.493 E-03 1	.84	3.442 E-03	0.35	1.48	3.444E-03	0.39	1.44
$fuel_macf_2$	7.970E-02	8.139E-02	2.13	$8.284 \text{E}{-}02$	3.94	1.74	8.307 E-02	4.22	2.01
diff_{-1}	$1.406\mathrm{E}{+00}$	1.381E+00 1	.77	$1.356\mathrm{E}{+}00$	3.55	1.84	$1.356\mathrm{E}{+}00$	3.55	1.84
diff 2	3.550E-01	3.583E-01 ().94	4.141E-01	16.66	13.48	4.176E-01	17.63	14.19
flux_l	9.018E-01	9.084E-01 ().73	9.128E-01	1.22	0.48	9.130E-01	1.24	0.50
$flux_2$	9.818E-02	9.156E-02 (3.74	8.724E-02	11.15	4.96	8.700 E-02	ll.38	5.24
Table 4. Comp	arison of cross-s	section values of BWR	t Benchma	urk configuration	exercises	-2.			
		SCALE	Error	SFRPENT-2	Error	SERPENT-2 vs	SERPENT-2	Error	SERPENT-2 vs
Output	Benchmark	(TRITON/NEWT)	(%)	(JEFF)	(%)	SCALE (%)	(ENDFB)	(%)	SCALE (%)
BWR_HFP_ur	ırodded								
$\rm k_assembly$	$1.076\mathrm{E}{+00}$	$1.080\mathrm{E}{+}00$	0.38	$1.081\mathrm{E}{+00}$	0.44	0.06	$1.081\mathrm{E}{+00}$	0.46	0.08
fuel_maca_1	6.880 E-03	6.934 E-03	0.79	6.848E-03	0.47	1.26	6.847 E-03	0.48	1.27
fuel_maca_2	$5.250 ext{E-02}$	5.196E-02	1.02	5.285 E-02	0.66	1.67	$5.304 \text{E}{-}02$	1.03	2.03
fuel_macf_1	1.780 E-03	1.817E-03	2.07	1.818E-03	2.13	0.06	1.818E-03	2.15	0.08
fuel_macf_2	2.610E-02	2.708E-02	3.75	$2.761 \text{E}{-}02$	5.77	1.91	2.770E-02	6.15	2.26
diff_{-1}	$1.768\mathrm{E}{+00}$	$1.645\mathrm{E}{+}00$	6.98	$1.611\mathrm{E}{+00}$	8.88	2.08	$1.611\mathrm{E}{+00}$	8.86	2.06
diff_2	$4.260 \text{E}{-}01$	4.117E-01	3.35	4.567 E-01	7.22	9.86	4.626E-01	8.58	10.99
flux_{-1}	8.015 E-01	7.765 E-01	3.12	7.905 E-01	1.37	1.77	7.911E-01	1.30	1.85
flux_2	1.985 E-01	2.235 E-01	12.59	2.095 E-01	5.53	6.69	2.089 E-01	5.24	6.99
BWR_HZP_un	urodded								
$\rm k_assembly$	$1.108\mathrm{E}{+00}$	$1.106\mathrm{E}{+00}$	0.20	$1.107\mathrm{E}{+00}$	0.07	0.13	$1.108\mathrm{E}{+00}$	0.04	0.16
${ m fuel_maca_1}$	7.150 ± 0.03	7.201E-03	0.71	7.098E-03	0.72	1.45	7.098E-03	0.73	1.45
fuel_maca_2	$5.460 ext{E-02}$	5.542 E-02	1.50	$5.621 \mathrm{E}{-}02$	2.95	1.41	5.646E-02	3.41	1.85
fuel macf 1	1.900E-03	1.909 E-03	0.47	1.912E-03	0.64	0.17	1.912 E-03	0.63	0.16

Table 4. (contin	nued).								
Output	Benchmark	SCALE (TRITON/NEWT)	Error (%)	SERPENT-2 (JEFF)	Error (%)	SERPENT-2 vs SCALE (%)	SERPENT-2 (ENDFB)	Error (%)	SERPENT-2 vs SCALE (%)
$fuel_macf_2$	2.910E-02	2.852E-02	1.99	2.894E-02	0.55	1.45	2.907E-02	0.10	1.90
diff_1	$1.496\mathrm{E}{+00}$	$1.455\mathrm{E}{+}00$	2.74	$1.387\mathrm{E}{+00}$	7.27	4.89	$1.388\mathrm{E}{+00}$	7.25	4.86
diff_2	3.320 E-01	3.367 E-01	1.40	3.646E-01	9.81	7.66	3.696E-01	11.31	8.91
flux_1	$7.330 \text{E}{-}01$	7.284E-01	0.62	7.411E-01	1.11	1.71	7.420E-01	1.22	1.83
flux_2 RWR HFP rod	2.670E-01 dded	2.716E-01	1.71	2.589 E-01	3.03	4.89	2.581E-01	3.36	5.24
k assembly	7.870E-01	7.689E-01	2.31	7.957E-01	1.11	3.37	7.950E-01	1.02	3.29
fuel maca 1	9.530E-03	9.121E-03	4.29	9.697 E-03	1.75	5.93	9.698E-03	1.76	5.95
fuel maca 2	7.100 E - 02	7.334E-02	3.30	7.210E-02	1.55	1.72	7.232E-02	1.86	1.41
fuel macf 1	1.830 E-03	1.750E-03	4.38	1.816E-03	0.74	3.67	1.817E-03	0.72	3.69
$fuel_macf_2$	$3.070 \text{E}{-}02$	3.097 E-02	0.87	$3.084 \text{E}{-}02$	0.46	0.41	$3.092 \text{E}{-}02$	0.70	0.16
diff_{-1}	$1.713\mathrm{E}{+00}$	$1.698\mathrm{E}{+}00$	0.87	$1.603\mathrm{E}{+}00$	6.44	5.96	$1.603\mathrm{E}{+}00$	6.43	5.94
diff_2	$4.680 \text{E}{-}01$	5.195E-01	11.01	4.892 E-01	4.53	6.19	$4.945 \text{E}{-}01$	5.65	5.07
$flux_1$	8.696E-01	8.827E-01	1.51	8.545 E-01	1.73	3.30	8.550 E-01	1.67	3.24
$flux_2$	1.304E-01	1.173E-01	10.10	1.455 E-01	11.55	19.40	1.450E-01	11.18	19.13
BWR_HZP_rod	dded								
$\rm k_assembly$	8.620 E-01	8.574E-01	0.53	8.667 E-01	0.54	1.06	8.662 E-01	0.49	1.02
fuel_maca_1	9.790 E-03	9.503E-03	2.93	9.901E-03	1.14	4.03	9.898E-03	1.10	3.99
fuel_maca_2	7.410E-02	7.622E-02	2.87	$7.590 ext{E-02}$	2.43	0.43	7.619 E-02	2.83	0.04
$fuel_macf_1$	1.970E-03	1.905E-03	3.31	1.932E-03	1.92	1.41	1.932E-03	1.91	1.43
$\mathrm{fuel}_\mathrm{macf}_2$	3.180 E - 02	$3.249 \text{E}{-}02$	2.16	3.313 E-02	4.20	1.95	$3.323 \text{E}{-}02$	4.49	2.23
diff_{-1}	$1.466\mathrm{E}{+00}$	$1.429\mathrm{E}{+}00$	2.50	$1.387\mathrm{E}{+00}$	5.37	3.03	$1.388\mathrm{E}{+00}$	5.34	2.99
$\operatorname{diff}_{-}2$	$3.390 \text{E}{-}01$	3.467E-01	2.27	$3.803 \text{E}{-}01$	12.19	8.84	3.848E-01	13.51	9.90
flux_1	8.029 E-01	8.091E-01	0.77	8.069 E-01	0.50	0.27	8.074E-01	0.56	0.21
flux 2	1.971E-01	1.909 E-01	3.14	1.931E-01	2.02	1.14	1.926E-01	2.28	0.88

Table 5. Comparison of KENO and NEWT k_{eff} assembly results.

	TRITON/KENO	TRITON/NEWT	Error (%)
PWR HFP unrodded	1.40031 ± 0.00036	$1.394E{+}00$	0.45
PWR HZP unrodded	1.41584 ± 0.00044	$1.411E{+}00$	0.34
PWR_HFP_rodded	1.02114 ± 0.00057	$1.025\mathrm{E}{+00}$	0.38
PWR_HZP_rodded	1.03528 ± 0.00047	$1.040\mathrm{E}{+00}$	0.45

BWR_HFP_unrodded Uncertainty Information

The relative standard deviation of K $_{\rm eff}$ (% $\Delta k/k$) due to cross-section covariance data is: 0.5358 % $\Delta k/k$ contributions to uncertainty in K $_{\rm eff}$ (% $\Delta k/k$) by individual energy covariance matrices:

Covarian	e Matrix	Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	3.0317E-01
²³⁵ U _{nubar}	²³⁵ U _{nubar}	1.7051E-01
²³⁸ U _{n,n'}	²³⁸ U _{n,n'}	1.6967E-01
235U	235 U a samma	1.4249E-01

Forward Calculation: K_{eff} = 1.08039129

Adjoint Calculation: K_{eff} = 1.08035940

BWR_HFP_rodded Uncertainty Information

The relative standard deviation of K_{eff} (%Δk/k) due to cross-section covariance data is: 0.6012 % Δk/k contributions to uncertainty in K_{eff} (%Δk/k) by individual energy covariance matrices:

Covarian	ce Matrix	Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁸ U _{n,n'}	²³⁸ U _{n,n'}	3.0300E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	2.9372E-01
²³⁵ U nubar	²³⁵ U _{nubar}	2.4778E-01
²³⁵ U _{chi}	²³⁵ U _{chi}	2.0790E-01

Forward Calculation: K_{eff} = 0.76879110 Adjoint Calculation: K_{eff} = 0.76878164

BWR_HZP_unrodded Uncertainty Information

The relative standard deviation of K_{eff} (% $\Delta k/k$) due to cross-section covariance data is: 0.5061 % $\Delta k/k$ contributions to uncertainty in K_{eff} (% $\Delta k/k$) by individual energy covariance matrices:

Covariand	e Matrix	Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	2.7926E-01
²³⁵ U _{nubar}	²³⁵ U _{nubar}	2.7820E-01
²³⁵ U _{n,gamma}	²³⁵ U _{n,gamma}	1.4214E-01
235U fission	235U fission	1.4189E-01

Forward Calculation: K_{eff} = 1.10596891 Adjoint Calculation: K_{eff} = 1.10593861

BWR_HZP_rodded Uncertainty Information

The relative standard deviation of K _{eff} (% $\Delta k/k$) due to cross-section covariance data is: 0.4936 % $\Delta k/k$ contributions to uncertainty in K _{eff} (% $\Delta k/k$) by individual energy covariance matrices:

		1
Covariance Matrix		Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁵ U _{nubar}	²³⁵ U _{nubar}	2.6840E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	2.5420E-01
²³⁵ U _{fission}	235U fission	1.4959E-01
238U	238U a.a.	1.3228E-01

Forward Calculation: K_{eff} = 0.85751511 Adjoint Calculation: K_{eff} = 0.85750102

Fig. 7. Most important contributor to uncertainty in k_{eff} (% $\Delta k/k$) in BWRs.

A comparison of the total computational time between SERPENT and TRITON/NEWT calculations, for PWR HFP unrodded configuration, returns:

- SERPENT total computational time (seconds): 935;

 TRITON/NEWT total computational time (seconds): 1358.

4.2 TRITON/KENO and SERPENT-2 results

From TRITON/KENO calculation, the best estimate system of k_{eff} in 3D Monte Carlo methodology is obtained, and it is possible to compare these results with deterministic values as shown in Table 5. In this table, the first column shows TRITON/KENO results, the k_{eff} values obtained in TRITON/NEWT figured in the second column, and its error with TRITON/KENO was finally reported.

Analysing Table 5 is possible to corroborate a good agreement through both SCALE models since the largest error calculated has been 0.45%.

4.3 TSUNAMI-2D results

TSUNAMI module is very useful for sensitivity analysis and propagation of uncertainties of cross-sections. TSUNAMI employs the Generalized Perturbation Theory and determines the sensitivities coefficient for each nuclide. Furthermore, the sensitivity coefficients for total reaction for each nuclide and mixture can be calculated with this module.

A list of four major contributors to the uncertainty in k_{eff} by individual energy covariance matrices is presented in Figure 7 for BWR calculations and in Figure 8 for PWR calculations.

In both figures, it is possible to find out that the list of major contributors does not vary greatly from case to case, and for all cases uranium seems to be responsible for the uncertainty of k_{eff} . In particular, $^{238}U_{n,\gamma}$, $^{235}U_{nubar}$, $^{238}U_{n,n'}$, are present at the top of contributors list.

Finally, looking at these results is possible to wise up the relative standard deviation due to cross-section covariance data. The relative standard deviation has been close to 0.50% in all cases, for both BWRs and PWRs calculations.

Likewise, it is interesting to analyse the sensitivity profiles of these major contributors, as shown in Figures 9 and 10. The

PWR_HFP_unrodded Uncertainty Information
The relative standard deviation of K $_{\rm eff}$ (% $\Delta k/k$) due to cross-section covariance data is: 0.4713 % $\Delta k/k$
contributions to uncertainty in K_{eff} (% $\Delta k/k$) by individual energy covariance matrices:

Covariance Matrix		Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
235U nubar	235U nubar	2.6798E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	2.5618E-01
²³⁵ U _{n,gamma}	²³⁵ U _{n,gamma}	2.0147E-01
235U firming	235U	1.0809E-01

Forward Calculation: $K_{eff} = 1.39405262$ Adjoint Calculation: $K_{eff} = 1.39403963$

PWR_HFP_rodded Uncertainty Information

The relative standard deviation of K _{eff} (% $\Delta k/k$) due to cross-section covariance data is: 0.5552 % $\Delta k/k$ contributions to uncertainty in K _{eff} (% $\Delta k/k$) by individual energy covariance matrices:

Covariance Matrix		Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁸ U _{n,n'}	²³⁸ U _{n,n'}	6.0318E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	3.5015E-01
²³⁵ U _{n,gamma}	²³⁵ U _{n,gamma}	3.0614E-01
235U chi	²³⁵ U chi	2.3938E-01

Forward Calculation: K_{eff} = 1.06274374 Adjoint Calculation: K_{eff} = 1.06274353 PWR_HZP_unrodded Uncertainty Information

The relative standard deviation of K_{eff} (% $\Delta k/k$) due to cross-section covariance data is: 0.4651 % $\Delta k/k$ contributions to uncertainty in K_{eff} (% $\Delta k/k$) by individual energy covariance matrices:

Covariance Matrix		Contributions to Uncertainty in K _{eff} (%Δk/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁵ U _{nubar}	235U nubar	2.6906E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	2.4790E-01
²³⁵ U _{n,gamma}	²³⁵ U _{n,gamma}	2.0090E-01
235U fission	235U n gamma	1.0834E-01

Forward Calculation: $K_{eff} = 1.41055640$ Adjoint Calculation: $K_{eff} = 1.41054241$

PWR_HZP_rodded Uncertainty Information

The relative standard deviation of K _{eff} (% $\Delta k/k$) due to cross-section covariance data is: 0.5273 % $\Delta k/k$ contributions to uncertainty in K _{eff} (% $\Delta k/k$) by individual energy covariance matrices:

Covariance Matrix		Contributions to Uncertainty in K _{eff} (%∆k/k)
Nuclide-Reaction	Nuclide-Reaction	Due to this Matrix
²³⁸ U _{n,n'}	²³⁸ U _{n,n'}	5.7584E-01
²³⁸ U _{n,gamma}	²³⁸ U _{n,gamma}	3.3700E-01
²³⁵ U _{n,gamma}	²³⁵ U _{n,gamma}	3.0398E-01
²³⁵ U _{chi}	²³⁵ U _{chi}	2.3223E-01

Forward Calculation: K_{eff} = 1.08104281 Adjoint Calculation: K_{eff} = 1.08104261

Fig. 8. Most important contributor to uncertainty in k_{eff} ($\Delta k/k$) in PWRs.



Fig. 9. Sensitivity profiles of most important contributor to uncertainty in k_{eff} , BWRs cases.

sensitivity per unit lethargy profiles looks similar, with peaks ranging from 0.28 to 0.35. Sensitivity at Hot Full Power state has emerged greater than sensitivity at Hot Zero Power state, for all cases of study. Moreover, sensitivity profile varies only slightly from case to case, and they do not lead to changes in the uncertainty of the k_{eff} . Figures 9 and 10 represent BWRs and PWRs calculations, respectively.

4.4 SAMPLER results

SAMPLER calculation provides not only estimates for the expected values of the data but also covariance data describing the correlated uncertainty. SAMPLER repeats perturbation steps for a specified number of samples to obtain a distribution of results that can be converted to a



Fig. 10. Sensitivity profiles of most important contributor to uncertainty in k_{eff} , PWRs cases.



Fig. 11. Histogram plot for k_{eff} and two cross-section calculations, for BWR HZP rodded case.



$PWR_HFP_unrodded_fast$ group and k_{eff}

Fig. 12. Histogram plot for $k_{e\!f\!f}$ and two cross-section calculations, for PWR HFP unrodded case.



Fig. 13. Samples population for k_{eff} and two cross-section calculations (fission and flux), for BWR_HZP with control rod.



PWR_HFP_unrodded_fast group and keff

Fig. 14. Samples population for k_{eff} and two cross-section calculations (fission and flux), for PWR_HFP without control rod.

standard deviation and correlation coefficients. The SCALE Criticality Safety Analysis Sequence (CSAS) with the 238-group nuclear data library was used for the computations.

Based on the Wilks' approach [14], the sample size for double tolerance limits with a 95% of uncertainty and with 95% of statistical confidence for the output variables is equal to 146 samples [15], which is the number of runs performed in this work.

To provide information on sampling convergence, layout response of SAMPLER results is presented. For every case run within SAMPLER, any number of responses can be extracted. For instance, the histogram plots that indicate the distribution of the k_{eff} and some cross-section computed values for these benchmark exercises are listed below. Figure 11 shows histogram results of BWR with control rod, at hot zero power state (HZP). The k_{eff} and two cross-sections (fission and flux) values for the fast and thermal group are represented.

While Figure 12 shows PWR histogram results, in unrodded configuration and at hot full power state (HFP).

In the same way, are represented the k_{eff} and fission, and flux cross-section results for the fast and thermal group.

According to the print flag set by the user, SAMPLER also prints a list of tables with interesting information, like average values and standard deviation, correlation matrices, and covariance matrices and so on. Another of interesting results of perturbed variables in SAMPLER layout is the running average, which represent average values and standard deviation for samples of the population during simulation. Moreover, in SAMPLER response it is possible to see that samples population are closer to the average value and almost entirely within its standard deviation, this is shown in the figures below. According to Wilks' theory, more than 95% of reliability is reached with 146 samples.

Figure 13 shows samples population with averaged values and standard deviation of BWR with control rod configuration, at hot zero power state (HZP). The k_{eff} and two cross-sections (fission and flux) values, for the fast and thermal group, are represented. While Figure 14 shows PWR results of samples population with averaged values and standard deviation, in unrodded configuration and at hot full power state (HFP). In the same way, are represented the k_{eff} and fission and flux cross-section results, for the fast and thermal group.

5 Conclusions

This work has been carried out in the framework of UAM-Benchmark Exercise I-1 Cell Physics and I-2 Lattice Physics. The two test cases (PB-2 BWR and TMI-1 PWR) have been analyzed in two different configurations and two different states, with the objective of quantifying the uncertainty in all step calculation and propagate uncertainties in the LWR whole system.

Transport calculations have been analyzed with the deterministic code TRITON/NEWT and stochastic code SERPENT-2.1.22 with the aim of comparing k_{eff} and

cross-sections results between both codes and with UAM-Benchmark reference values.

Sensitivity calculations have been performed with TSUNAMI module, which uses Generalized Perturbation Theory, and SAMPLER for the perturbed cross-section with stochastic sampling techniques.

The following significant conclusions can be highlighted:

- TRITON/NEWT is a solid, validated code that has performed well the UAM-Benchmark calculations but spent more computational times comparing against the SERPENT-2 results. Even though the slight differences comparing SERPENT with validated and reference results, this code could be a good transport calculation code, ongoing testing and validation;
- TSUNAMI module was adopted to estimate sensitivity and uncertainty analysis, the impact of the uncertainties in the basic nuclear data on the calculation of the multiplication factor and microscopic and macroscopic cross-sections. Uncertainties were found to be $\approx 0.5\%$ on k_{eff} . The particular $^{238}U_{n,\gamma}$, $^{235}U_{nubar}$, $^{238}U_{n,n'}$ were found to be the most important contributors to the uncertainty in these exercises. The deterministic solutions were compared with SAMPLER response, and good agreement was found for this exercise.

This work contains findings produced within the OECD/NEA UAM-Benchmark.

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