

Document downloaded from:

<http://hdl.handle.net/10251/141652>

This paper must be cited as:

Tanov, S.; Pachano-Prieto, LM.; Andersson, Ö.; Wang, Z.; Richter, M.; Pastor, JV.; García-Oliver, JM.... (01-2). Influence of spatial and temporal distribution of Turbulent Kinetic Energy on heat transfer coefficient in a light duty CI engine operating with Partially Premixed Combustion. Applied Thermal Engineering. 129:31-40.
<https://doi.org/10.1016/j.applthermaleng.2017.10.006>



The final publication is available at

<https://doi.org/10.1016/j.applthermaleng.2017.10.006>

Copyright Elsevier

Additional Information

1 **Influence of Spatial and Temporal Distribution of Turbulent Kinetic**
2 **Energy on Heat Transfer Coefficient in a Light Duty CI Engine Operating**
3 **with Partially Premixed Combustion**

4

5 **Authors:**

6 Slavey Tanov¹, Leonardo Pachano², Öivind Andersson¹, Zhenkan Wang³,
7 Mattias Richter³, José V. Pastor², José M. García-Oliver², Antonio García²

8 **Affiliation:**

9 1. Division of Combustion Engines, Lund University, Sweden

10 2. CMT Motores Térmicos – Universitat Politècnica de València, Camino de Vera s/n,
11 46022, Valencia, Spain.

12 3. Division of Combustion Physics, Lund University, Sweden

13

14 **Corresponding author**

15 Dr. Antonio García

16 CMT - Motores Térmicos /Universitat Politècnica de València

17 angarma8@mot.upv.es

18

19

20

21

22

23

24

25

26 **Abstract**

27 Emission regulations together with the need of more fuel-efficient engines have driven the
28 development of promising combustion concepts in compression ignition (CI) engines. Most of
29 these combustion concepts, lead toward a lean and low temperature combustion potentially
30 suitable to achieve lower emission and fuel consumption levels compared to conventional
31 diesel combustion. In this framework, Partially Premixed Combustion (PPC) using gasoline
32 as fuel is one of the most accepted concepts. There are numerous studies focused on studying
33 concepts such as PPC from the emissions point of view. Nonetheless, there is a lack of
34 knowledge regarding changes in heat transfer introduced by the use of these combustion
35 concepts. It is worth noting that heat transfer can be considered as a key aspect behind
36 possible engine performance improvements. Thus, the reliable estimation of this parameter is
37 of considerable importance. Additionally, a better understanding of how events such as
38 injection and combustion might affect heat transfer is also relevant.

39 To gain insight into gasoline PPC heat transfer coefficient, its evolution during late
40 compression and early expansion were studied. In particular, this work aims to analyze
41 Turbulent Kinetic Energy (TKE) spatial and temporal evolution influence on heat transfer
42 coefficient. The analysis is based on experimental TKE maps derived from Particle Image
43 Velocimetry (PIV) data. For the heat transfer coefficient estimation a modified Woschni
44 correlation has been used. Results from several injection strategies and a reference motored
45 case have been analyzed. It has been found that injection strategy has a considerable
46 influence on the TKE field and hence on heat transfer coefficient evolution.

47 **1 Introduction**

48 Low Temperature Combustion (LTC) technologies, like Homogenous Charge Compression
49 Ignition (HCCI) [1–3], Partially Premixed Combustion (PPC) [4,5] and Reactivity Controlled
50 Compression Ignition (RCCI) [6–9] have been widely studied in the last decade because they
51 are expected to achieve low engine-out emissions along with high thermal efficiency. In this
52 context, gasoline PPC has received increasingly attention due to its potential for simultaneously
53 reducing fuel consumption and NO_x emissions in gasoline Spark Ignition (SI) engines, and its
54 capability to avoid soot and NO_x emissions in CI diesel engines. Several experimental and
55 numerical investigations of this combustion type have been performed over the recent years.
56 PPC is achieved by controlling the injection events, inlet temperature and pressure, and
57 composition of the fuel-air mixture, so that it ignites close to Top Dead Center (TDC). It is

58 similar in nature to HCCI, with significantly early fuel injection in the cycle and a combustion
59 process that occurs as a sequence of auto-ignition events. The main difference between HCCI
60 and PPC is that the goal of PPC is to control auto-ignition timing with moderately early fuel
61 injections (~20 CAD bTDC) by manipulating the in-cylinder charge stratification level.

62 PPC is highly dependent on the level of in-cylinder fuel stratification at the start of combustion
63 (SOC). Recent work performed by Izadi et al [10] focused on investigation of combustion
64 stratification using single injections. They concluded that combustion stratification is low and
65 almost independent of start of injection (SOI) for early injections, while there is a remarkable
66 reverse correlation between combustion timing and stratification level for the late injections (-
67 45 to -12.5 CAD). Lee and Reitz [11] investigated characteristics of PCCI with single early
68 injection. Their results demonstrated that combustion performance and emissions are strongly
69 affected by injection timing. They also indicated that spray targeting at the surface of piston
70 bowl directly influenced emissions formation. Spray targeting point, which was located near
71 the edge of piston bowl, was considered as the optimum in PCCI combustion through different
72 engine operating conditions because the squish flow would promote mixture preparation when
73 spray is injected at this location.

74 Multiple injection strategies with early injection pulses are usually employed to promote
75 different levels of fuel stratification. Early study [12] on combustion stratification using
76 different injection strategies had shown that the combustion following triple injection is more
77 homogeneous compared to single and double injection. Manente et al. [13] employed triple
78 pulse injections to control excessive fuel stratification that causes unacceptable pressure rise
79 rates in a heavy duty diesel engine. They operated the engine at various loads using fixed SOI
80 timings for the three pulses. Kalghatgi et al. [14] tested the effectiveness of triple injection on
81 controlling heat release rates as well as improving engine performance in a small-bore diesel
82 engine fueled with RON84 gasoline. Sellnau et al. [15–17] applied a triple-injection method to
83 improve fuel economy in a light duty (LD) diesel engine operated with RON91 gasoline at 6
84 bar IMEP. They also showed that a triple injection strategy allowed the use of lower injection
85 pressures compared to single and double-injection strategies.

86 Studies on multiple injection strategies conducted with experimental fluid dynamics have rarely
87 been reported. The systematic study by the authors [12,18] in the field of experimental fluid
88 dynamics has focused on in-cylinder flow pattern and temporal evolution of turbulence level
89 under PPC conditions.

90 In view of the potential of PPC and aiming at a more comprehensive study of this combustion
91 mode, the analysis of wall heat transfer coefficient is of particular importance. It has been
92 reported that wall heat transfer affects in-cylinder physical phenomena such as droplet
93 evaporation, auto-ignition and flame-wall interaction [19]. Therefore, wall heat transfer has a
94 deep impact on the overall engine performance. On the one hand, heat losses through cylinder
95 walls reduce energy available to be converted into useful mechanical work affecting indicated
96 efficiency. On the other hand, changes in gas and surface temperature due to heat transfer might
97 also affect pollutant formation [20].

98 The aim of the current work is to gain insight into heat transfer under PPC operating conditions.
99 To that end, a modified Woschni correlation is used to estimate wall heat transfer coefficient.
100 The modified correlation is not only a function of thermodynamic conditions, but also of the
101 in-cylinder velocity and TKE fields. By using experimental data measured by means of high-
102 speed PIV [12] TKE can be derived from the velocity field. Temporal and spatial TKE
103 distribution and its influence on the heat transfer coefficient is studied. The analysis is carried
104 out for three injection strategies (i.e. single, double and triple injection) as well as for a reference
105 case under motored conditions.

106 **2 Experimental Setup**

107 **2.1 Experimental facility**

108 Experiments were performed in a Bowditch-designed single-cylinder engine modified from a
109 Volvo D5 LD diesel engine.

110 Table 1 shows the engine specifications. Due to the large top ring-land crevice required for
111 side-view imaging, the geometric compression ratio of this optical engine is lower than the
112 target of 16, typical of PPC combustion systems. The engine further allows the intake swirl to
113 be adjusted by a swirl control valve. It was operated at a swirl ratio of 2.6 through this work. A
114 Bosch common rail fuel injection system and a 5-hole solenoid Bosch injector were used for
115 fuel injection.

116

117

118

119

Engine base type	Volvo D5
Number of cylinders	5
Number of valves	4
Bore	81 mm
Stroke	93.2 mm
Connecting rod	147 mm
Displacement	0.48 L
Compression ratio (metal configuration)	1:16
Compression ratio (optical configuration)	1:11.3
Swirl ratio	2.6
Fuel Injection	
Type	Common Rail
Fuel injector Type	Solenoid
Hydraulic flow	360 cc/30s @ 100 bar
Umbrella angle	140°
Orifice diameter	0.159 mm
Number of holes	5
Hole conicity	1.5

120

121

Table 1. Engine and injection system specifications.

122 **2.2 PIV measurement system**

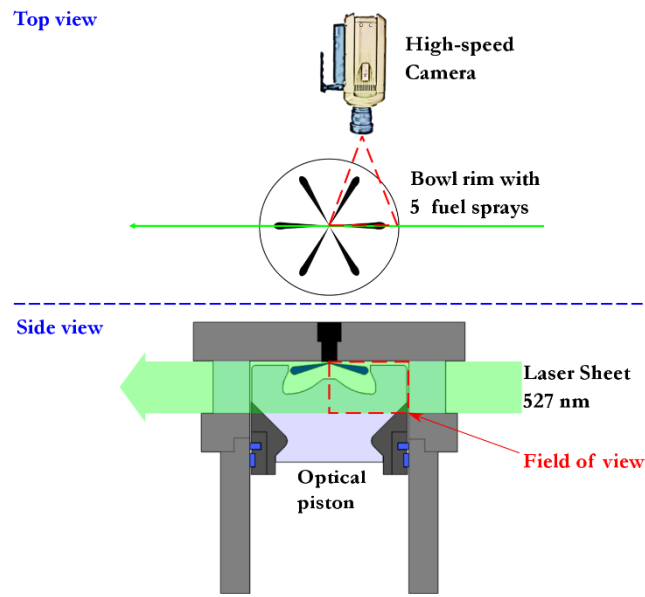
123 An Nd: YLF diode pumped dual cavity laser from Dantec Dynamics (model type: DualPower
124 30-1000) was used as the light source. Its wavelength is 527 nm and can reach a maximum 30
125 mJ power per pulse at a running repetition rate of 1 kHz. At 800 rpm engine speed the laser
126 operates at 2.4 kHz with 13 mJ energy per pulse. The light sheet created by the optics unit was
127 aligned with the injector tip installed in the optically accessible engine. The engine is equipped
128 with a production-like optical piston and an optical cylinder liner with height of 25 mm shown
129 in Figure 1. The light sheet was focused in the area between the injector tip and the liner inner
130 surface with a height of around 3 cm through the field of view.

131 Titanium Dioxide (TiO₂) powder was used as PIV seeding particles, which have a mean particle
132 diameter from 2 to 3 μm and a density of 4260 kg/m³. Assuming Stokes drag, the particle time
133 constant (τ_s) representing the response time to changes in the flow is roughly 40 μs at TDC-like
134 thermodynamic conditions. At these conditions, τ_s is slightly larger than the estimated
135 Kolmogorov time scale and thus the seeding particles are able to follow most of the turbulent
136 structures. Seeding particles were introduced from a cylindrical container fed by a swirl airflow
137 of around 20 l/min, which is precisely controlled by a mass flow meter. The seeding flow was

138 then mixed with the intake stream inside the intake manifold. The TiO_2 powder was baked over
139 24 hours before seeding it into the engine to efficiently prevent particle agglomeration. In
140 addition, water vapor was added to the intake stream, as part of EGR gases, to reduce the
141 electrostatic charge build-up, and eventually reduce the chance of adhesion of seeding particles
142 to the engine's optical components. These two procedures were proved essential for in-cylinder
143 PIV measurements.

144 Images were acquired using a Dantec Dynamics high-speed CMOS camera (SpeedSense 710).
145 A Nikkor 105 mm lens with an extension ring (Nikkor PK-13) was also used and the lens
146 aperture was closed as much as possible (f#16) to focus the full field of view. The image format
147 was cropped to 1040x440 pixels to increase the maximum camera frame rate. The exposure
148 time for all image pairs was 63 μs for the first image and around 350 μs for the second image
149 (depending on the engine rotation speed). The long exposure for the second image was due to
150 the time required to readout the first image from camera sensor. The time between the laser
151 pulses was set to 20 μs . According to [21] this time delay between laser pulses is a good
152 compromise between resolving velocity and being able to perform the cross-correlation. The
153 maximum displacement was less than one quarter of the side of the interrogation area as the
154 error increases with increasing displacement [22,23].

155 Despite the combustion process, TiO_2 powder proved to be a suitable seeding element.
156 Although soot luminosity is known to be a challenge for PIV measurements under reacting
157 conditions, PPC efficiently helps to overcome this issue. No filter was used in these
158 measurements, since little soot luminosity from PPC was observed and thus the Mie scattering
159 light from particles was dominant. The optical piston with a realistic bowl geometry brings
160 significant optical distortion to the acquired image. Considering that the glass piston thickness
161 is irregular in the optical path and that the distortion might depend on the piston position with
162 respect to the camera lens, the distortion is almost impossible to be compensated by adding
163 additional optics between the engine and the camera. Consequently, distortion was handled with
164 a code for image dewarping.



165
166

Figure 1. Optical setup.

167 3 Methodology

168 3.1 Operating conditions

169 Throughout this study, the engine was running at 800 rpm, with an injection pressure of 600
 170 bar for 44 continuously fired cycles. A blended fuel, PRF 70, consisting of 30% n-heptane
 171 and 70% iso-octane in volume percentage was used for the experiment to achieve suitable
 172 ignition delay prior to auto-ignition for PPC. Injection strategy was the main variable for
 173 comparison, and three different injection patterns (single, double and triple injection) were
 174 investigated. In these experiments the engine load was kept constant at 4 bar IMEP. At SOC,
 175 the fuel-air mixture was stratified enough to achieve stable ignition and controlled heat
 176 release. The inlet air temperature was kept at 73 °C and the intake air pressure was 1.14 bar
 177 (absolute). The experiments were performed in a randomized order after the engine reached a
 178 steady thermal state, which corresponded to a cooling water temperature of 65 °C. The
 179 exhaust of an industrial fuel-oil burner was used to decrease the oxygen level of the intake
 180 charge to 17%. The operating conditions are summarized in

181 Table 2, while heat release rate and pressure traces for each of the injection strategies
 182 proposed are shown in Figure 2.

183

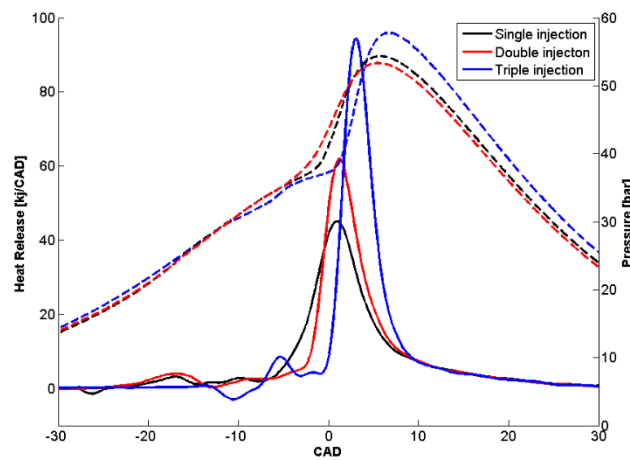
184

Engine parameters		
Intake pressure	1.14 bar	
Intake temperature	73 °C ±2 °C	
Injection pressure	600bar ±2 bar	
Swirl ratio	2.6	
O ₂	17 vol %	
λ	1.75	
Cooling water temperature	65 °C ±2 °C	
Liner wall temperature	80 °C	
Fuel type	PRF 70	
Seeding particles	TiO ₂	
Injection timing	SOI [CAD]	Duration [CAD]
Single	15	2.9
Double	60/16.5	2.9/1.9
Triple	62.5/29.5/17	2.3/1.7/1.7

186

187

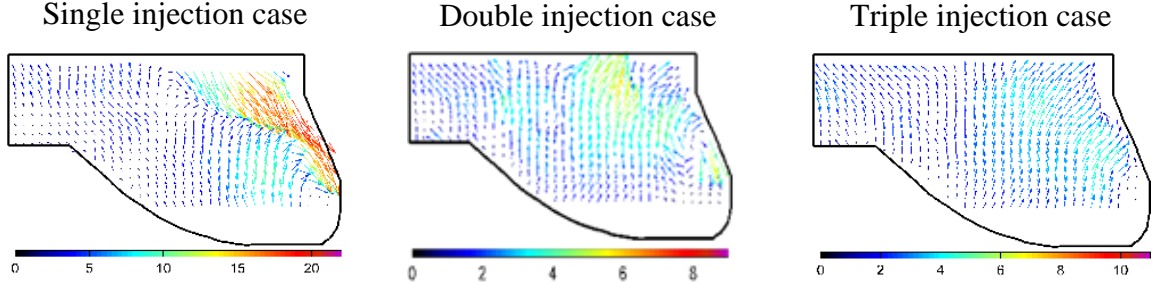
Table 2. Engine operation parameters and injection timing.



188

189 Figure 2. Rate of heat release (thick line) and pressure traces (dashed line) for different injection strategies.

190 In reference to PIV measurements, Figure 3 shows ensemble average flow fields measured at
 191 12 CAD bTDC for different injection strategies. More detailed explanation on ensemble
 192 average flow fields are reported in a previous work [12].



193 *Figure 3. Ensemble Average velocity field in m/s at -12 CAD for single, double and triple injection.*

194 **3.2 Theoretical considerations**

195 **3.2.1 Turbulent Kinetic Energy calculation**

196 Ensemble average velocity fields (coming from expressions such as Eq. (1)) are helpful to
 197 analyze the overall evolution of the flow field. In that expression, θ indicates the crank-angle
 198 and the index i refers to the cycle number (44 cycles for each crankshaft angular position).
 199 Nonetheless, the estimation of velocity fluctuations based on this approach is strongly
 200 influenced by cycle-to-cycle variations [10,12].

$$\bar{U}_{EA}(\theta) = \frac{1}{N} \sum_{i=1}^N \bar{U}(\theta, i) \quad (1)$$

201 To overcome this issue, flow field variations can be estimated within a temporal window.
 202 Following such approach, velocity field variations can be solved on a cycle-resolved basis.
 203 Consequently, for a given direction in the flow the cycle-resolved average can be solved at
 204 each pixel in the domain following Eq. (2):

$$\bar{U}(\theta, i) = \frac{1}{M} \sum_{j=1}^M C(j) \cdot U\left(\theta + \frac{2j - M - 1}{2} \Delta\theta\right) \quad (2)$$

205 The temporal window size is defined by the scalar M , which was set to five for this study i.e.
 206 the cycle-resolved average velocity field at TDC is a function of the instantaneous velocity
 207 fields measured at -4, -2, 0, 2 and 4 CAD. The variable $\Delta\theta$ is the temporal resolution of the
 208 measurements (2 CAD). Finally, the term $C(j)$ is a sine weight factor calculated according to
 209 Eq. (3):

$$C(j) = \frac{\sin\left(\frac{j}{M+1}\pi\right)}{\sum_{j=1}^M \sin\left(\frac{j}{M+1}\pi\right)} \quad (3)$$

210 Where: $\sum_{j=1}^M C(j) = 1$.

211 Once the cycle-resolved average field has been solved, the instantaneous fluctuations can be
 212 computed as the standard deviation within the temporal window following Eq. (4):

$$u'(\theta, i) = \sqrt{\frac{1}{M} \sum_{j=1}^M \left[U\left(\theta + \frac{2j - M - 1}{2} \Delta\theta\right) - \bar{U}(\theta, i) \right]^2} \quad (4)$$

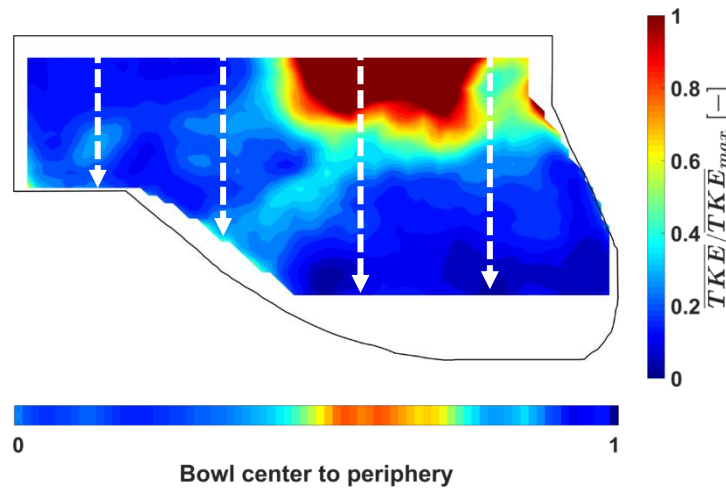
213 The instantaneous fluctuating components of velocity are then averaged over the total number
 214 of measurements (N is equal to the number of measured cycles) as expressed in Eq. (5):

$$\bar{u}'(\theta) = \frac{1}{N} \sum_{i=1}^N u'(\theta, i) \quad (5)$$

215 Lastly, the TKE associated to the fluctuations of velocity derived from PIV measurements is
 216 computed for every pixel in the domain using Eq. (6):

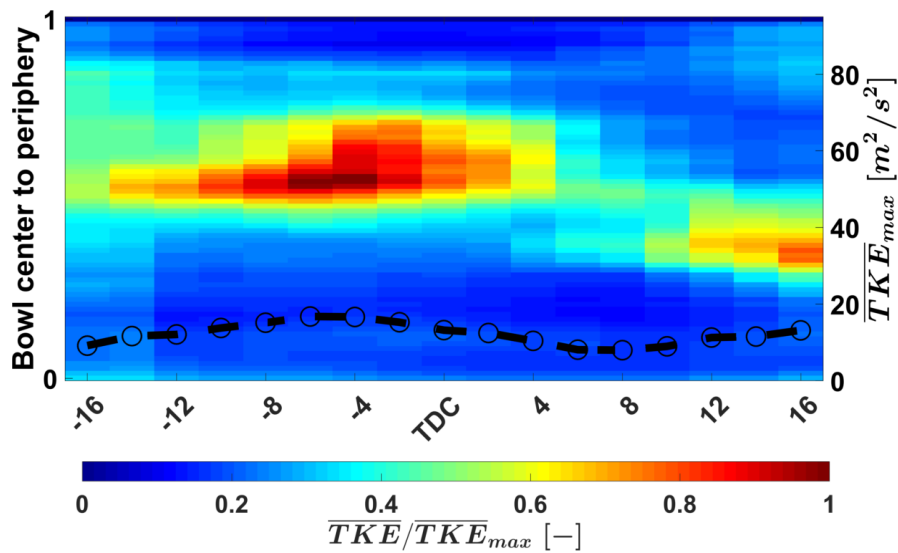
$$TKE(\theta) = \frac{1}{2} (\bar{u}'(\theta)^2 + \bar{v}'(\theta)^2) \quad (6)$$

217 Based on the previous theoretical considerations, a particular methodology has been
 218 developed to comprehensively analyze spatial and temporal evolution of TKE. Figure 4 shows
 219 the TKE field at TDC under motored conditions. The contour of the bowl is marked with a
 220 solid black line. TKE is then averaged along the axial direction (marked with white dashed
 221 arrows). This averaged TKE is plotted at the bottom of the figure where 0 and 1 mark the
 222 bowl center and periphery, respectively. TKE has been normalized by the maximum value in
 223 the domain. This normalization allows for a more direct comparison of TKE spatial evolution
 224 under firing conditions.



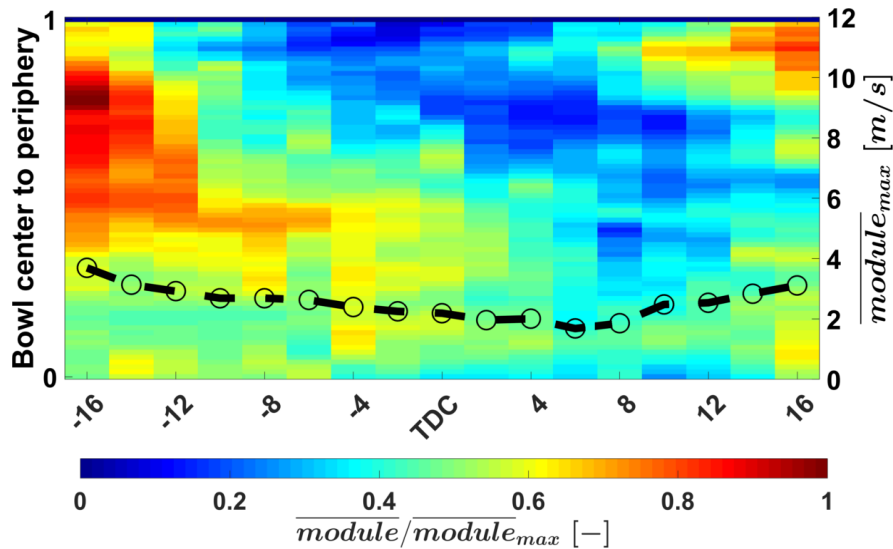
225 *Figure 4. Normalized TKE at TDC under motored condition.*

226 The same procedure described in Figure 4 is applied at all crank-angles at which PIV data is
 227 available. Normalized mean TKE plots are then rearranged as shown in Figure 5. Each
 228 column shows the spatial distribution of normalized TKE averaged over the axial direction.
 229 Then, on the right side of Figure 5 the temporal evolution of the maximum TKE value in the
 230 domain is also included. Under motored conditions, high TKE zones are located around the
 231 bowl middle area (between bowl center and periphery) as piston approaches TDC. In terms
 232 of magnitude, TKE maximum level also follows piston motion with increasing levels during
 233 late compression stroke. Once piston has reached TDC maximum levels decrease as
 234 expansion takes place with a moderate increase after 10 CAD aTDC (consistent with an
 235 increase in velocity magnitude, see Figure 6).



236 *Figure 5. Normalized mean cycle-resolved TKE for motored condition case.*

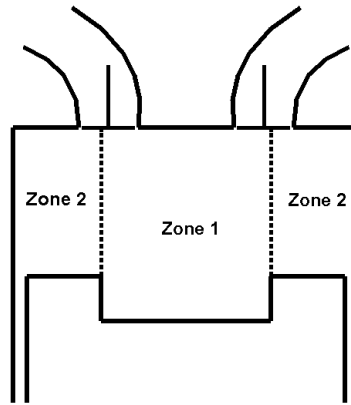
237 As in the case for TKE, velocity module has also been normalized and averaged in order to
 238 analyze the overall spatial and temporal evolution of the velocity field. Figure 6 shows the
 239 normalized mean ensemble average velocity module along with a plot (dashed black line) of
 240 the maximum value at each crank-angle. In this regard, maximum velocity values does not
 241 change significantly. In terms of spatial distribution, highest velocities take place at 16 CAD
 242 bTDC predominantly around bowl middle zone and periphery. Towards TDC, airflow
 243 decelerates at bowl periphery with predominant high velocities near bowl center. Throughout
 244 early expansion module of velocity starts raising again with dominant high velocity at bowl
 245 periphery.



246 *Figure 6. Normalized mean ensemble average velocity module for motored condition case.*

247 **3.3 1D model for tangential velocity estimation**

248 It has been stated in the previous section that the PIV setup used allows to measure the in-
 249 cylinder radial (V) and axial (U) components of velocity. Nonetheless, if the flow field is to
 250 be fully characterized the tangential component of velocity (swirl velocity) should be also
 251 studied. In that sense, the average evolution of this component is predicted by means of a 1D
 252 model [24]. The model is developed under the assumption that the in-cylinder volume can be
 253 divided into two zones as shown in Figure 7. Zone 1 considers the volume inside and above
 254 piston bowl meanwhile zone 2 considers the annular volume above piston bowl. It is also
 255 assumed that charge inside these volumes follows a solid-body rotation movement. Based on
 256 these statements, evolution of tangential velocity can be obtained resolving the change of
 257 angular momentum over time. This change can be solved for both volumes in terms of the
 258 increase and decrease of momentum flux. Increase of momentum flux is considered to be
 259 caused by mass flow coming into the volumes and momentum exchange between them due to
 260 squish flow and viscous shear forces. On the other hand, decrease of momentum flux is
 261 modeled by wall friction losses and mass flow going out of the volumes. The theoretical
 262 model has been validated with experimental PIV data measured in a horizontal section in an
 263 optical engine. The optical engine used as reference to validate 1D model predictions is a
 264 single-cylinder LD CI optical diesel engine. Engine displacement is 544.8 cm³ with a
 265 cylindrical bowl. For validation, engine was operated under motored conditions and PIV
 266 measurements were carried out from 10 CAD bTDC up to 20 CAD aTDC with a 10 CAD
 267 temporal resolution.



268

269

Figure 7. In-cylinder volumes used in the 1D model for tangential velocity prediction.

270

Solid line in Figure 8 shows the time evolution of tangential velocity (zone 1) predicted by the

271

1D model. Meanwhile, prediction of tangential velocity for the reference engine is plotted

272

with a dashed line. Predicted values show proper agreement with experimental data marked

273

with squares. For both cases, the increase in tangential velocity as piston approaches TDC is

274

an expected trend. During compression, mass flow is confined inside the bowl inducing an

275

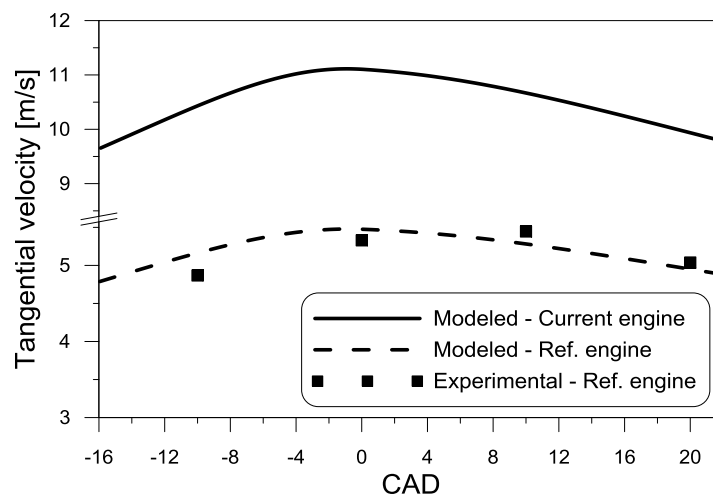
increase in tangential velocity. Differences in tangential velocity between reference engine

276

and the engine used for the current work are mainly caused by differences in cylinder head

277

and intake port geometry as well as operating conditions.



278

279

Figure 8. Tangential velocity evolution predicted by the 1D model.

280

3.4 Experimental wall heat transfer coefficient calculation

281

Based on the flow field characterization described in the previous sections an estimation of

282

the convective wall heat transfer coefficient can be achieved. For that purpose, a variation of

283

Woschni's approach [25] is used following Eq. (7)

$$h(\theta) = 0.012 \cdot D^{-0.2} \cdot p(\theta)^{0.8} \cdot T(\theta)^{-0.53} \cdot V_c(\theta)^{0.8} \quad (7)$$

284 In the former expression, wall heat transfer coefficient can be estimated over time (θ makes
 285 reference to crank-angle) as a function of engine geometry (D is the bowl diameter), in-
 286 cylinder thermodynamic conditions assumed to be uniform (p and T are pressure and
 287 temperature, respectively) and the mean flow velocity. In the original empirical correlation
 288 presented by Woschni the mean flow velocity was assumed to be a function of piston mean
 289 velocity and the gas velocity that accounted for combustion influence. This approach
 290 introduces two additional constants that need to be fitted with experimental data. In this work
 291 the flow field influence on heat transfer coefficient is accounted by means of a characteristic
 292 velocity (V_c) as described in [24]. Although Woschni's empirical correlation is based on a
 293 global one-zone view of heat transfer it has been reported to be applicable for conventional
 294 diesel combustion as well as for low temperature combustion [26,27]. In the context of this
 295 work, Woschni's modified correlation is suitable to analyze trends and to study the influence
 296 of the injection pattern and the combustion process on the heat transfer coefficient. The
 297 aforementioned velocity is calculated following Eq. (8):

$$V_c(\theta) = \sqrt{V_{module}^2(\theta) + V_{tang}^2(\theta) + 2\overline{TKE}(\theta)} \quad (8)$$

298 Velocity module calculated with Eq. (9):

$$V_{module} = \sqrt{\overline{U}_{EA}^2 + \overline{V}_{EA}^2} \quad (9)$$

299 Magnitudes for V_{module} and \overline{TKE} are derived from the PIV measurements and averaged over
 300 a layer with thickness equal to 5% of a characteristic length of the bowl (maximum radius of
 301 the bowl for this study) [28].

302 Experimentally derived magnitudes are functions of radial and axial components of velocity.
 303 On the other hand, V_{tang} is predicted by the 1D model previously described since the PIV
 304 setup used does not allow for this component of velocity to be measured. The use of this
 305 characteristic velocity leads to a more comprehensive estimation of wall heat transfer
 306 coefficient since flow field experimental data is used.

307 **4 Results and discussion**

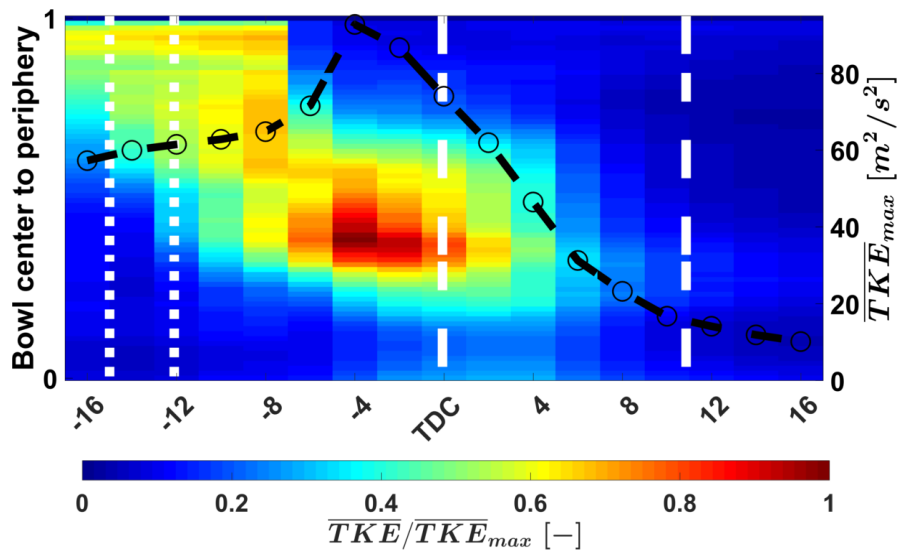
308 **4.1 Injection strategy effect on TKE**

309 In this section, TKE spatial and temporal evolution is analyzed. The influence of injection and
 310 combustion events is studied for single and multiple injection strategies. Then, results from

311 the estimation of wall heat transfer temporal evolution are presented. At the end of the
 312 section, a comprehensive analysis of injection and combustion influence on TKE and wall
 313 heat transfer coefficient is made.

314 4.1.1 Single injection

315 TKE spatial and temporal evolution for the single injection strategy is presented in Figure 9.
 316 Dotted lines mark SOI (15 CAD bTDC) and EOI (12.1 CAD bTDC). In a similar way, dashed
 317 lines mark the combustion event at the crank-angle where 10% of the total heat release has
 318 been reached (CA10) and the crank-angle at which 90% of the total heat release has been
 319 reached (CA90). In addition, temporal evolution of the maximum TKE level in the domain is
 320 plotted on the right-side axis of the figure.

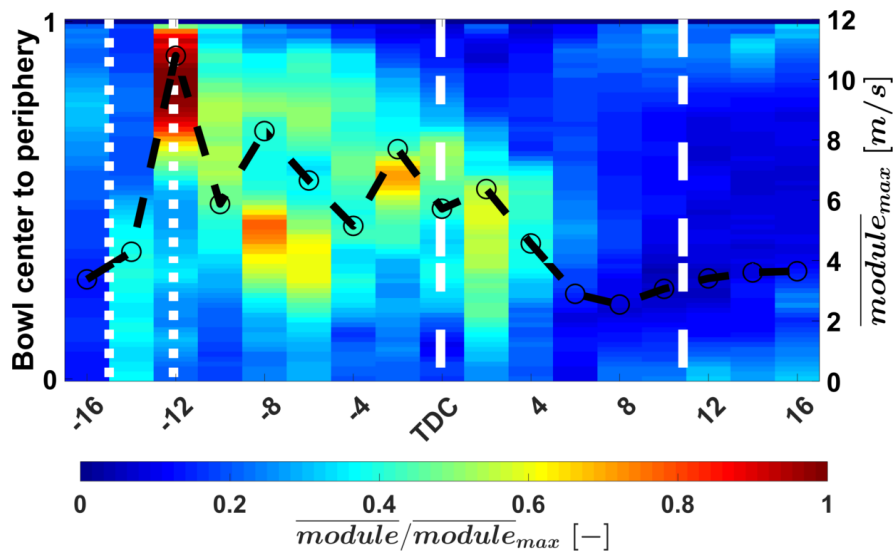


321 *Figure 9. Normalized mean cycle-resolved TKE for single injection strategy.*

322 In terms of spatial distribution, the injection event seems to induce high TKE levels especially
 323 near bowl periphery. This high TKE zones also seem to be broader up to 8 CAD bTDC. After
 324 this point, high TKE zones are located closer to bowl middle zone (between bowl center and
 325 periphery) as the squish flow becomes predominant. During the combustion process, the high
 326 TKE zone remains centered until it finally dissipates after 8 CAD aTDC. In terms of
 327 magnitude, injection event increases TKE level compared to the levels for motored case (see
 328 Figure 6). Additionally, squish flow seems to cause the peak value at the latest stage of the
 329 compression stroke. Concerning combustion, no TKE increase is introduced by this process.
 330 On the contrary, TKE keeps to continuously dissipate as far as 16 CAD aTDC.

331 Figure 10 shows a global view of how injection and combustion processes influence the flow
 332 field. As it was expected, before the start of the injection spatial distribution and maximum

333 value of velocity are consistent with results from the motored case. Once injection takes
 334 place, it is evident that this event strongly affects the flow field pattern. In terms of spatial
 335 distribution, the highest velocity level is reached around bowl periphery at 12 CAD bTDC.
 336 Beyond this point, moderate velocity levels are observed with the lowest levels taking place
 337 during expansion. Regarding temporal evolution of the maximum module of velocity (black
 338 dashed line), it can be seen that it is highly influenced by fuel injection with a sharp raise at
 339 the end of this event. Following fuel injection, maximum module of velocity increases and
 340 decreases alternatively at the same temporal range where high TKE levels were observed in
 341 Figure 9. In reference to combustion, after an initial raise of velocity around CA10 maximum
 342 velocity values drop to levels similar to the motored case levels.



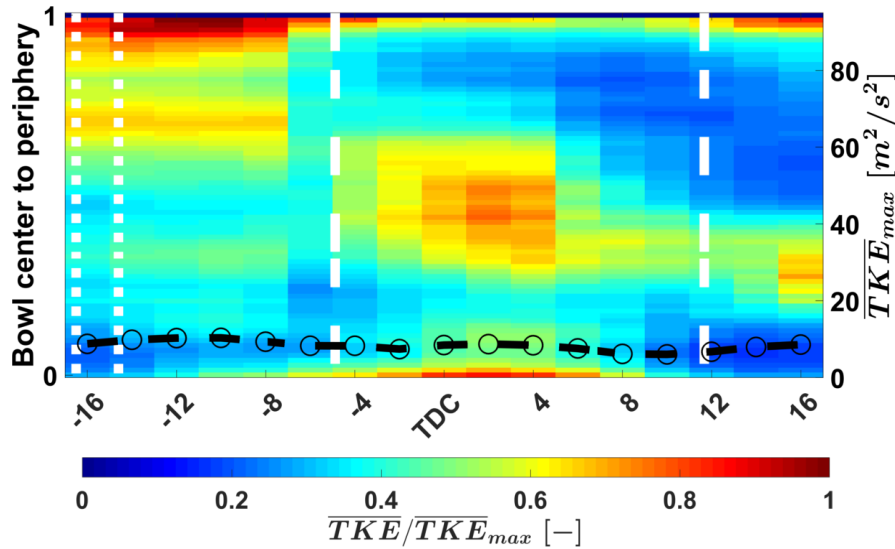
343 *Figure 10. Normalized mean ensemble average velocity module for single injection strategy.*

344 **4.1.2 Double injection**

345 The response of the in-cylinder TKE field to a double injection strategy was also evaluated.
 346 For that purpose, fuel was first injected at 60 CAD bTDC during 2.9 CAD with a second
 347 injection at 16.5 CAD with a 1.9 CAD duration (1 CAD shorter than the one for the single
 348 injection strategy). SOI and EOI of the last fuel injection are marked with dotted lines in
 349 Figure 11. On the other hand, dashed lines mark CA10 and CA90 for the combustion event.

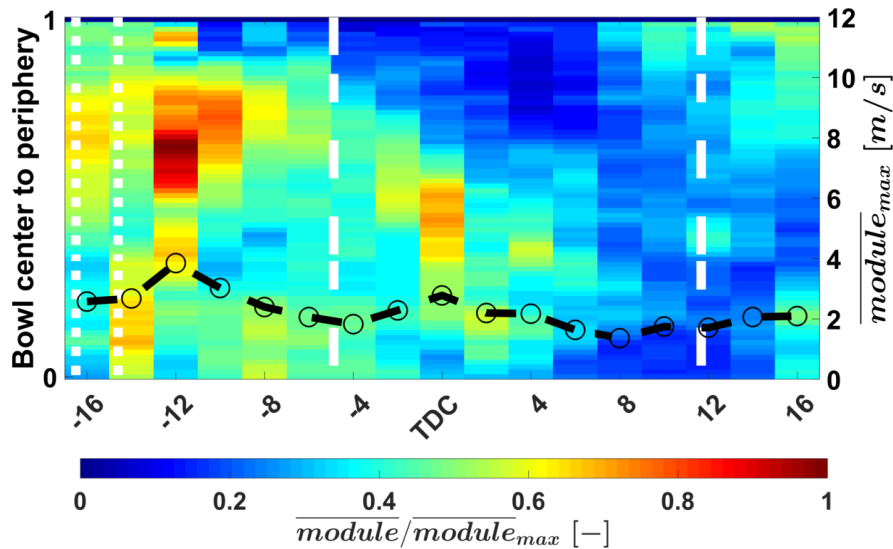
350 Globally, TKE spatial distribution follows a close trend compared to the one observed with a
 351 single injection. The highest TKE levels are located at bowl periphery with more moderate
 352 levels towards bowl center. Such distribution remains unchanged as far as 8 CAD bTDC. This
 353 last observation differs from what was observed when a single fuel injection occurred. For
 354 that case, height of the moderate TKE zone only increased as piston approached TDC. This
 355 behavior is mainly related with the fact that with a single injection pattern, no TKE inducing

356 event takes place before fuel injection. On the other hand, in the double injection strategy
 357 TKE induced by the early injection at 60 CAD bTDC still seems to contribute to reach
 358 moderate TKE levels closer to bowl center. As for combustion, high TKE levels are reached
 359 in the middle area after CA10 and dissipate beyond 8 CAD aTDC with a moderate increase
 360 near bowl center at 16 CAD aTDC. Finally, the black dashed line shows that the highest TKE
 361 level in the domain does not change drastically over time following a smooth evolution
 362 independent from injection and combustion events.



363 *Figure 11. Spatial and temporal evolution of TKE for double injection strategy.*

364 Normalized mean velocity module is shown in Figure 12. In this case, the second injection
 365 does not seem to heavily influence flow field contrary to what was observed for the single
 366 injection pattern. Spatially, velocity levels are more uniformly distributed up to CA10. During
 367 combustion, higher velocities take place closer to bowl center. In terms of magnitude, levels
 368 and temporal evolution are similar to results from the motored case.

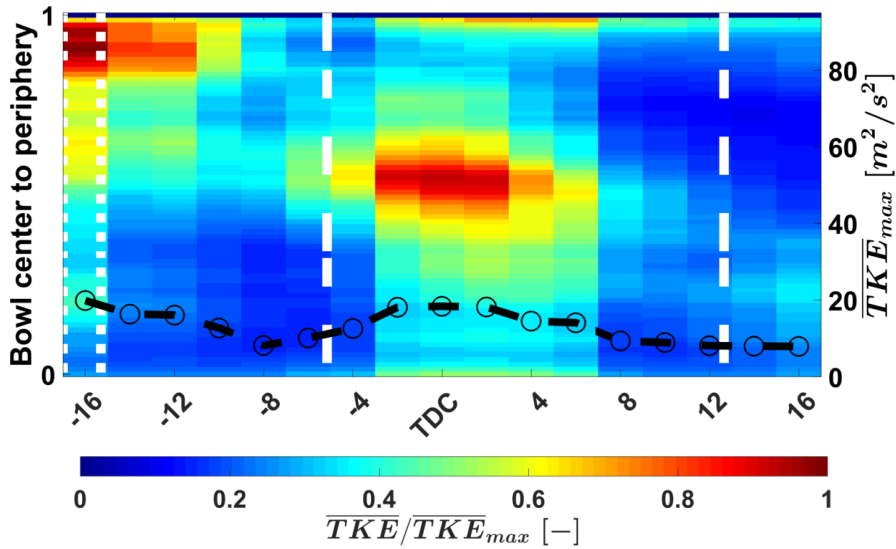


369 *Figure 12. Normalized mean ensemble average velocity module for double injection strategy.*

370 **4.1.3 Triple injection**

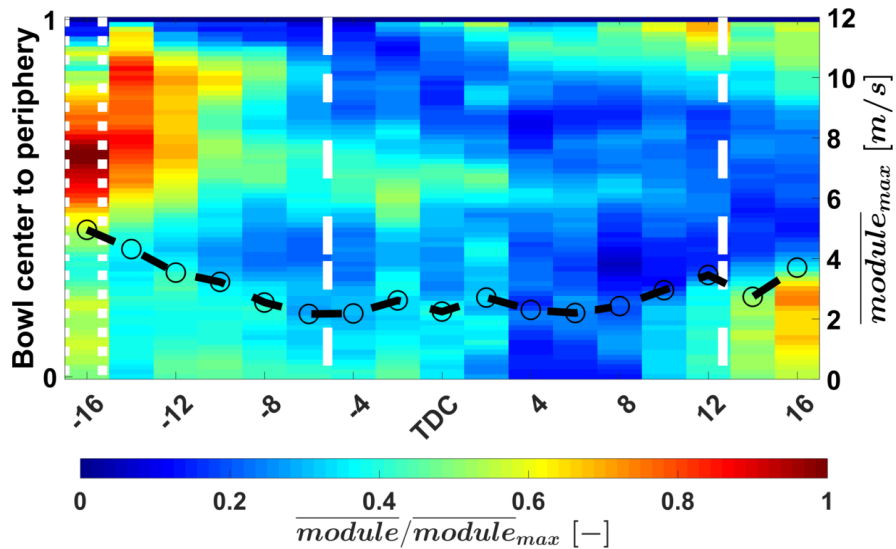
371 As part of the multiple injection approach, a triple injection strategy was also assessed. First
 372 and second injection start as early as 62.5 and 29.5 CAD bTDC with a duration of 2.3 and 1.7
 373 CAD, respectively. The third injection also has a duration of 1.7 CAD with SOI and EOI
 374 marked with dotted lines in Figure 13. Similarly, combustion process is marked at CA10 and
 375 CA90 with dashed lines.

376 Injection and combustion events clearly are major contributors to the TKE field. In this
 377 matter, injection events predominantly introduce high TKE zones at bowl periphery (narrow
 378 red zone) at 16 CAD bTDC. At this same crank-angle, some moderate levels of TKE (mostly
 379 contributions from previous injections) can be observed. As compression progresses, TKE
 380 dissipates up to the point where combustion becomes predominant. After CA10, it promotes
 381 high TKE zones in the middle area of the domain with moderate levels evenly distributed
 382 throughout the rest of the bowl. After 6 CAD aTDC TKE starts to dissipate from bowl
 383 periphery towards its center up to 16 CAD aTDC. Under this injection strategy, the temporal
 384 evolution of the maximum TKE level (black dashed line) shows smooth raises at timings
 385 where injection and combustion dominates the flow.



386 *Figure 13. Spatial and temporal evolution of TKE for triple injection strategy.*

387 In a similar way (in reference to double injection), the flow field under the triple injection
 388 pattern achieves a more evenly distributed velocity field compare to the single injection
 389 pattern. In Figure 14, it is seen that the third injection introduces high velocity levels around
 390 bowl periphery followed by moderate velocity levels uniformly distributed in the domain up to
 391 CA10. During early stages of combustion, higher velocities take place near bowl center
 392 contrary to what is observed between 4 CAD aTDC and CA90. In terms of magnitude,
 393 maximum velocity levels drop after the third injection and slightly increase beyond 8 CAD
 394 aTDC, as seem for all cases.

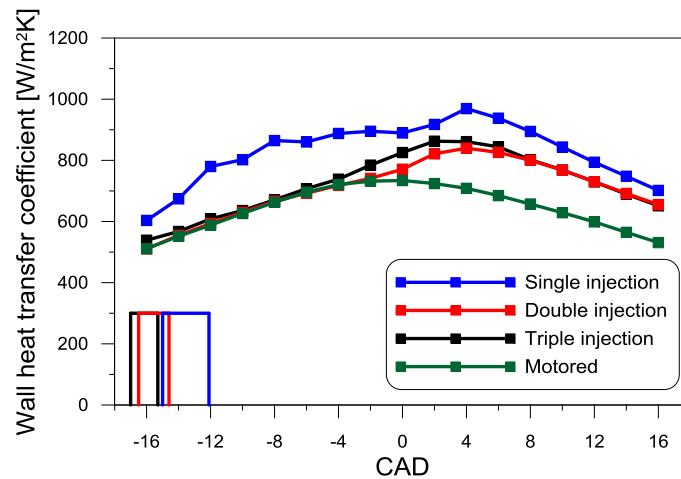


395 *Figure 14. Normalized mean ensemble average velocity module for triple injection strategy.*

396 **4.2 Wall heat transfer coefficient temporal evolution**

397 A modified Woschni's empirical correlation has been used to estimate how wall heat transfer
 398 coefficient changes over time. As stated under the theoretical considerations section of this

399 work, the modified approach leads to a more comprehensive analysis since experimentally
 400 derived flow field parameters are taken into account. Figure 15 shows the
 401 temporal evolution of wall heat transfer coefficient for the motored case as well as for the
 402 firing cases. The injection pulse signal (arbitrary units) has also been plotted as reference. The
 403 injection pulse for the single injection pattern is plotted with a blue rectangle. Meanwhile, the
 404 injection pulse for the last injection in the double and triple injection patterns are plotted with
 405 red and black rectangles, respectively. It is evident that wall heat transfer coefficient temporal
 406 evolution highly depend upon injection pattern and consequently on combustion. To gain
 407 insight into how flow field interacts with fuel injection and combustion and how this
 408 interaction influences wall heat transfer coefficient, an analysis is made based on the
 409 methodology developed around TKE spatial and temporal evolution. The analysis is made in
 410 terms of how the different injection strategies introduce changes into the TKE field and how it
 411 differs from the motored case.



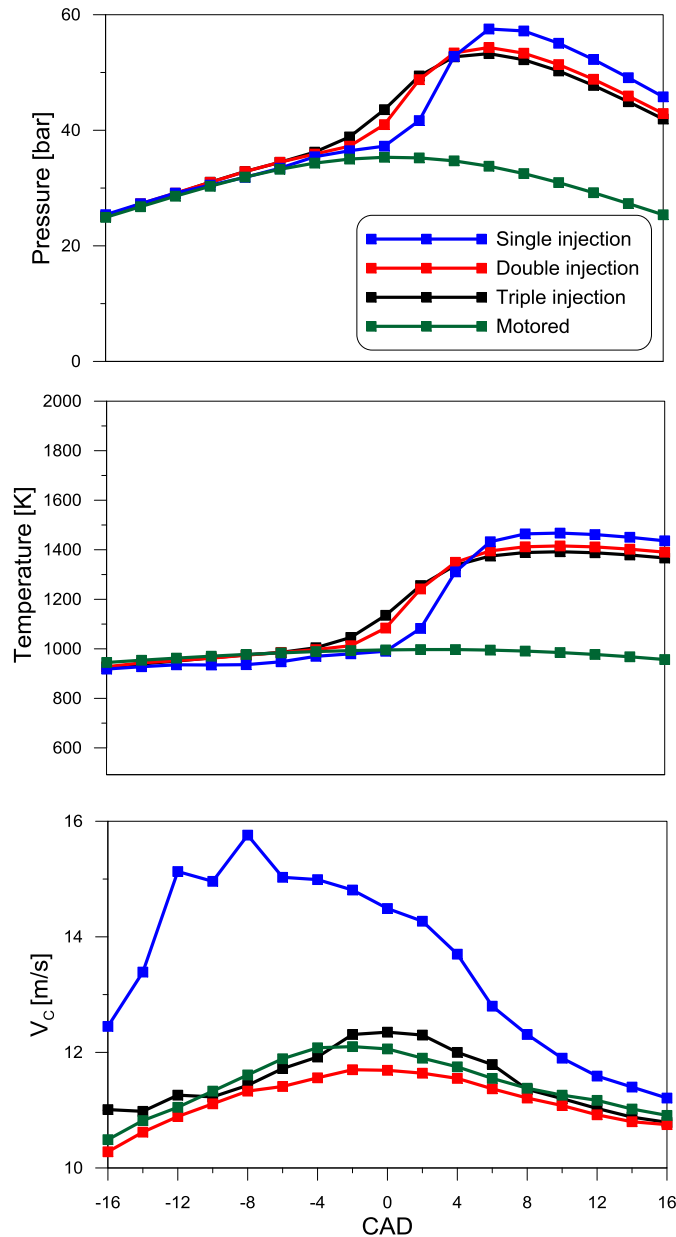
412

413

Figure 15. Wall heat transfer coefficient temporal evolution.

414 Figure 16 (from top to bottom) show in-cylinder pressure, temperature and characteristic
 415 velocity. These three magnitudes are the main terms driving wall heat transfer according to
 416 Eq. (7). For the single injection pattern, in-cylinder thermodynamic conditions does not differ
 417 from the motored case. Nonetheless, the characteristic velocity strongly differs from the
 418 motored case influenced by the injection event. Based on these observations, differences in
 419 wall heat transfer coefficient for the single injection pattern (in reference to the motored case)
 420 are not only driven by combustion (as it was expected), but also for the injection event that
 421 strongly influences the flow field. As for the multiple injection patterns, in-cylinder
 422 thermodynamic conditions remain close to the motored case before the start of combustion.
 423 Unlike the single injection pattern, in these cases the characteristic velocity does not strongly

424 differs from the motored case showing a less strong influence of injection on the flow field.
 425 As a result, wall heat transfer coefficient match the values from the motored case up to the
 426 start of combustion after which changes are mainly driven by thermodynamic conditions.



427

428 *Figure 16. In-cylinder pressure, temperature and characteristic velocity for motored case and three different*
 429 *injection patterns.*

430 To better understand influence of injection and combustion on the flow field (and hence on
 431 wall heat transfer coefficient), TKE fields for the different injection patterns have been
 432 normalized by subtracting the motored case TKE field. This normalization isolates
 433 contributions of injection and combustion allowing for a more suitable comparison. As in the
 434 previously presented TKE fields, in Figure 17 dotted lines are drawn at SOI and EOI.
 435 Similarly, dashed lines mark CA10 and CA90 for combustion. For every injection pattern,

436 wall heat transfer coefficient differences respect to motored case are plotted with black dashed
437 lines.

438 Fuel injection for the single injection pattern seems to cause TKE production at bowl
439 periphery (see top part in Figure 17) up to EOI. Beyond that point, squish induced TKE seems
440 to become predominant at bowl middle zone. Initial stages of combustion after CA10
441 introduce some differences that are quickly dissipated in the entire domain at crank-angles
442 beyond 6 CAD aTDC.

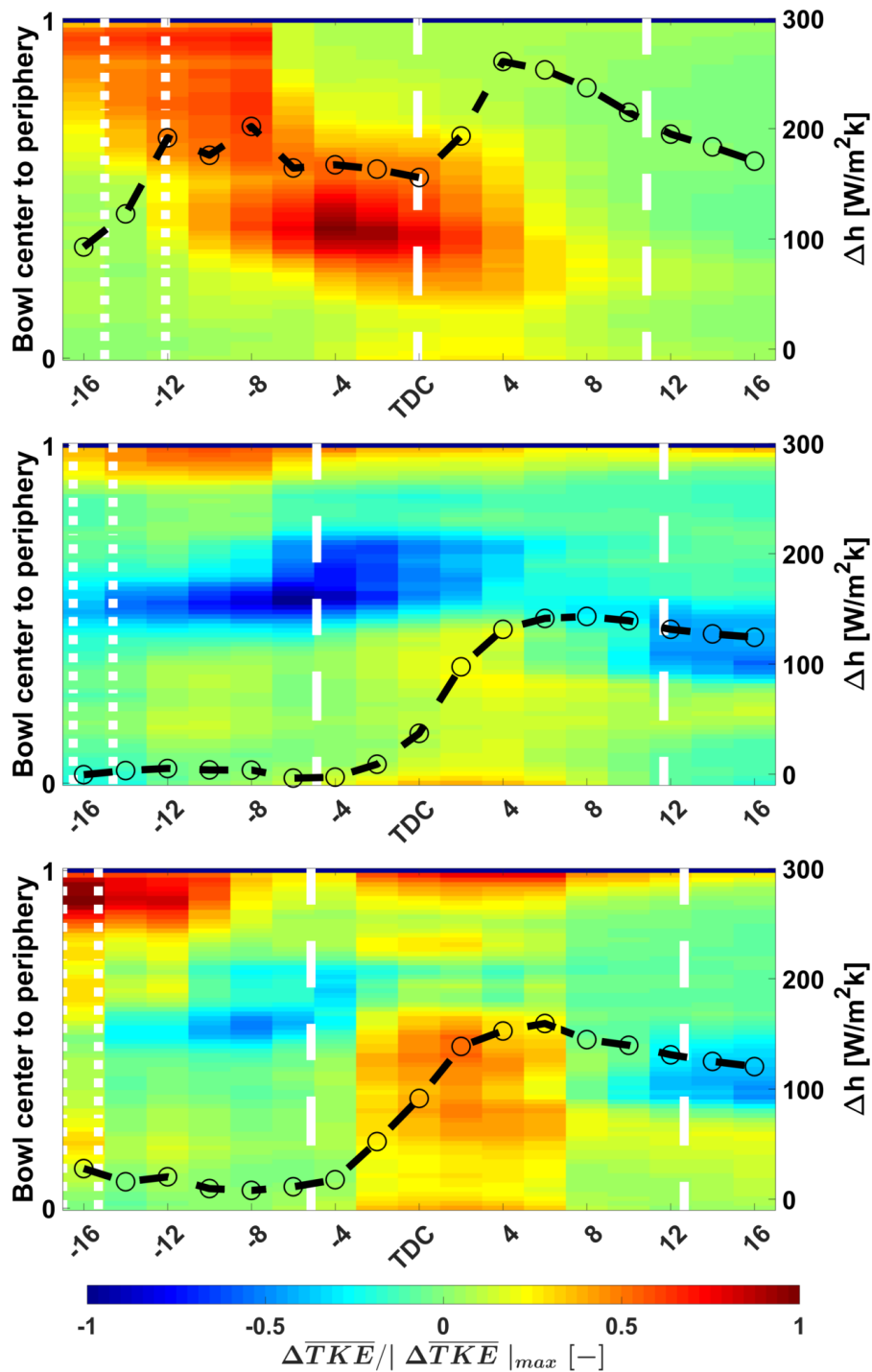
443 As for multiple injection patterns, second injection under the double injection strategy
444 introduces higher TKE levels at bowl periphery compared to motored condition as it was
445 expected. Nonetheless, at bowl middle zone influence of squish is weaker under this injection
446 pattern compared to motored condition (blue area). As previously observed for the single
447 injection case, combustion predominantly influence the TKE field from bowl middle zone
448 towards its center.

449 Once more, the last injection under the triple injection pattern promotes higher TKE levels
450 (compared to motored condition) at bowl periphery. Although, unlike double injection, this
451 third injection event seems to induce similar TKE levels compared to those reached under
452 motored condition. As a result, the blue zone observed at bowl middle zone in the TKE field
453 for double injection is smaller and weaker (in terms of intensity) in the TKE field for triple
454 injection (bottom of the figure). Regarding combustion, it mainly influences bowl middle
455 zone towards its center. However, under this injection pattern combustion also introduces
456 high TKE levels around bowl periphery (narrow red zones). As in the other two injection
457 patterns, TKE is dissipated beyond 6 CAD aTDC.

458 Finally, it is evident that wall heat transfer coefficient temporal evolution is strongly
459 influenced by injection pattern. Black dashed lines show wall heat transfer coefficient
460 differences respect to motored case. Top of Figure 17 shows how the single injection case
461 differs the most respect to the reference case. Both injection and combustion processes induce
462 sharp raises. It is noticeable that after EOI wall heat transfer coefficient remains close to the
463 same level reached at EOI. However, during combustion peak level at 4 CAD aTDC is
464 followed by a continuous drop.

465 For multiple injection cases, contrary to what was observed for the single injection pattern,
466 wall heat transfer coefficient temporal evolution follows the motored case to a greater extend.
467 After the last injection event, and up to CA10, the triple injection pattern seems to differ more

468 if compared to the double injection pattern. It is consistent with greater differences in TKE
 469 levels previously described. As for combustion, just after CA10 wall heat transfer coefficient
 470 for the triple injection pattern shows a slightly sharper increase respect to the double injection.
 471 However, longer combustions (multiple injection cases) come with wall heat transfer
 472 coefficient smoother temporal evolutions and lower peak values compared to the single
 473 injection case.



474 *Figure 17. Overall spatial and temporal evolution of TKE difference between injection strategies and motored*
 475 *case. From top to bottom, single, double and triple injection.*

476 **5 Conclusions**

477 Influence of spatial and temporal evolution of TKE on wall heat transfer coefficient has been
478 studied. TKE has been computed based on PIV data measured in a vertical plane inside the
479 combustion chamber of an optical engine with real bowl geometry. A methodology has been
480 developed to represent maps with TKE temporal and spatial evolution. Wall heat transfer
481 coefficient has also been computed based on a modification of the empirical correlation
482 proposed by Woschni. Finally, results for several fuel injection strategies have been compared
483 against the result from a reference case under motored conditions. Main remarks can be
484 summarized as follows:

- 485 • Under motored conditions, piston movement seems to modulate TKE with peak levels
486 towards TDC and around upper bowl area induced by squish flow stronger influence.
- 487 • Injection events introduce high TKE levels at bowl periphery. On the other hand,
488 during combustion, high TKE levels take place at bowl middle zone.
- 489 • The single injection pattern introduces the greatest differences in the TKE field (in
490 reference to motored case) compared to differences introduced by multiple injection
491 patterns.
- 492 • For the single injection pattern, differences in wall heat transfer coefficient (in
493 reference to motored case) before the start of combustion are mainly driven by
494 changes in the flow field. After the start of combustion, differences in in-cylinder
495 thermodynamic conditions also contribute to reach higher wall heat transfer
496 coefficient levels than in the motored condition case.
- 497 • For multiple injection patterns, the greatest differences in wall heat transfer (in
498 reference to motored case) take place after the start of combustion and are mainly
499 driven by changes in the in-cylinder thermodynamic conditions since no significant
500 changes are observed in the characteristic velocity.

501 **Acknowledgements**

502 The authors gratefully acknowledge the Swedish Energy Agency and the Competence Center
503 for Combustion Processes KCFP. The authors also would like to acknowledge that part of the
504 work has been partially funded by the Spanish government under the grant “José Castillejo”
505 (JC2015/0036). Research leading to this work has also receive funding from Universitat
506 Politècnica de València through the FPI program.

507

508 **References**

- 509 [1] M. Christensen, B. Johansson, The Effect of Combustion Chamber Geometry on HCCI
510 Operation, Sae Tech. Pap. Ser. 2002 (2002). doi:10.4271/2002-01-0425.
- 511 [2] R. Hasegawa, H. Yanagihara, HCCI Combustion in DI Diesel Engine, SAE Tech. Pap.
512 (2003). doi:10.4271/2003-01-0745.
- 513 [3] J.E. Dec, W. Hwang, M. Sjöberg, An Investigation of Thermal Stratification in HCCI
514 Engines Using Chemiluminescence Imaging, (2006). doi:10.4271/2006-01-1518.
- 515 [4] V. Manente, B. Johansson, W. Cannella, Gasoline partially premixed combustion, the
516 future of internal combustion engines?, Int. J. Engine Res. 12 (2011) 194–208.
- 517 [5] V. Manente, Gasoline Partially Premixed Combustion An Advanced Internal
518 Combustion Engine Concept Aimed to High Efficiency, Low Emissions and Low
519 Acoustic Noise in the Whole Load Range, Lund University, 2010.
- 520 [6] J. Benajes, S. Molina, A. García, E. Belarte, M. Vanvolsem, An investigation on RCCI
521 combustion in a heavy duty diesel engine using in-cylinder blending of diesel and
522 gasoline fuels, Appl. Therm. Eng. 63 (2014) 66–76.
523 doi:10.1016/j.applthermaleng.2013.10.052.
- 524 [7] J. Benajes, A. García, J. Monsalve-Serrano, V. Boronat, Gaseous emissions and
525 particle size distribution of dual-mode dual-fuel diesel-gasoline concept from low to
526 full load, Appl. Therm. Eng. 120 (2017) 138–149.
527 doi:10.1016/j.applthermaleng.2017.04.005.
- 528 [8] J. Benajes, A. García, J.M. Pastor, J. Monsalve-Serrano, Effects of piston bowl
529 geometry on reactivity controlled compression ignition heat transfer and combustion
530 losses at different engine loads, Energy. 98 (2016) 64–77.
531 doi:10.1016/j.energy.2016.01.014.
- 532 [9] A.H. Kakaee, A. Nasiri-Toosi, B. Partovi, A. Paykani, Effects of piston bowl geometry
533 on combustion and emissions characteristics of a natural gas/diesel RCCI engine, Appl.
534 Therm. Eng. 102 (2016) 1462–1472. doi:10.1016/j.applthermaleng.2016.03.162.
- 535 [10] M. Izadi Najafabadi, S. Tanov, H. Wang, B. Somers, B. Johansson, N. Dam, Effects of
536 Injection Timing on Fluid Flow Characteristics of Partially Premixed Combustion
537 Based on High-Speed Particle Image Velocimetry, SAE Int. J. Engines. 10 (2017).

- 538 doi:10.4271/2017-01-0744.
- 539 [11] S. Lee, R.D. Reitz, Spray Targeting to Minimize Soot and CO Formation in Premixed
540 Charge Compression Ignition (PCCI) Combustion with a HSDI Diesel Engine, SAE
541 World Congr. 2006 (2006). doi:10.4271/2006-01-0918.
- 542 [12] S. Tanov, Z. Wang, H. Wang, M. Richter, B. Johansson, Effects of Injection Strategies
543 on Fluid Flow and Turbulence in Partially Premixed Combustion (PPC) in a Light
544 Duty Engine, SAE Tech. Pap. (2015). doi:10.4271/2015-24-2455.
- 545 [13] V. Manente, B. Johansson, P. Tunestal, M. Sonder, S. Serra, Gasoline partially
546 premixed combustion: high efficiency, low NO_x and low soot by using an advanced
547 combustion strategy and a compression ignition engine, *Int. J. Veh. Des.* 59 (2012)
548 108–128. doi:10.1504/IJVD.2012.048689.
- 549 [14] G. Kalghatgi, L. Hildingsson, B. Johansson, Low NO_x and low smoke operation of a
550 diesel engine using gasolinelike fuels, *J. Eng. Gas Turbines Power.* 132 (2010) 92803.
- 551 [15] M. Sellnau, M. Foster, K. Hoyer, W. Moore, J. Sinnamon, H. Husted, Development of
552 a Gasoline Direct Injection Compression Ignition (GDCI) Engine, *SAE Int. J. Engines.*
553 7 (2014) 835–851. doi:10.4271/2014-01-1300.
- 554 [16] M. Sellnau, W. Moore, J. Sinnamon, K. Hoyer, M. Foster, H. Husted, GDCI Multi-
555 Cylinder Engine for High Fuel Efficiency and Low Emissions, *SAE Int. J. Engines.* 8
556 (2015) 775–790. doi:10.4271/2015-01-0834.
- 557 [17] M. Sellnau, J. Sinnamon, K. Hoyer, H. Husted, Gasoline Direct Injection Compression
558 Ignition (GDCI) - Diesel-like Efficiency with Low CO₂ Emissions, (2011).
559 doi:10.4271/2011-01-1386.
- 560 [18] Z. Wang, S. Tanov, H. Wang, M. Richter, B. Johansson, M. Alden, High-Speed
561 Particle Image Velocimetry Measurement of Partially Premixed Combustion (PPC) in a
562 Light Duty Engine for Different Injection Strategies, (2015). doi:10.4271/2015-24-
563 2454.
- 564 [19] H. Fridriksson, B.S. aen, S. Hajireza, M.T. aer, CFD Investigation of Heat Transfer in a
565 Diesel Engine with Diesel and PPC Combustion Modes, *SAE Pap.* 2011-01–18 (2011).
566 doi:10.4271/2011-01-1838.
- 567 [20] A. Broatch, P. Olmeda, A. García, J. Salvador-Iborra, A. Warey, Impact of swirl on in-

- 568 cylinder heat transfer in a light-duty diesel engine, *Energy*. (2016) 1–14.
569 doi:10.1016/j.energy.2016.11.040.
- 570 [21] B. Petersen, P. Miles, PIV Measurements in the Swirl-Plane of a Motored Light-Duty
571 Diesel Engine, *SAE Int. J. Engines*. 4 (2011) 1623–1641. doi:10.4271/2011-01-1285.
- 572 [22] K. Kiger, Introduction of Particle Image Velocimetry, (2010) 57.
573 <http://www2.cscamm.umd.edu/programs/trb10/presentations/PIV.pdf>.
- 574 [23] J. Westerweel, Fundamentals of digital particle image velocimetry, *Meas. Sci. Technol.*
575 8 (1997) 1379–1392. doi:10.1088/0957-0233/8/12/002.
- 576 [24] C. Schubert, A. Wimmer, F. Chmela, Advanced Heat Transfer Model for CI Engines,
577 *SAE Tech. Pap. 2005* (2005). doi:10.4271/2005-01-0695.
- 578 [25] G. Woschni, A Universally Applicable Equation for the Instantaneous Heat Transfer
579 Coefficient in the Internal Combustion Engine, *SAE Int.* (1967) 19.
580 doi:10.4271/670931.
- 581 [26] F. Payri, X. Margot, A. Gil, J. Martin, Computational Study of Heat Transfer to the
582 Walls of a DI Diesel Engine, *SAE Tech. Pap. 2005* (2005). doi:10.4271/2005-01-0210.
- 583 [27] J. Chang, O. Güralp, Z. Filipi, D.N. Assanis, T.-W. Kuo, P. Najt, R. Rask, New Heat
584 Transfer Correlation for an HCCI Engine Derived from Measurements of Instantaneous
585 Surface Heat Flux, (2004). doi:10.4271/2004-01-2996.
- 586 [28] T. Morel, P. Blumberg, E. Fort, R. Keribar, Methods for heat analysis and temperature
587 field analysis of the insulated diesel, NASA CR-174783,(DOE/NASA/0342-1). (1984).

588 **Abbreviations**

589 CA50: Crank Angle at 50% mass fraction burned

590 CA90: Crank Angle at 90% mass fraction burned

591 CAD: Crank Angle Degree

592 CI: Compression Ignition

593 EA: Ensemble average

594 EOI: End of Injection

595 *h*: Wall heat transfer coefficient

- 596 HCCI: Homogenous Charge Compression Ignition
- 597 LTC: Low Temperature Combustion
- 598 NO_x: Nitrogen Oxides
- 599 PIV: Particle Image Velocimetry
- 600 PPC: Partially Premixed Combustion
- 601 RCCI: Reactivity Controlled Compression Ignition
- 602 SI: Spark Ignition
- 603 SOC: Start of Combustion
- 604 SOI: Start of Injection
- 605 TDC: Top Dead Center
- 606 TKE: Turbulent Kinetic Energy
- 607 V_c : Characteristic velocity
- 608 V_{module} : Module of velocity in the vertical plane
- 609 V_{tang} : Tangential component of velocity